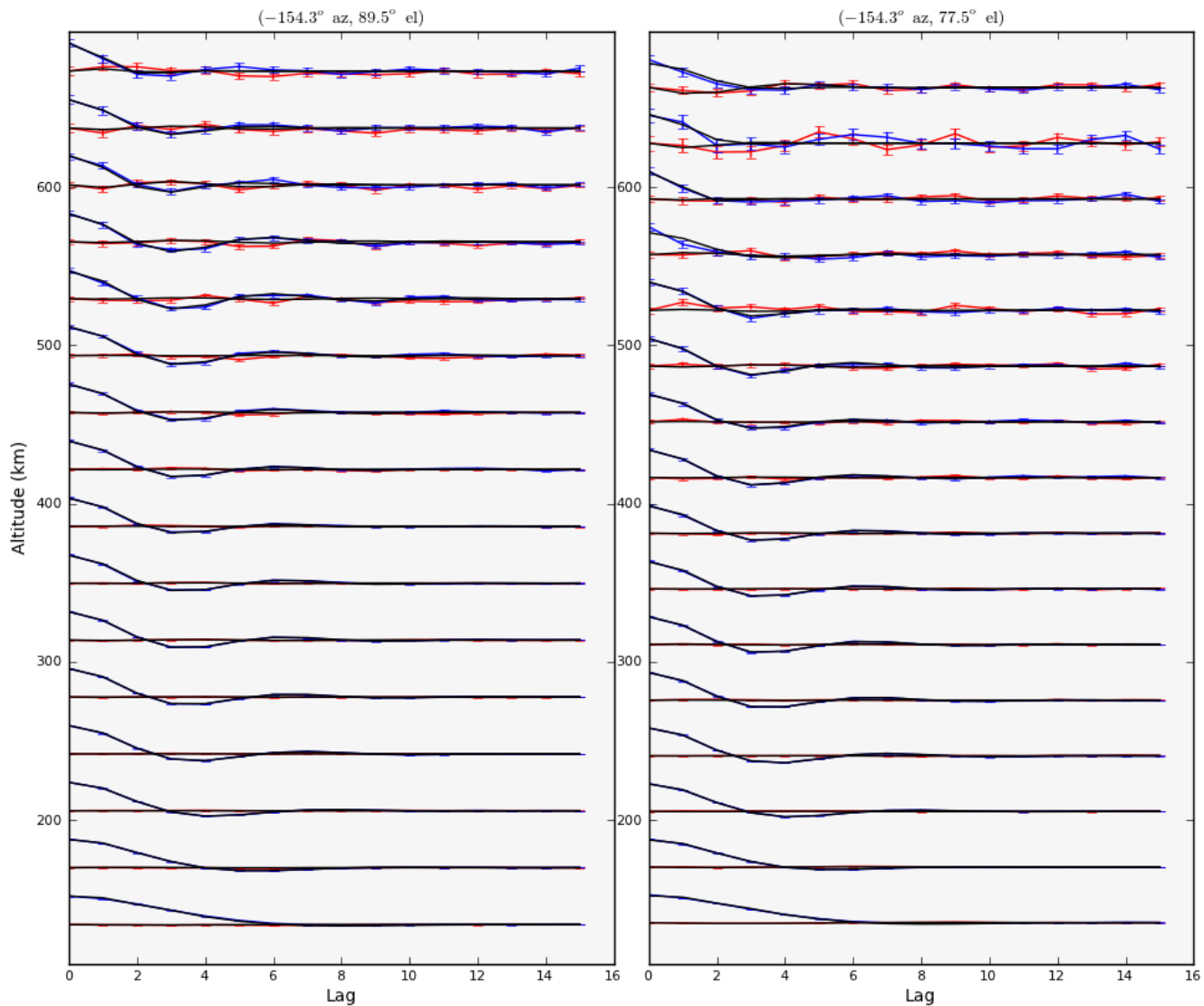




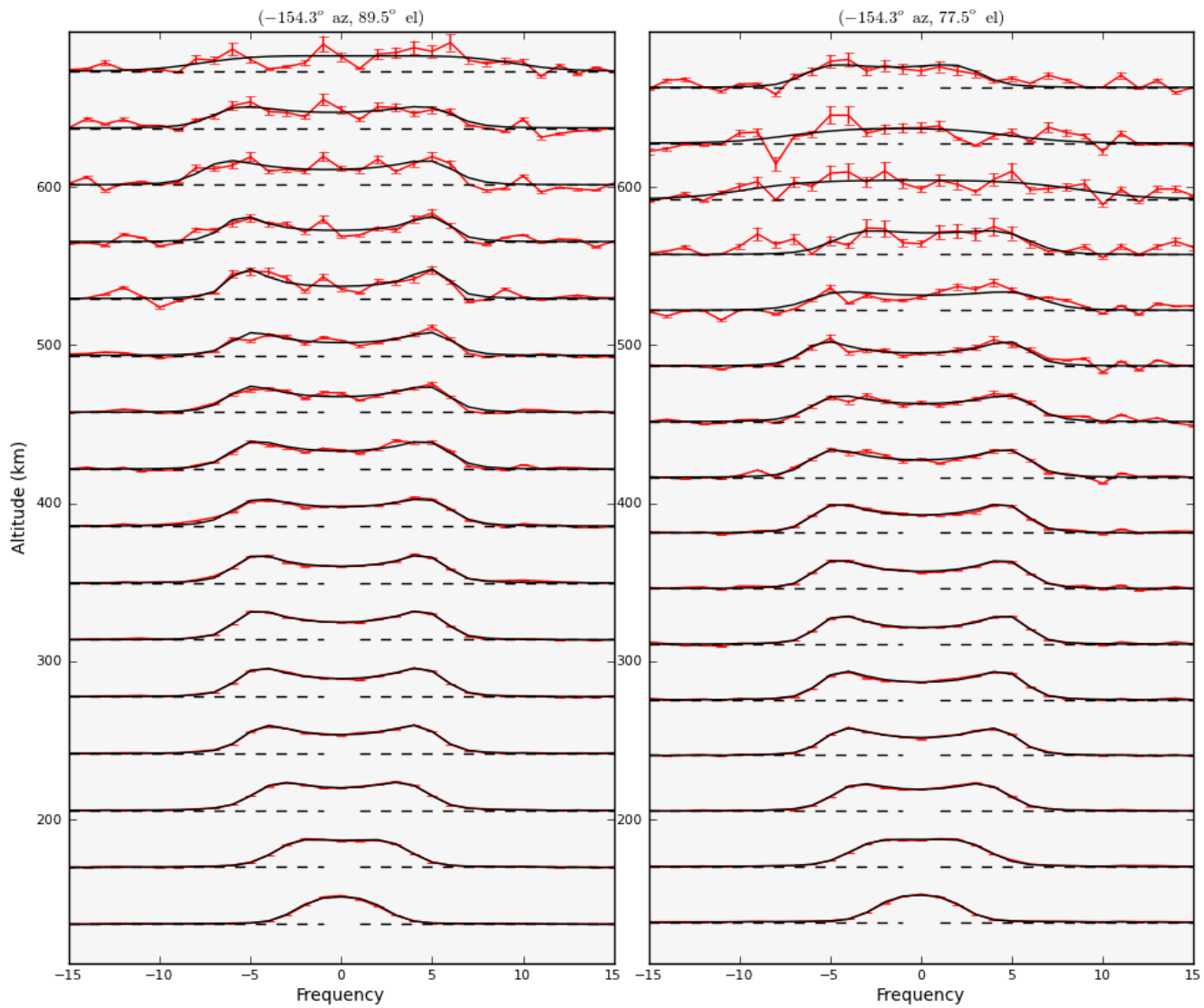
**Incoherent Scatter Radar Spectra
(and other stuff)**

8-1-2012 3.5143 53.0

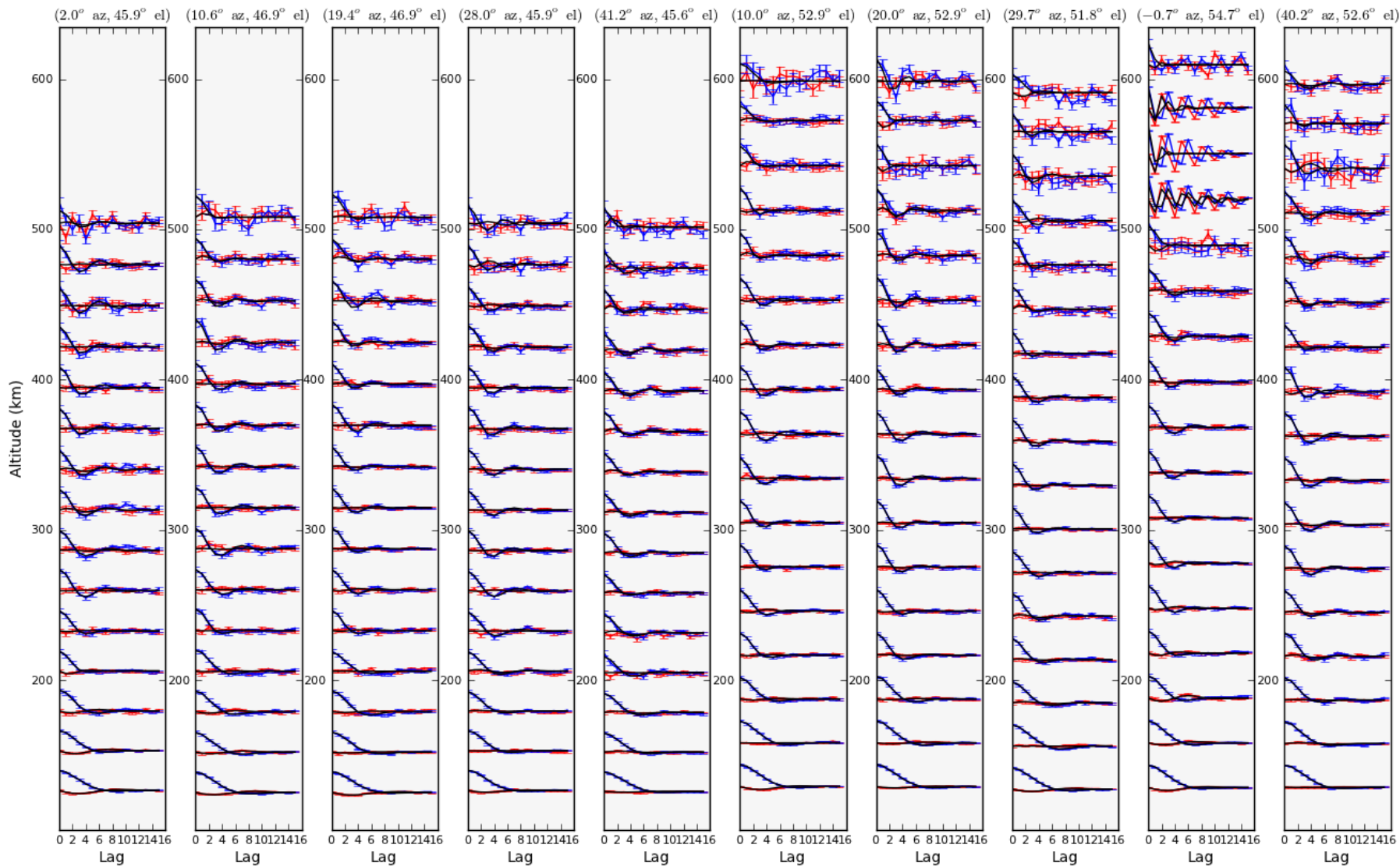
Additional Plots Available



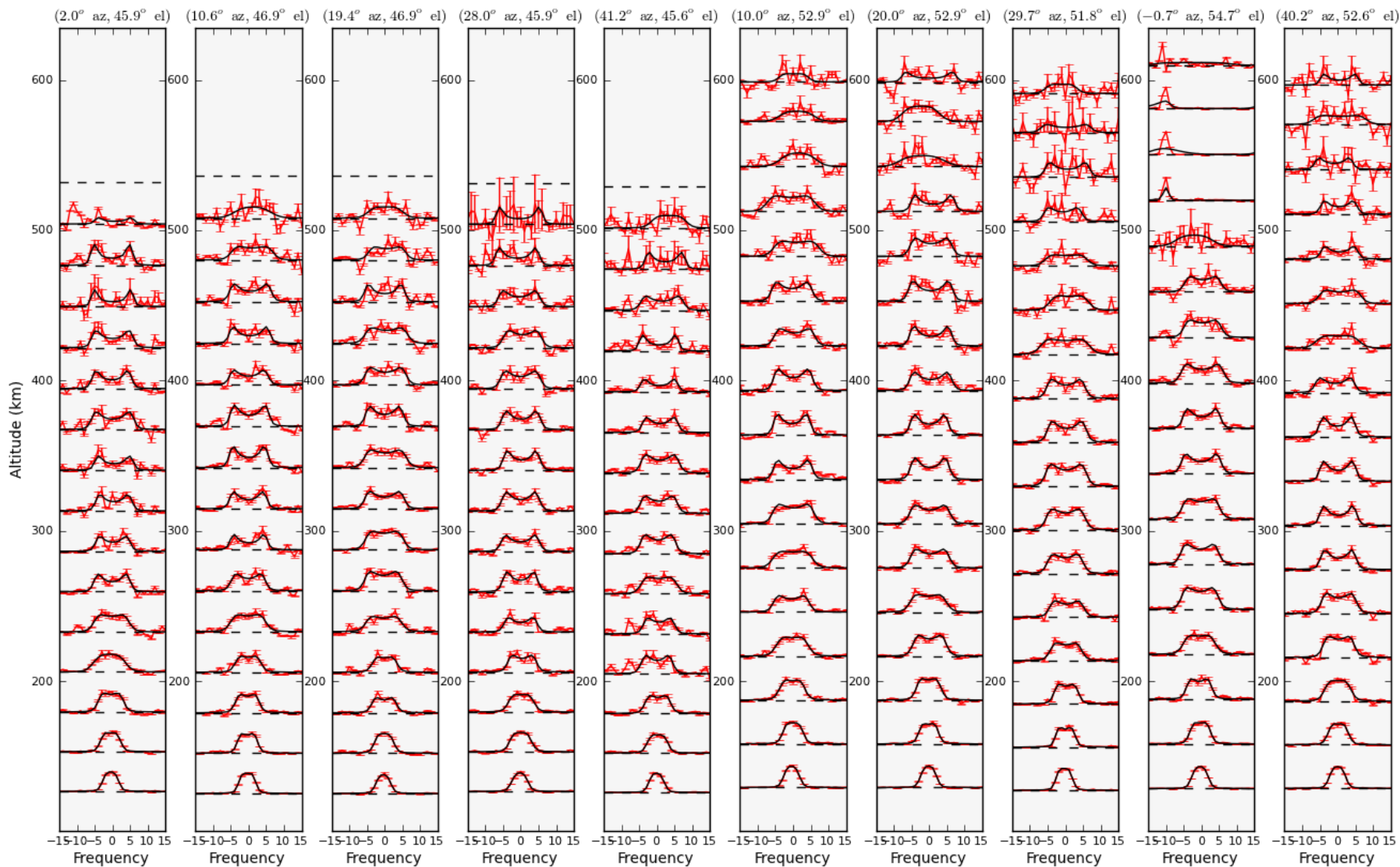
8-1-2012 3.514-3.534 UT



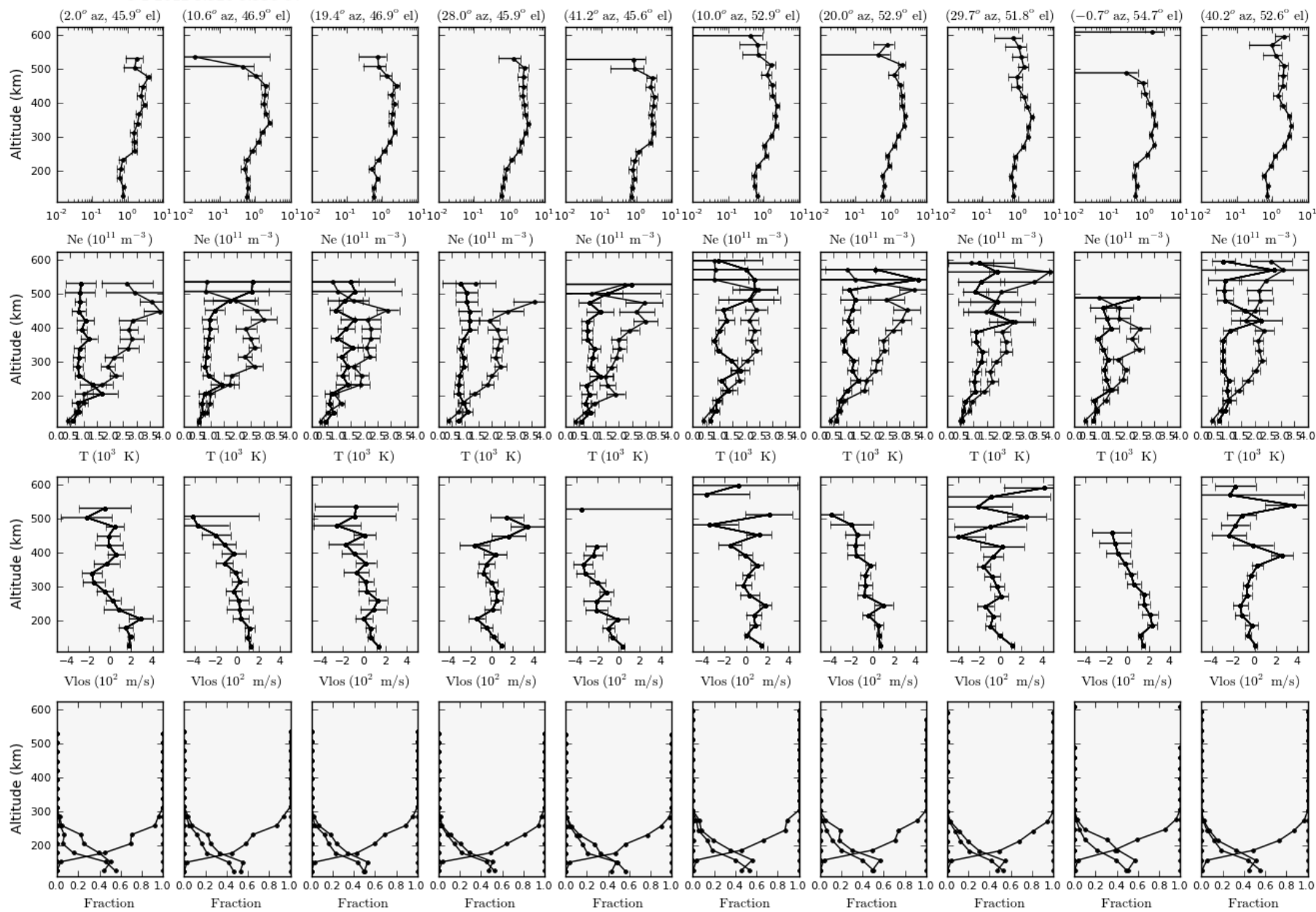
8-1-2012 9.520-9.538 UT



8-1-2012 9.520-9.538 UT



8-1-2012 9.520-9.538 UT



N_e From Ion Line vs. Plasma Line

Ion Line

$$N_e(r) = (1 + k^2 \lambda_D^2)(1 + k^2 \lambda_D^2 + T_r) \frac{C_s}{P_T \tau_p} P_r(r)$$

$$k = \frac{2\pi}{\lambda_{TX}} + \frac{2\pi}{\lambda_{RX}}$$

$$\lambda_D = \sqrt{\frac{\epsilon_0 k_B T_e}{q_e^2 N_e}}$$

So N_e is proportional to the received power from the ion line (a noise-like signal) and is affected to a very minimal extent by λ_D .

N_e From Ion Line vs. Plasma Line

Plasma Line

$$f_{pl}^2 \approx f_{pe}^2 + \frac{3k^2}{4\pi^2} \frac{k_B T_e}{m_e}$$

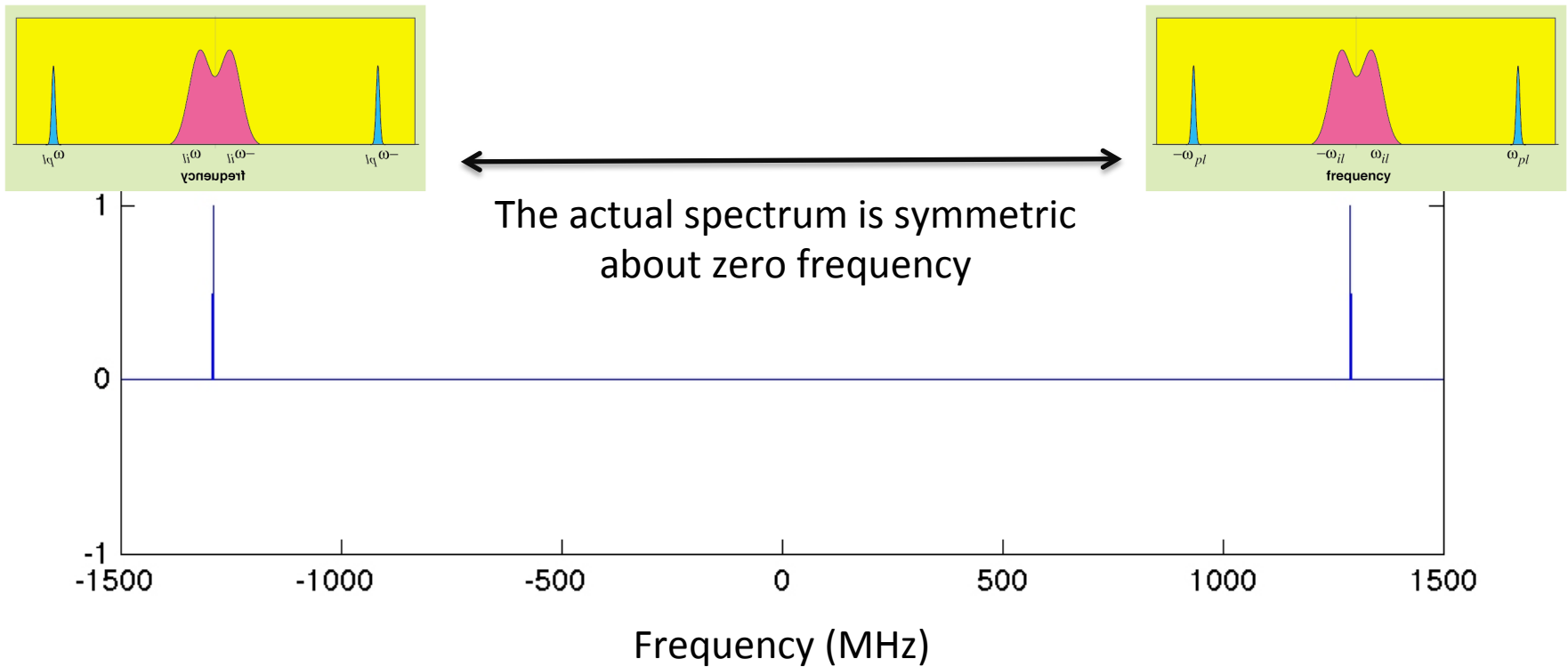
$$f_{pe} = \frac{1}{2\pi} \sqrt{\frac{N_e q_e^2}{m_e \epsilon_0}}$$

$$k = \frac{2\pi}{\lambda_{TX}} + \frac{2\pi}{\lambda_{RX}}$$

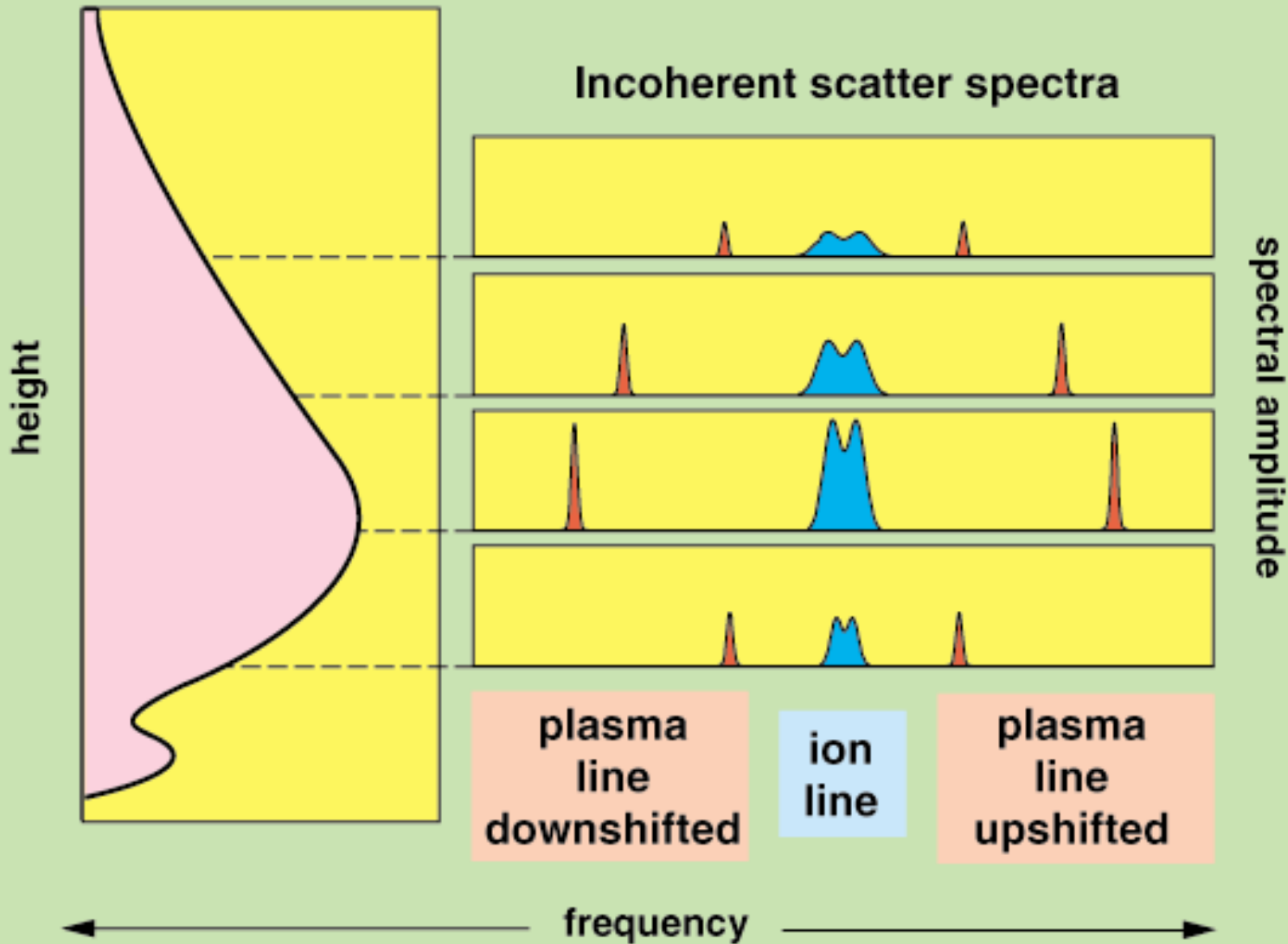
So N_e is determined by T_e and the frequency of the plasma line. Note that here λ_{RX} and λ_{TX} are significantly different for up-shifted vs. down-shifted plasma lines!

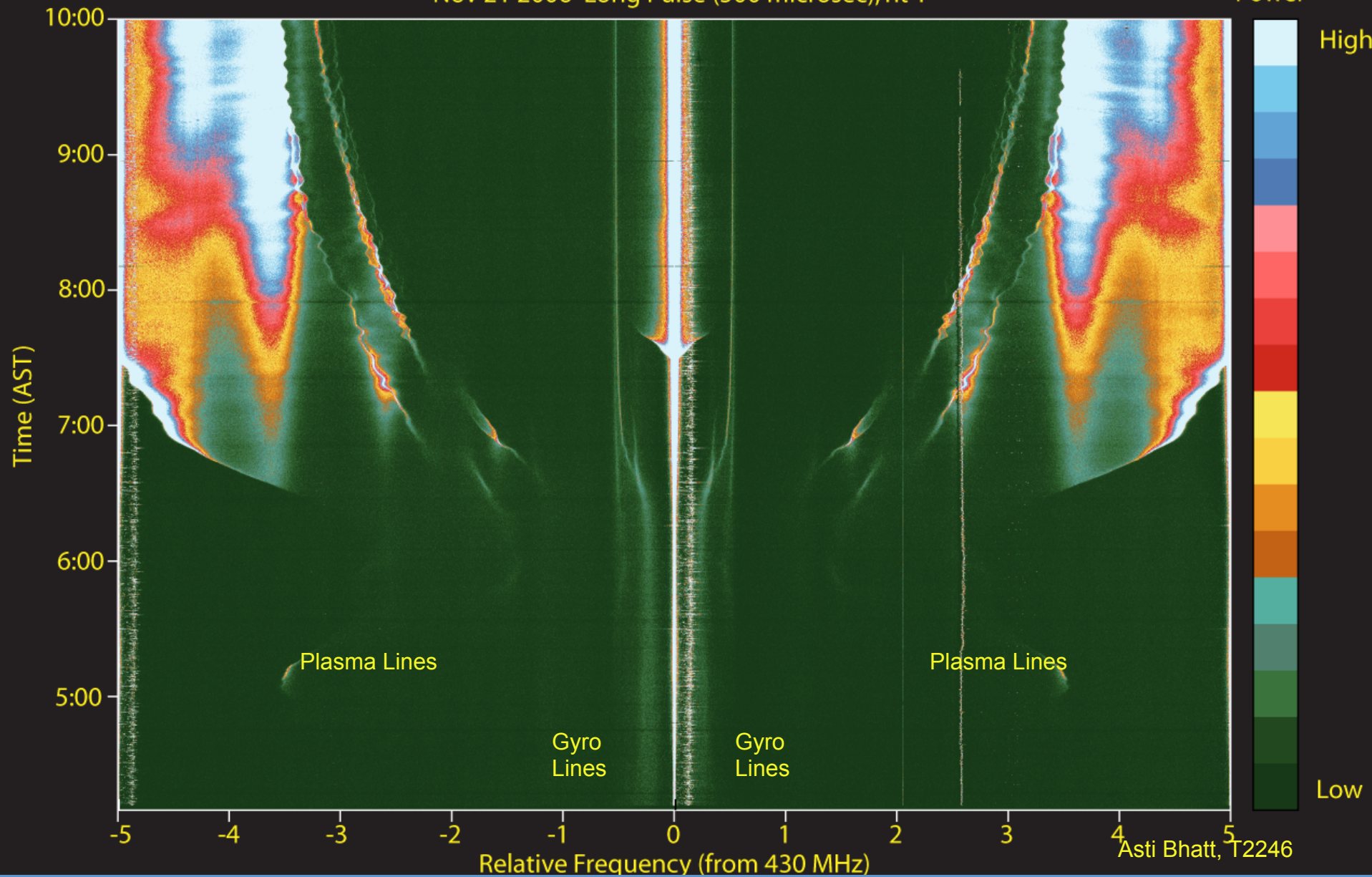
IS Received Spectrum

The individual sides are **not** symmetric about their center frequencies!



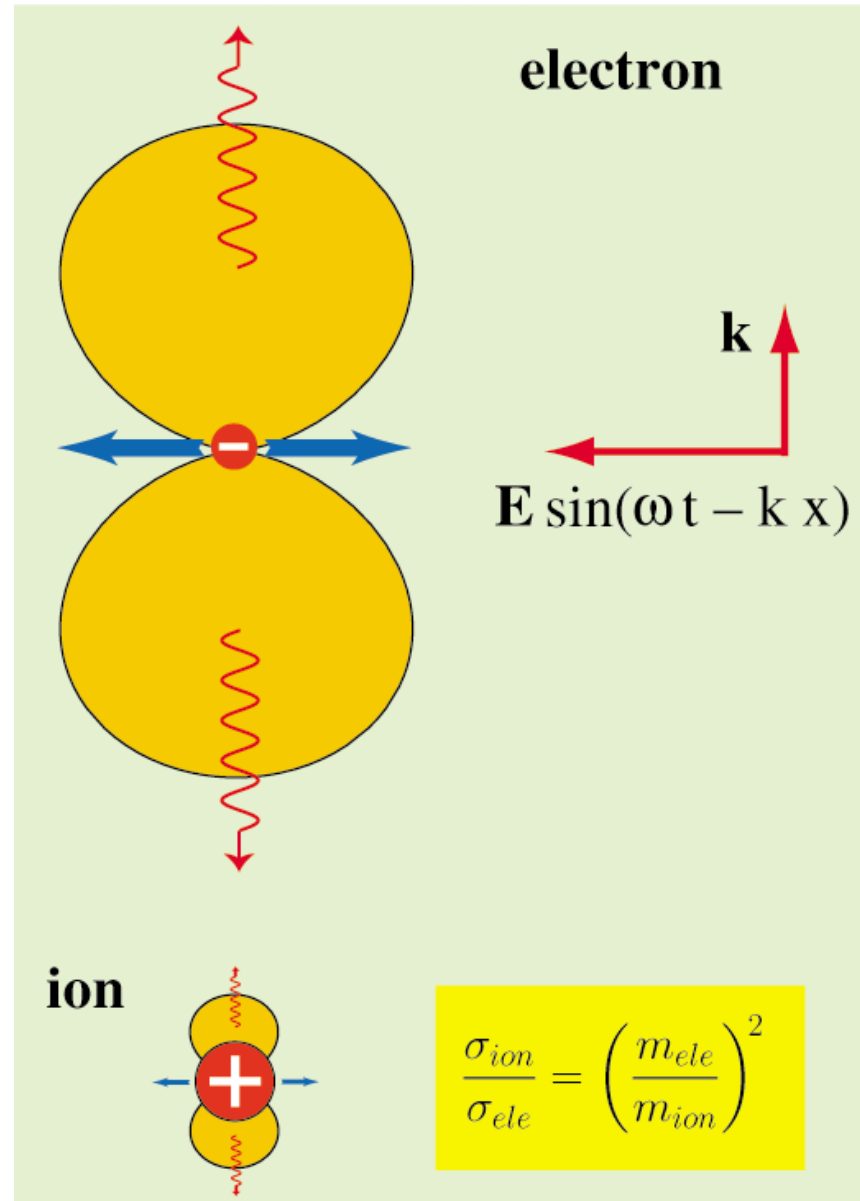
electron density profile





Arecibo Sensitivity: The 305 m dish, 2.5 MW of power, and T_{sys} of about 80 K (condition dependent) provide high time resolution on even weak features such as the gyro line. The data above are centered on the E region. The strong plasma line after sunrise is “leakage” from the low F region. The complicated behavior of the gyro line is probably due to multiple layers, but is not completely understood.

Thomson Scattering



Thomson Scattering

$$E_x = E_0 e^{j(\omega t - kx)}$$

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$$

$$v_x = -j \frac{q_e E_0}{m_e \omega} e^{j\omega t}$$

$$E_\phi = \frac{\mu_0 q_e^2}{4\pi m_e} \frac{\sin \phi}{r} e^{-jkr} E_0$$

$$\begin{aligned} \sigma_e &= 4\pi \left(\frac{\mu_0 q_e^2}{4\pi m_e} \right)^2 \sin^2 \phi = 4\pi r_e^2 \sin^2 \phi \\ &\approx 10^{-28} \sin^2 \phi \quad (\text{m}^2) \end{aligned}$$

Plasma Response to a Stationary Test Charge

For a neutral Maxwellian plasma ($q\phi \ll k_B T$), with a static potential ϕ , the density fluctuations will look like

$$n_{e,i} = n_0 \exp\left(-\frac{q_{e,i}\phi}{k_B T_{e,i}}\right) \approx n_0 \left(1 - \frac{q_{e,i}\phi}{k_B T_{e,i}}\right)$$

Solving Poisson's equation for a test charge at the origin, we obtain

$$\Delta n_e = n_e - n_0 \approx n_0 \frac{q_e^2}{4\pi\epsilon_0 r} \exp\left(-\frac{r}{\lambda}\right)$$
$$\frac{1}{\lambda^2} = \frac{1}{\lambda_{Di}^2} + \frac{1}{\lambda_{De}^2} = -n_0 \frac{q_e^2}{\epsilon_0 k_B} \left(\frac{1}{T_i} + \frac{1}{T_e}\right)$$

where λ_D is the Debye length of the plasma.

Plasma Response to a Stationary Test Charge (cont' d)

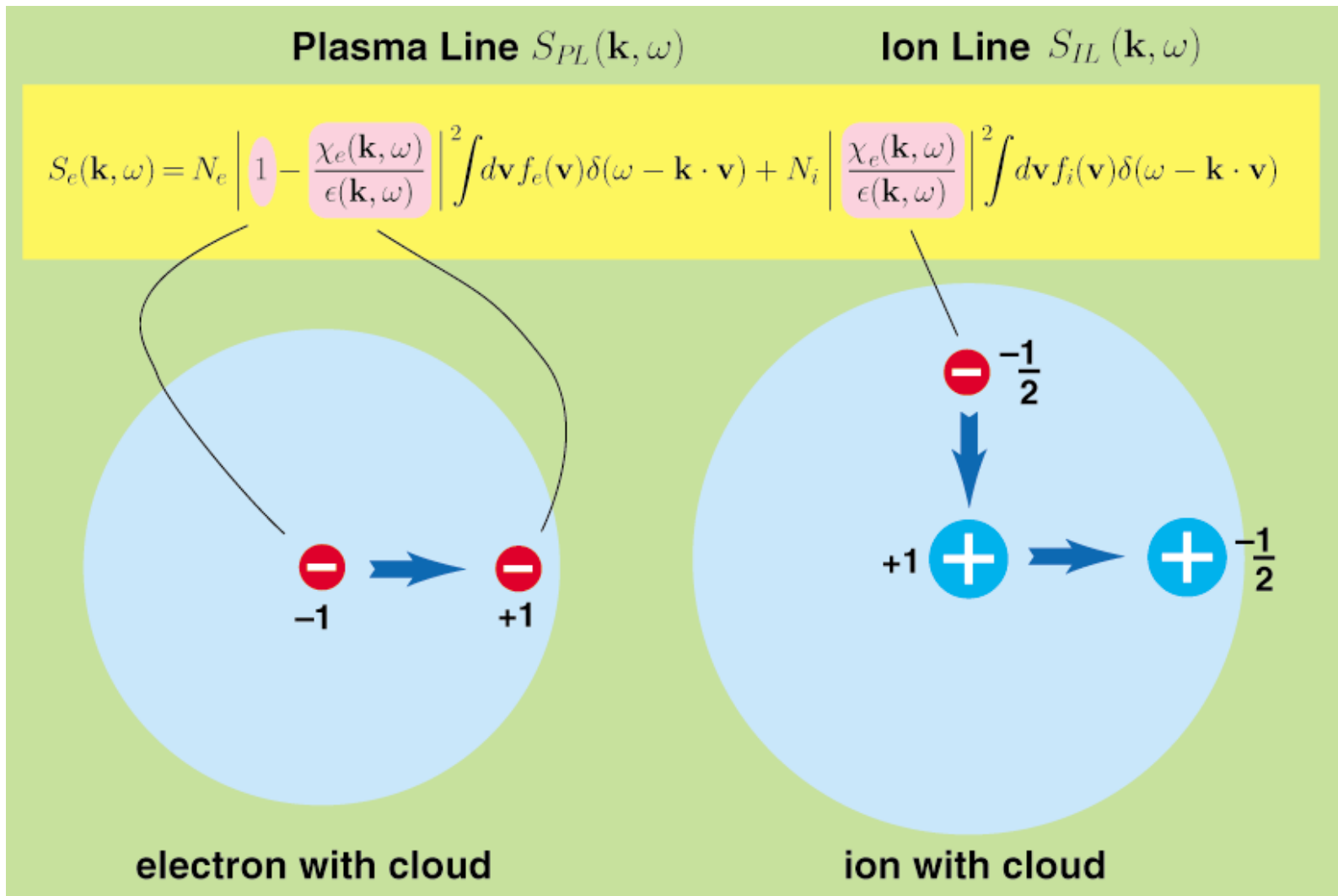
Integrating over all space, we can calculate the total number of additional electrons and ions due to this test charge

$$\iiint \Delta n_e dV = \frac{T_i}{T_e + T_i} = \frac{1}{1 + T_r}$$

$$\iiint \Delta n_i dV = -\frac{T_i}{T_e + T_i} = -\frac{1}{1 + T_r}$$

So the test ion is neutralized by attracting half of an electron and repelling half of an ion when the ion and electron temperatures are equal.

Plasma Response to a Stationary Test Charge (cont' d)



What



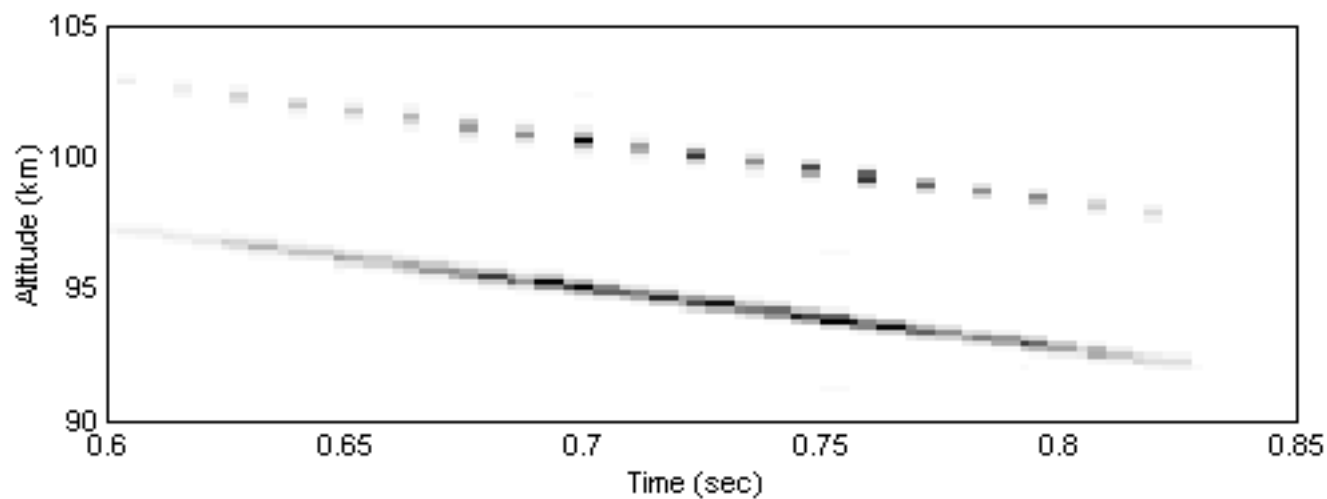
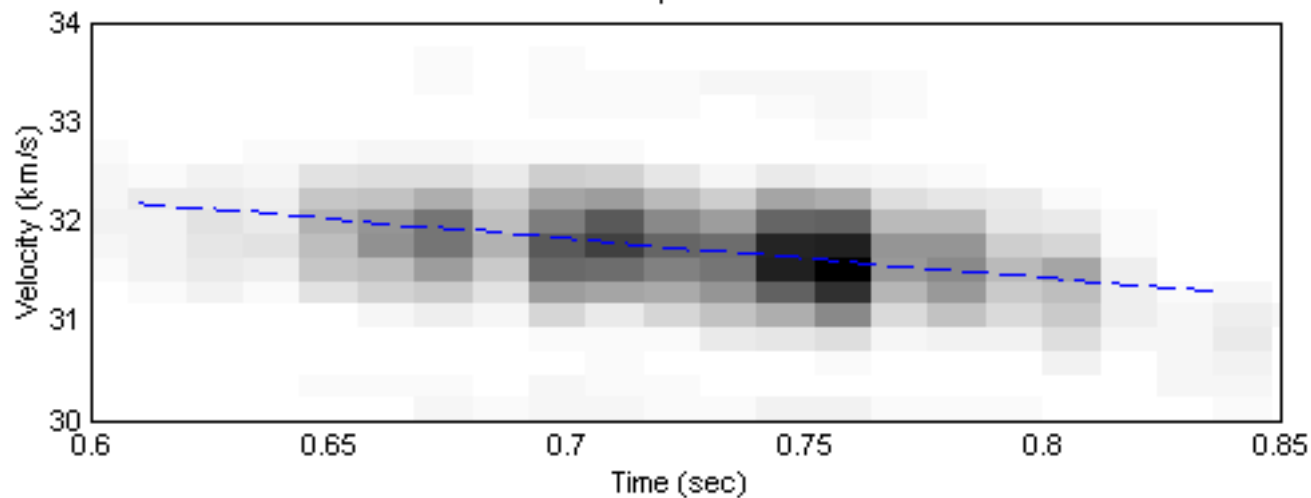
ong?



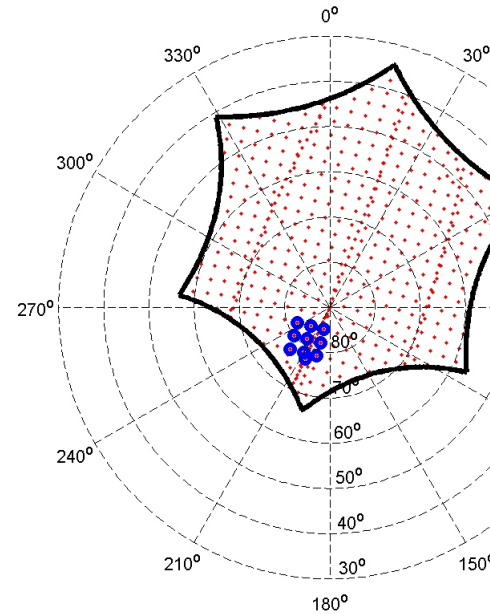
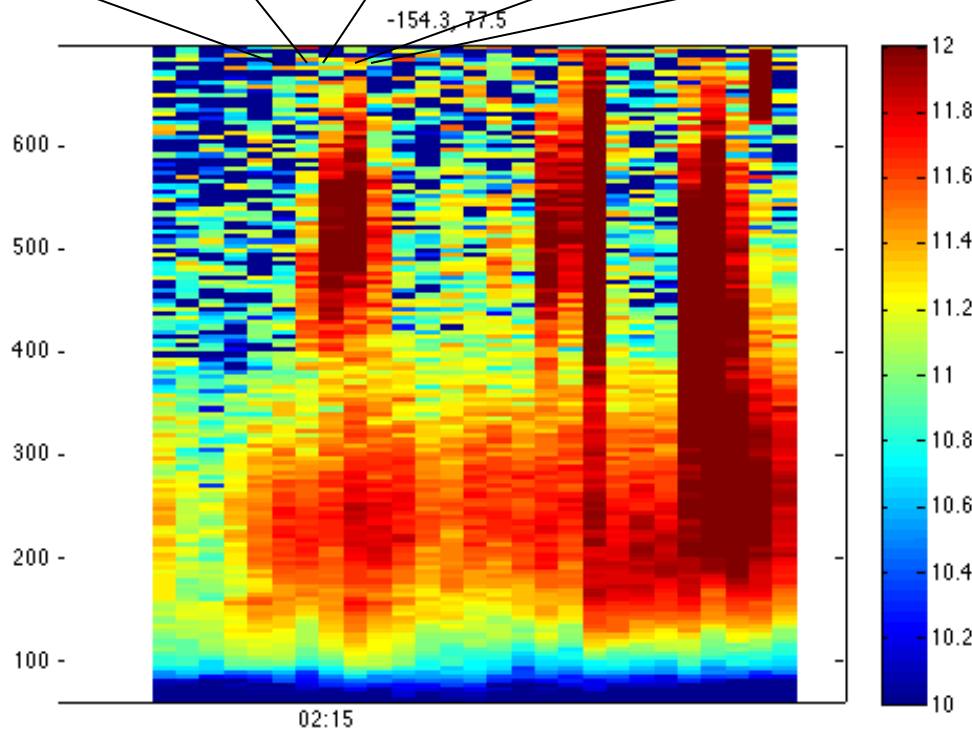
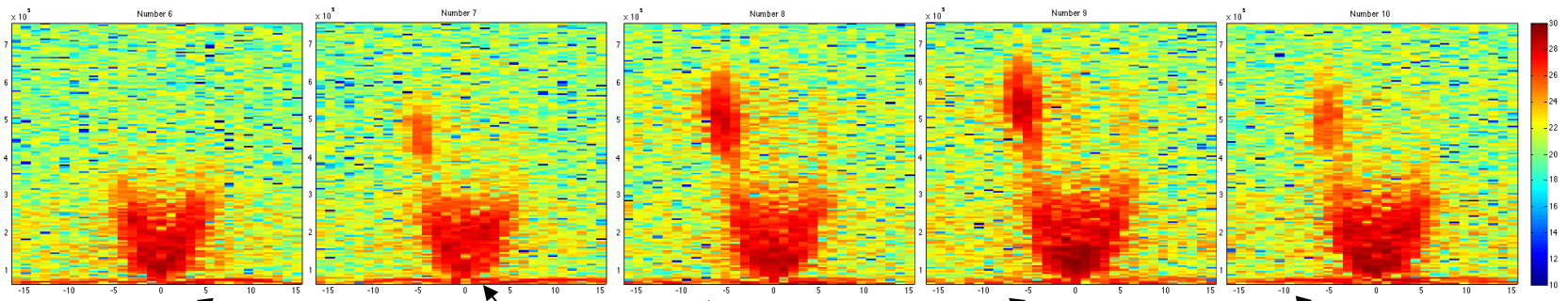
Tao Berman, 19

Echoes from Meteors

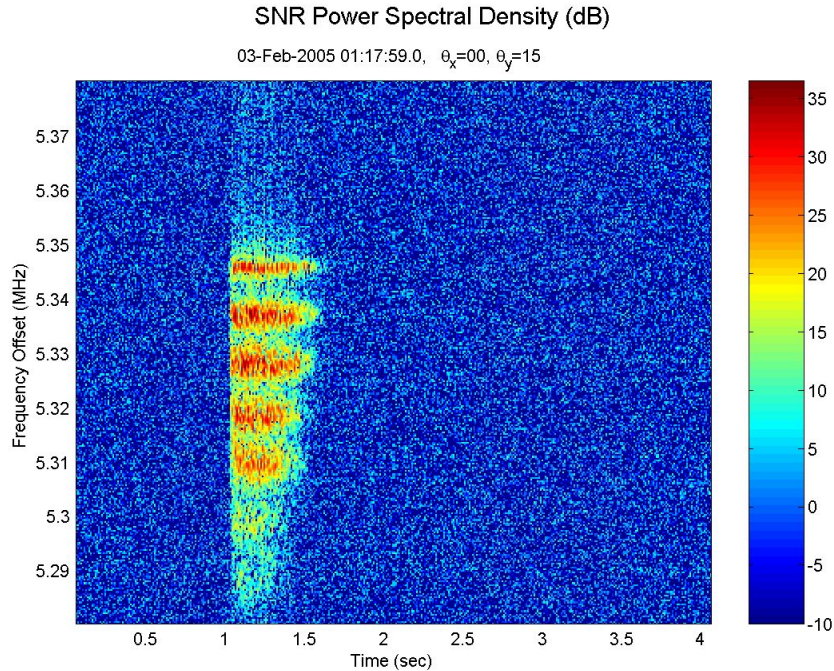
20010607, 135249.2 UT



Naturally Enhanced Ion Acoustic Lines

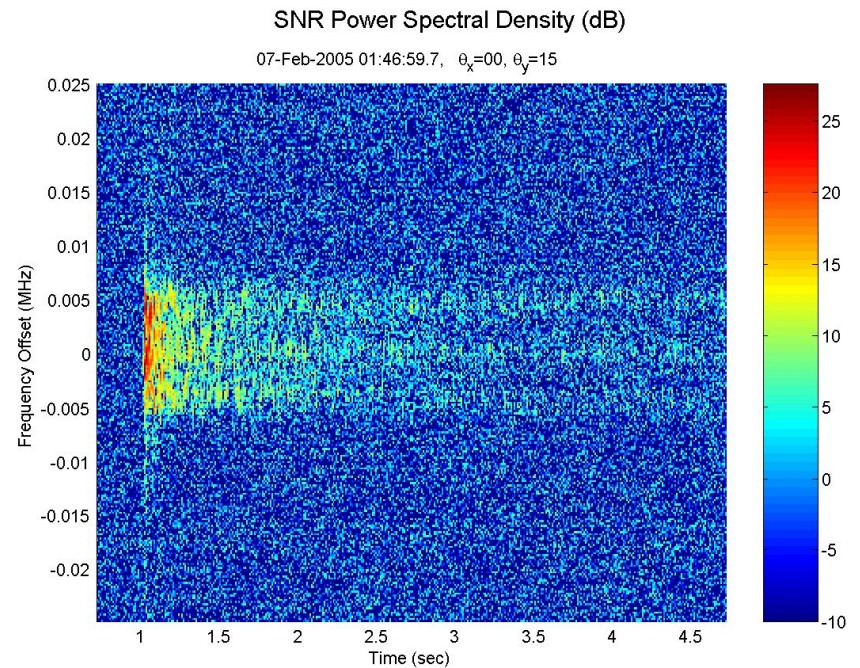


Ionospheric Heating, HAARP, Feb 2005



Enhanced Ion Line

$F_{\text{heater}}=4.25$ MHz, O-mode, CW, on at t=1 sec



Enhanced Plasma Line

$F_{\text{heater}}=5.35$ MHz, O-mode, CW, on at t=1 sec