



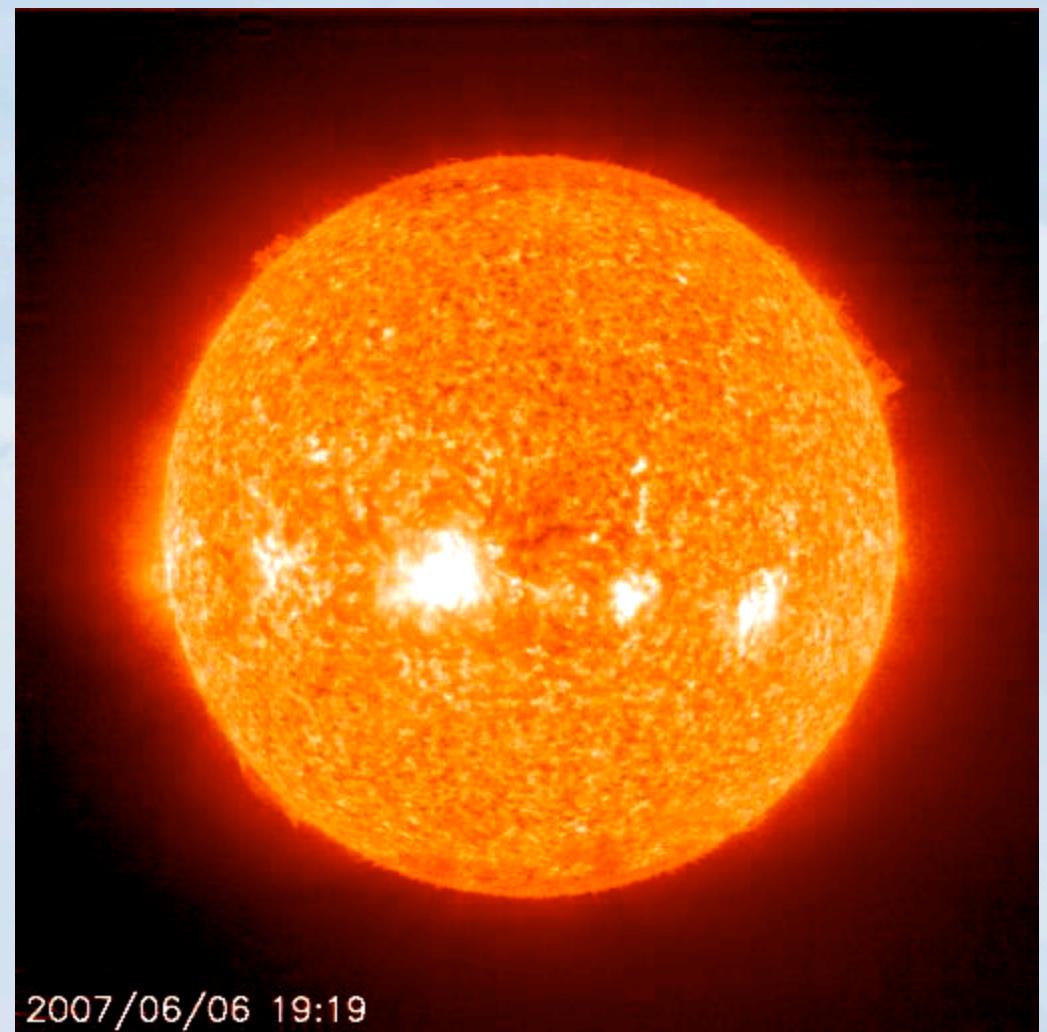
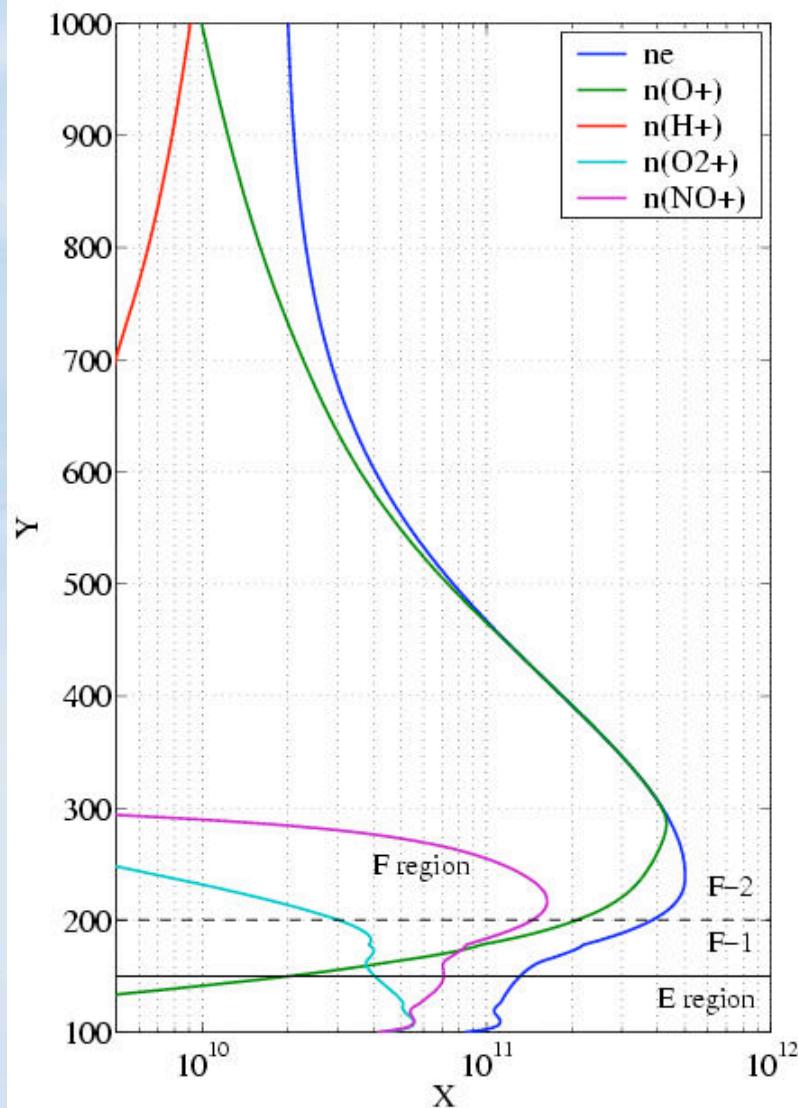
The short introduction to Incoherent Scatter (IS) Theory

Anja Strømme
SRI International



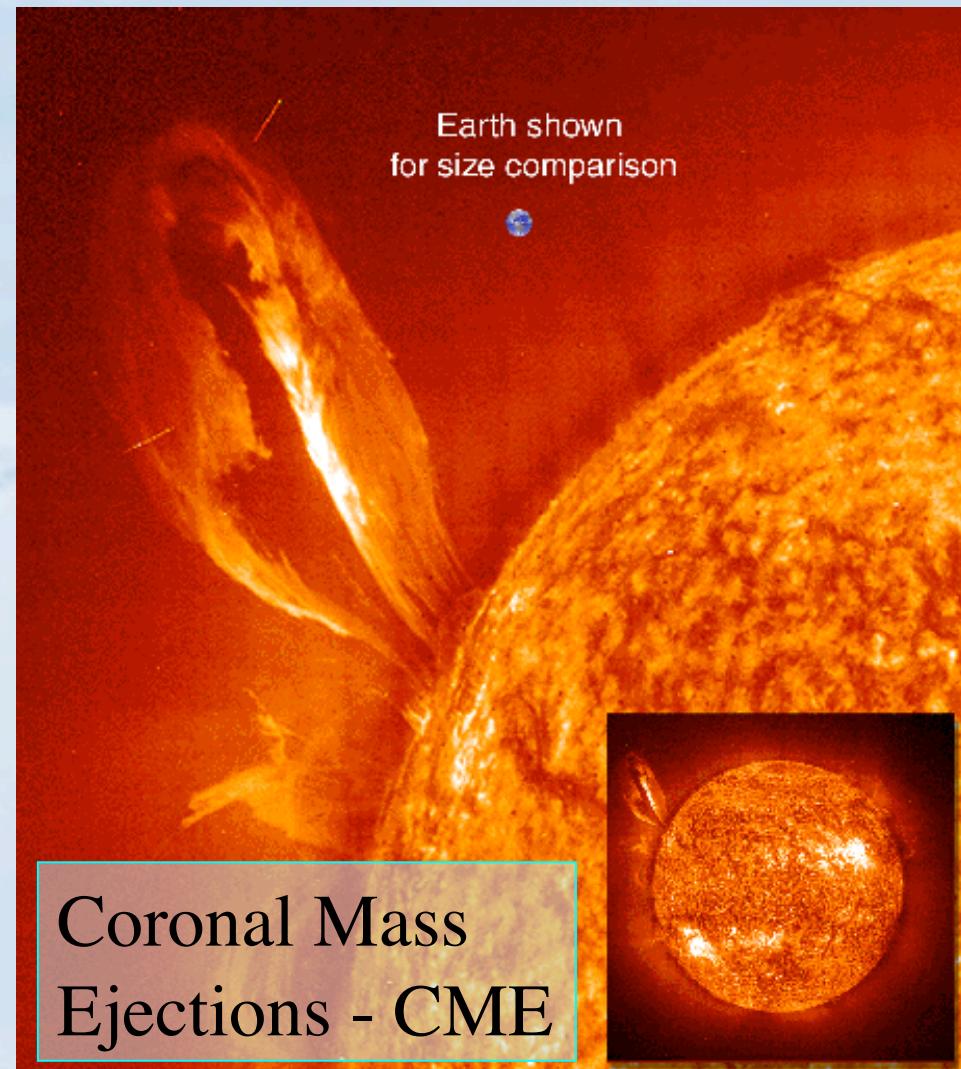
First: We need an
Ionsphere...

The Earth's Ionosphere

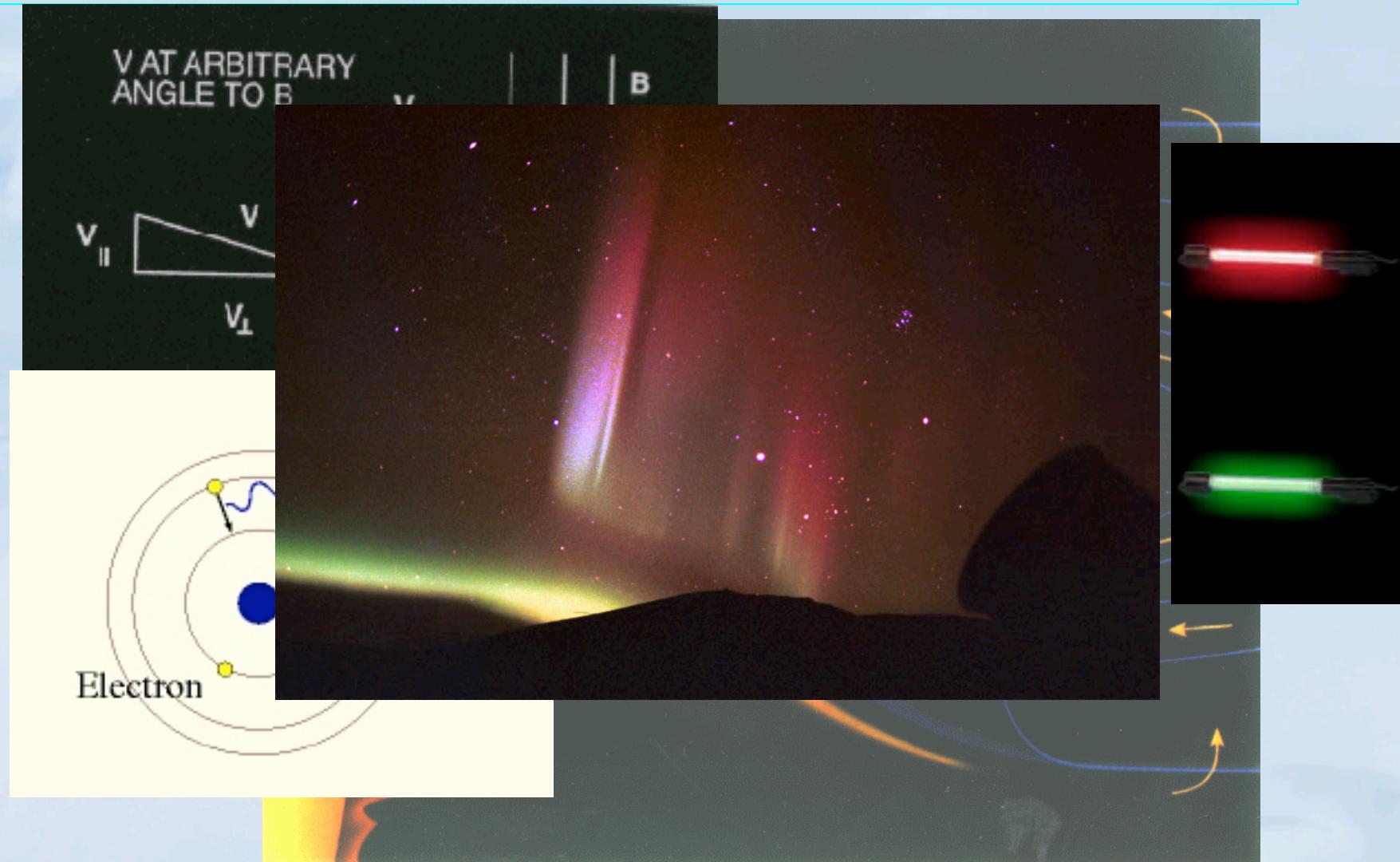


Additional ionization at high latitudes

- The Sun evaporates at supersonic velocities
- The Sun's atmosphere is very hot and ionises the escaping gas
- The expanding 'Solar Wind' carries the surface magnetic fields with it
- The Solar Wind is frequently distorted by huge eruptions which generate vast shock waves



At high latitudes electron (and proton) with solar wind origin creates additional ionization

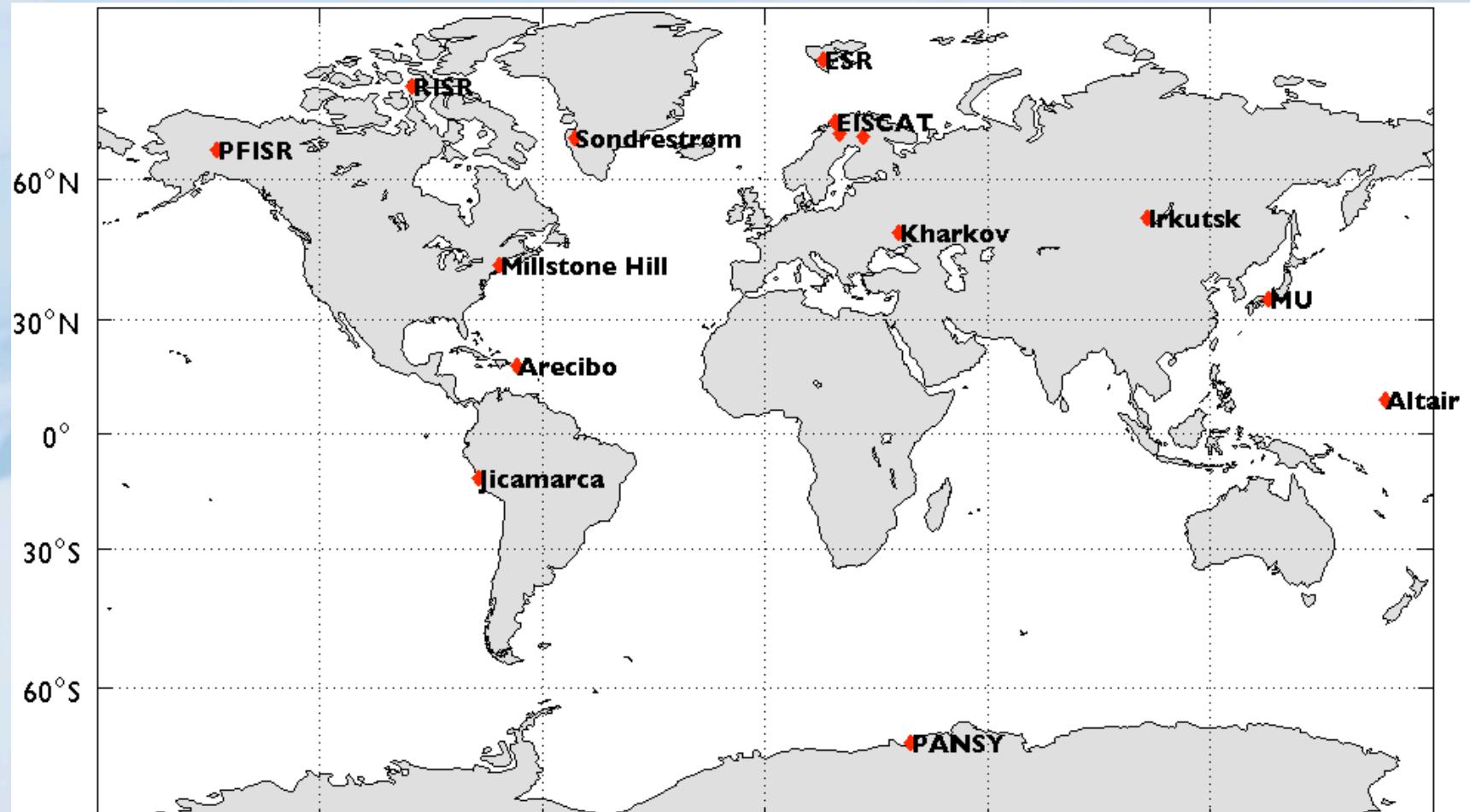




Now we have an ionosphere -
let's add the Incoherent Scatter
Radar (ISR) to probe it!

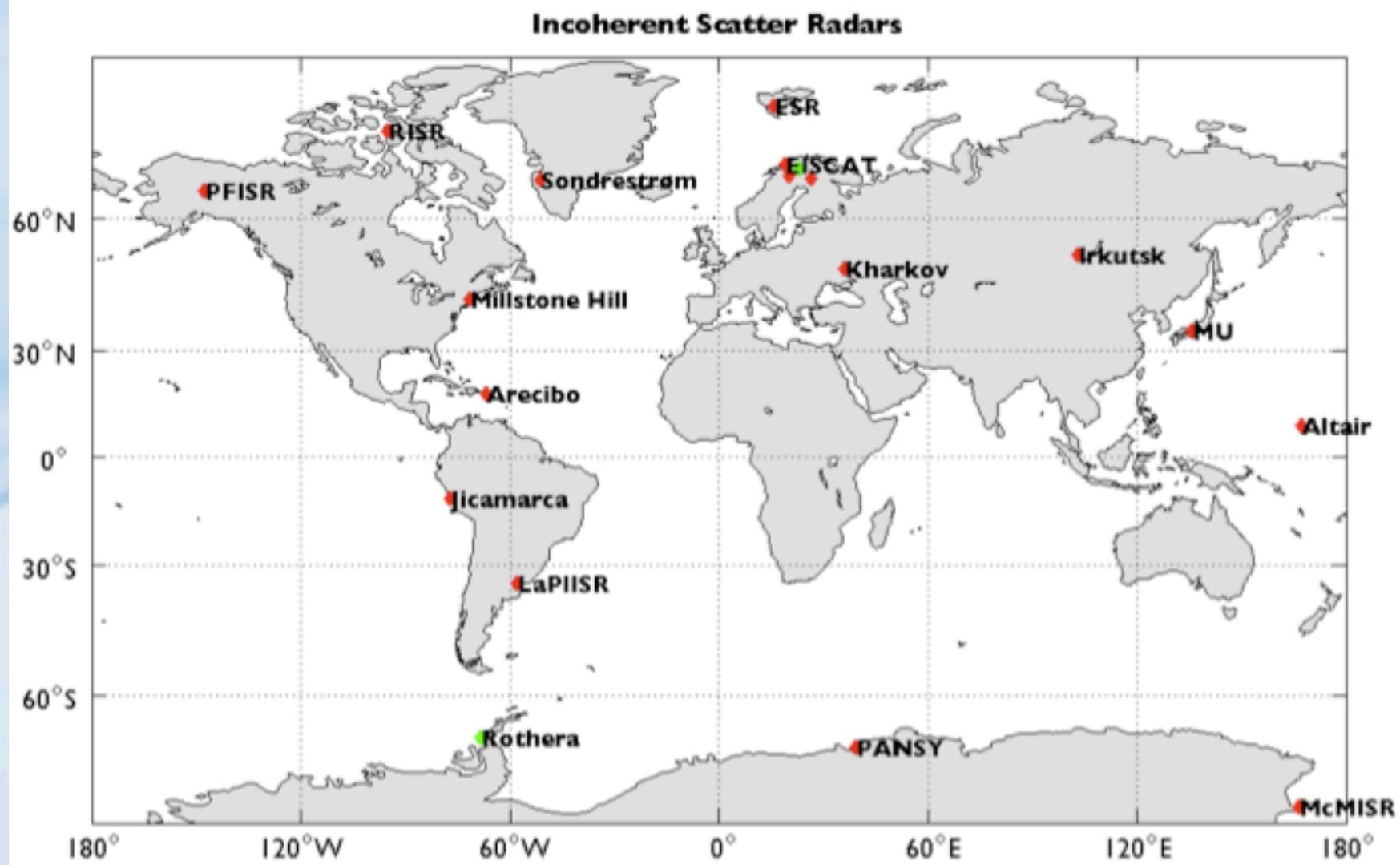


Incoherent Scatter Radars of the World



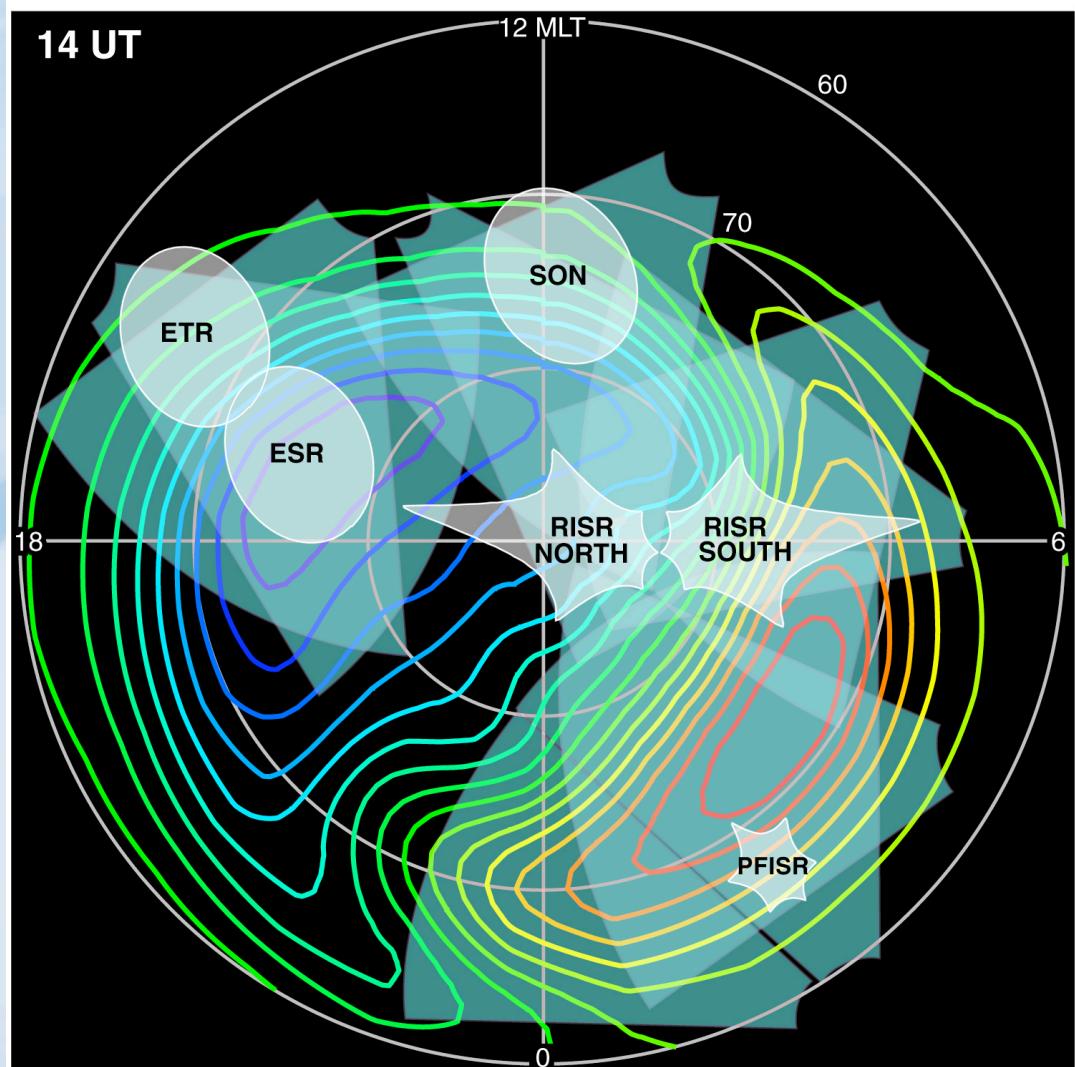


Incoherent Scatter Radars of the World





Map of the north...





High latitude Incoherent Scatter Radars....



PFISR (Poker Flat Incoherent Scatter Radar) (Resolute Bay Incoherent Scatter Radar) AMISRs currently at





Mid-Latitude Incoherent Scatter Radars





Low-Latitude Incoherent Scatter Radars

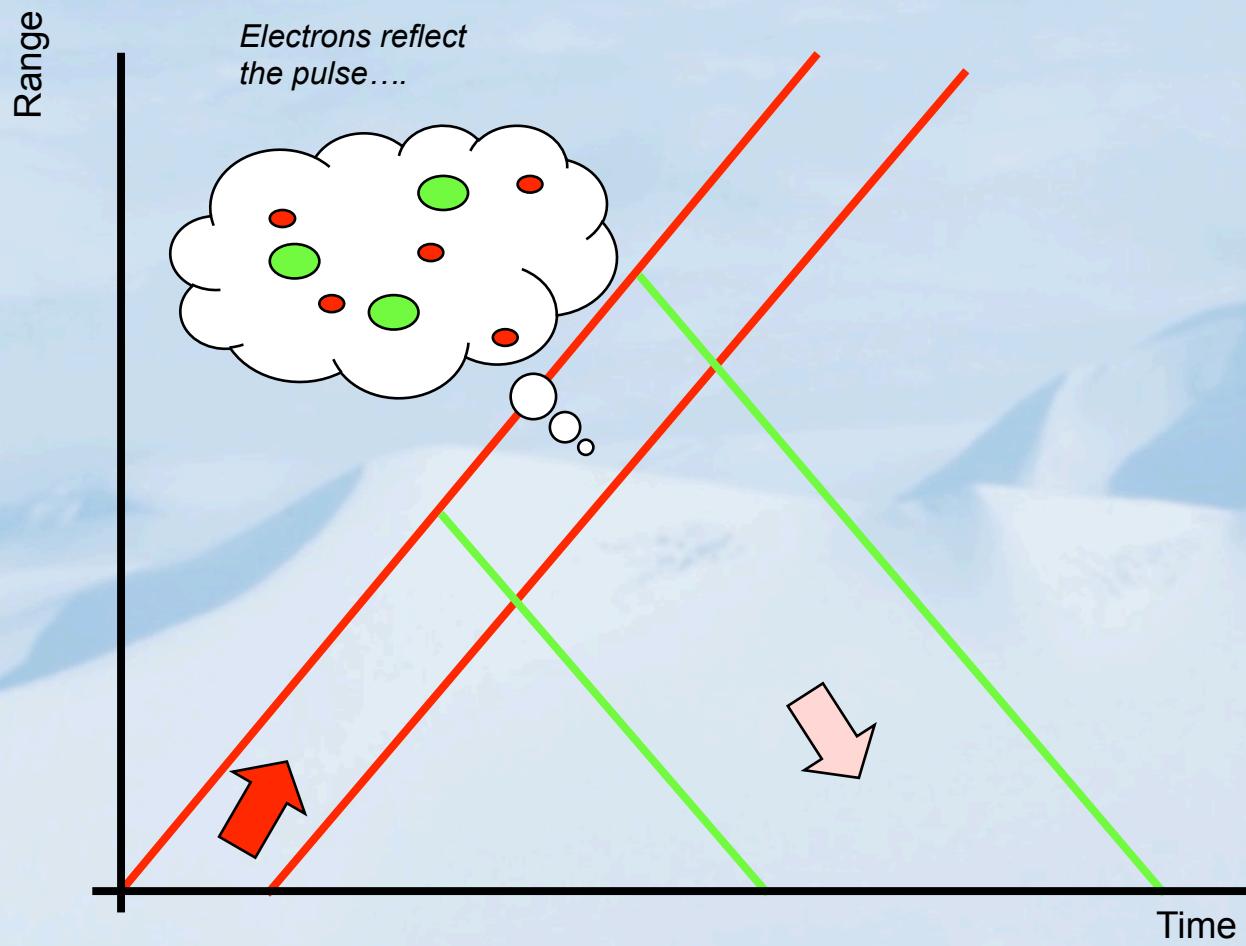




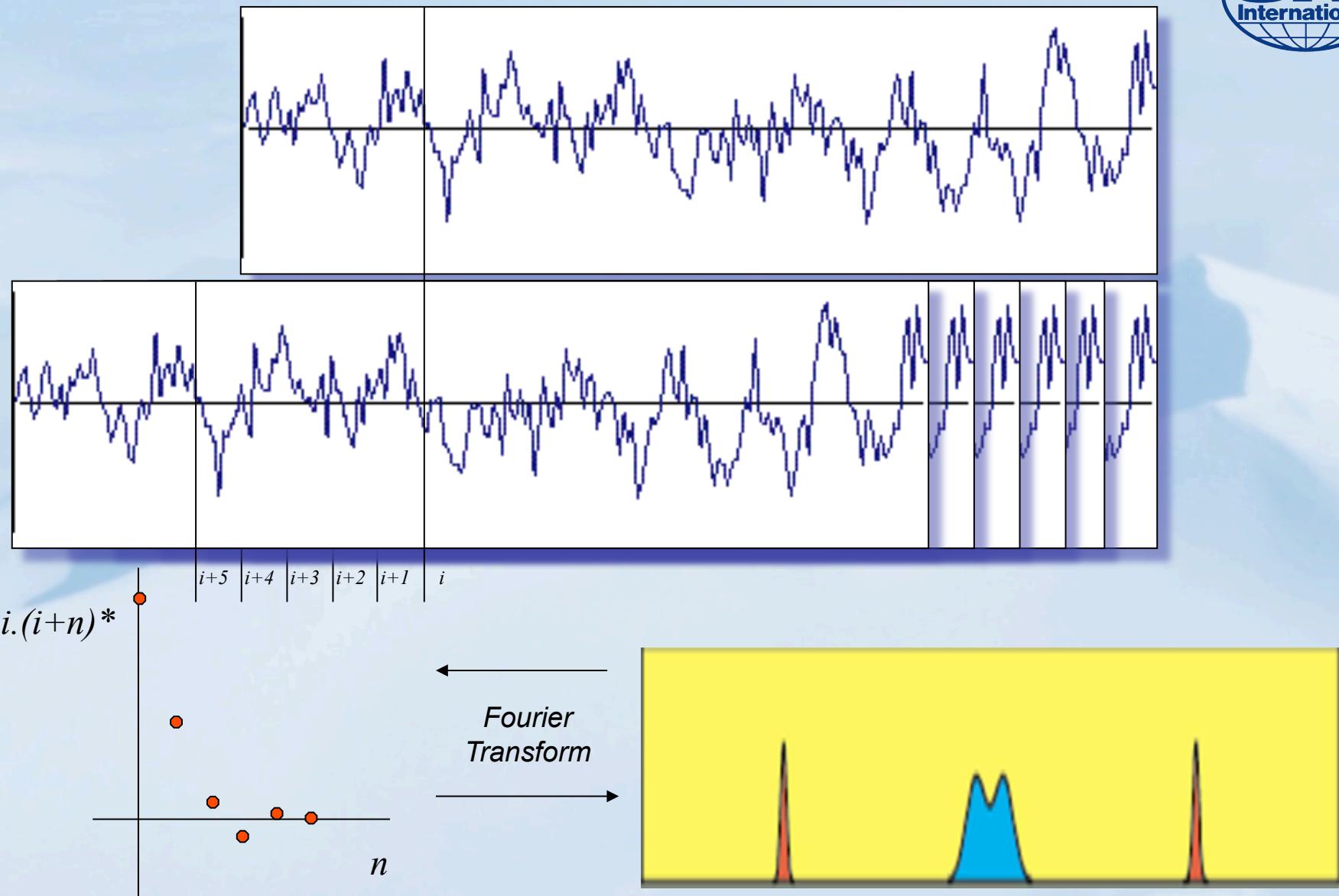
Questions you might have now:

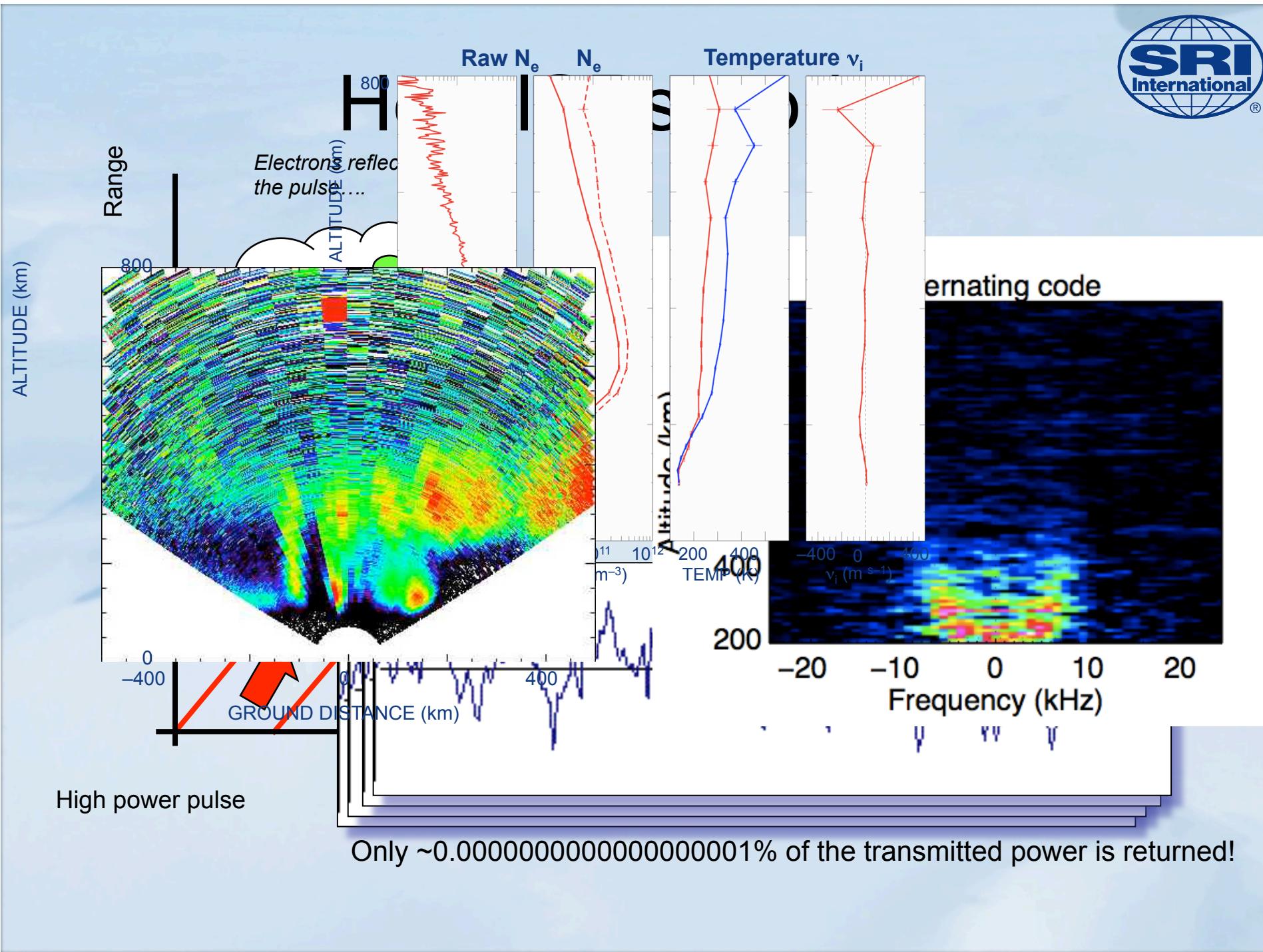
- Why are incoherent scatter radars (ISRs) so big? Is it a status-thing?
- Why is it called *incoherent* scattering?
- What do the ISR returns look like and why?
- What can ISRs measure?
- Can we get through this before lunch?

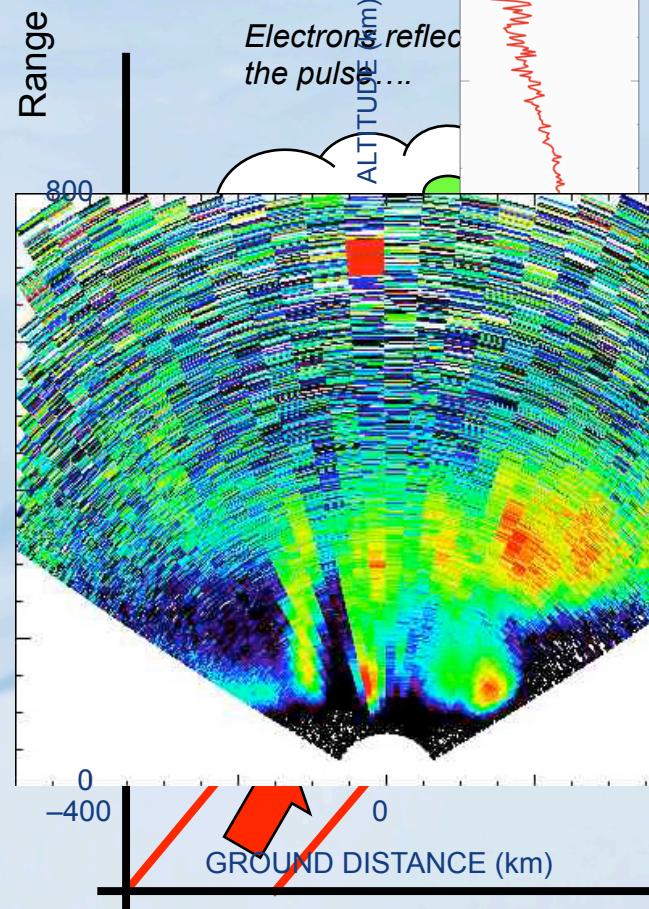
How ISRs work...



Only $\sim 0.0000000000000001\%$ of the transmitted power is returned!

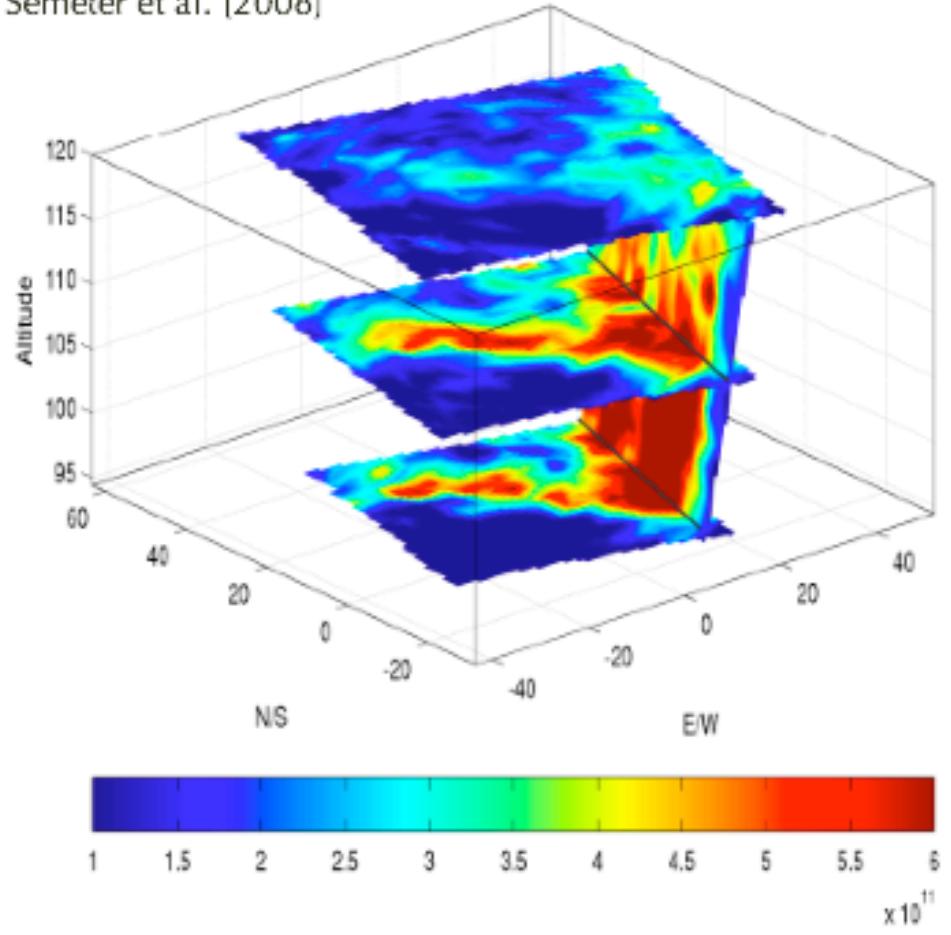






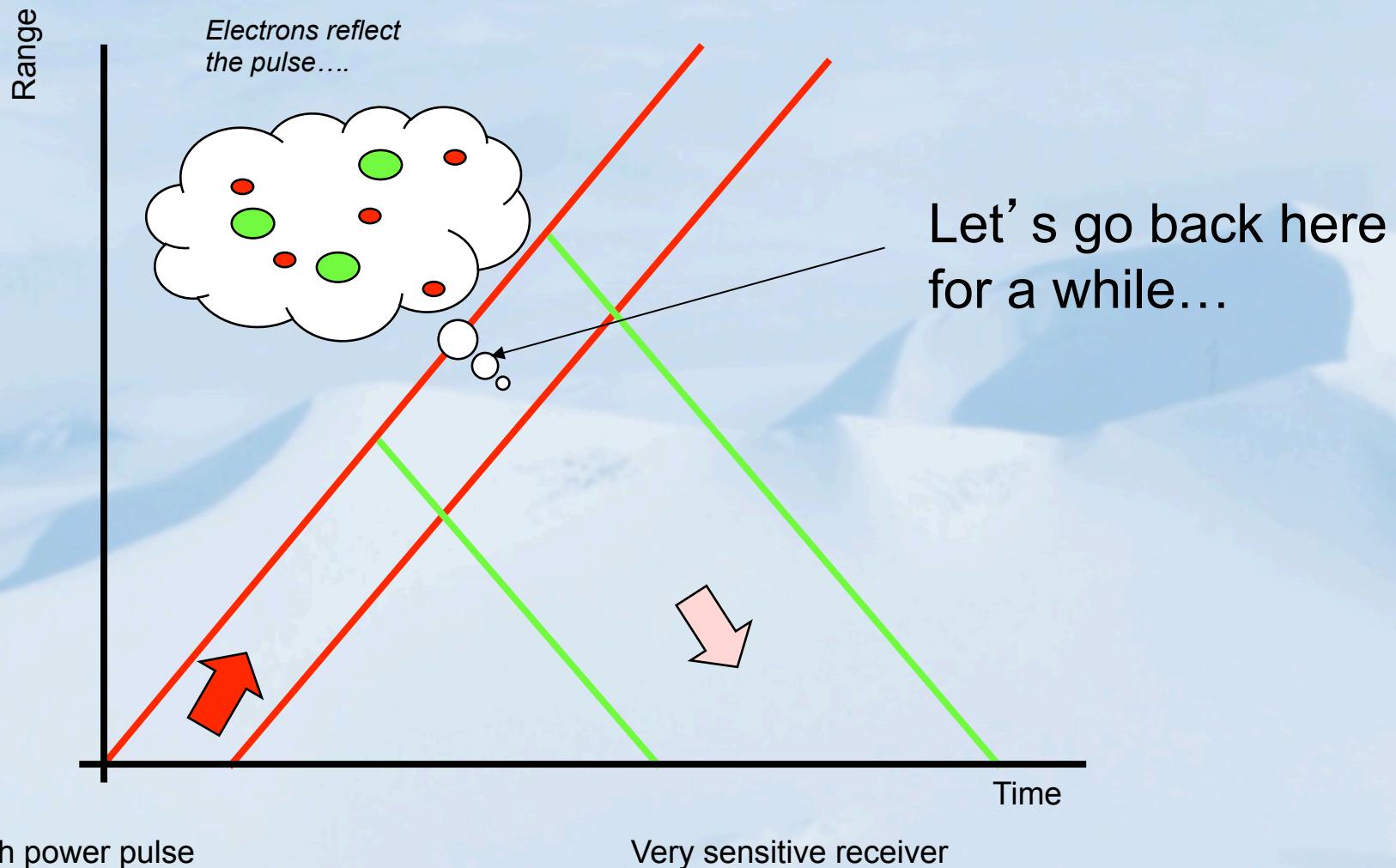
Raw N_e **N_e** **Temperature v_i**

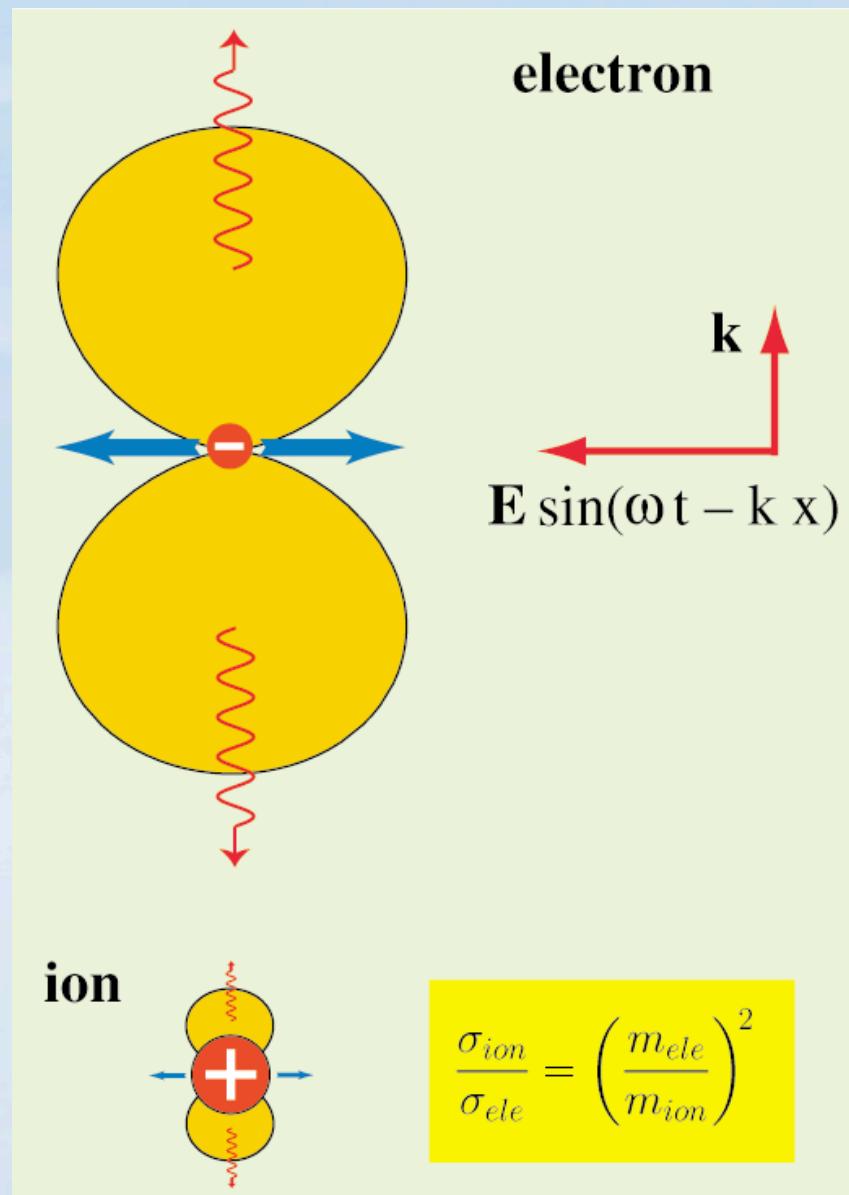
N_e 10-Nov-2007 09:43:51 - 09:44:50
Semeter et al. [2008]



Only $\sim 0.0000000000000001\%$ of the transmitted power is returned!

How ISRs work...

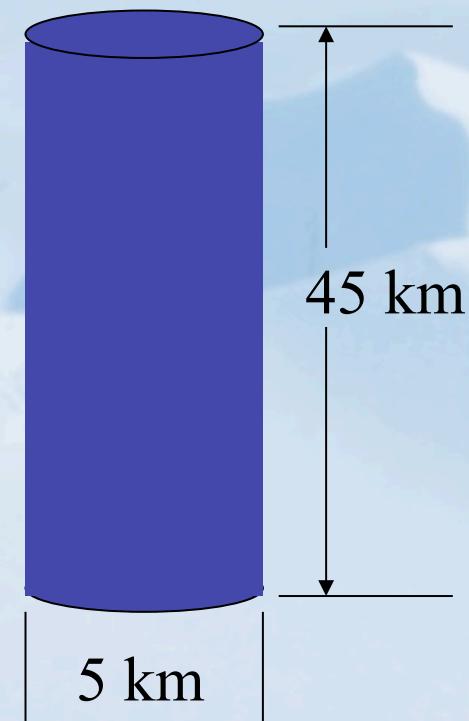




Total Cross-Section Estimate

Consider an antenna with a 1-degree beam measuring the ionospheric plasma at 300 km range and using a 300 microsecond pulse.

If the electron density is 10^{12} m^{-3} , the total number of electrons scattering into a given measurement is $\sim 8.8 \times 10^{23}$. This yields a total cross-section of 88 mm^2 – we need a big radar!

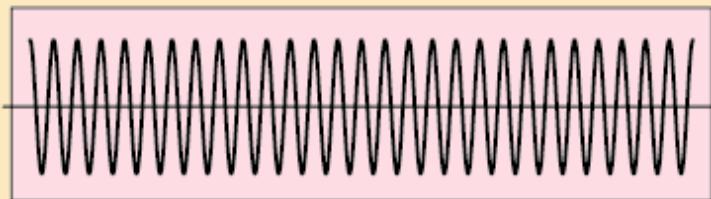


Incoheret scattering - the short story

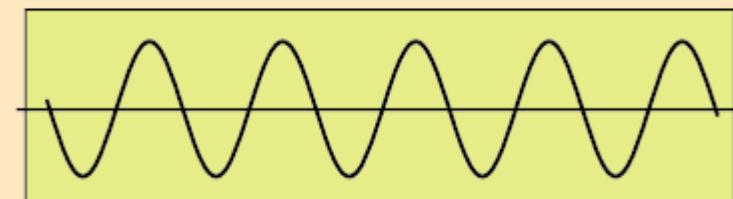


Incoherent scattering - the short story

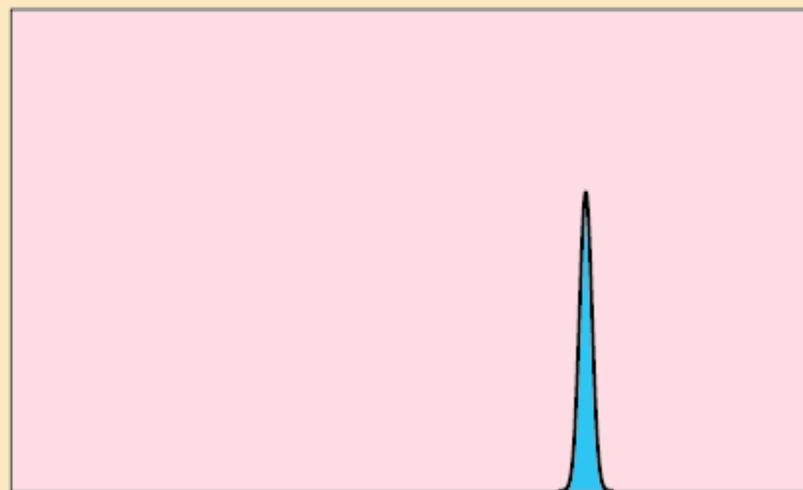




time



time

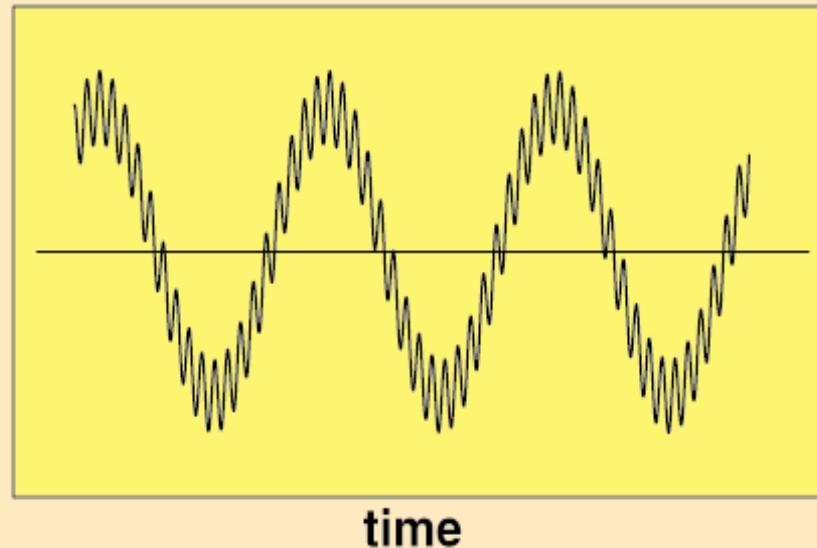


frequency



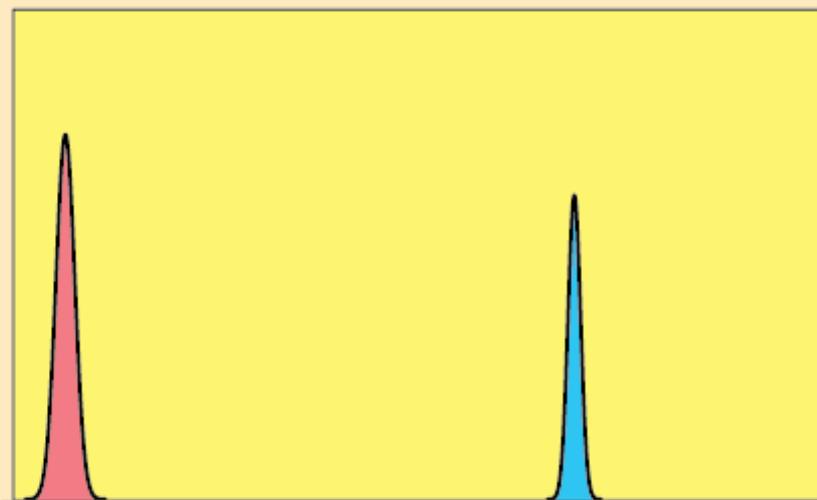
frequency

$$\text{v} \text{ } \text{ } \text{ } \text{ } + \text{ } \text{ } \text{ } \text{ } =$$



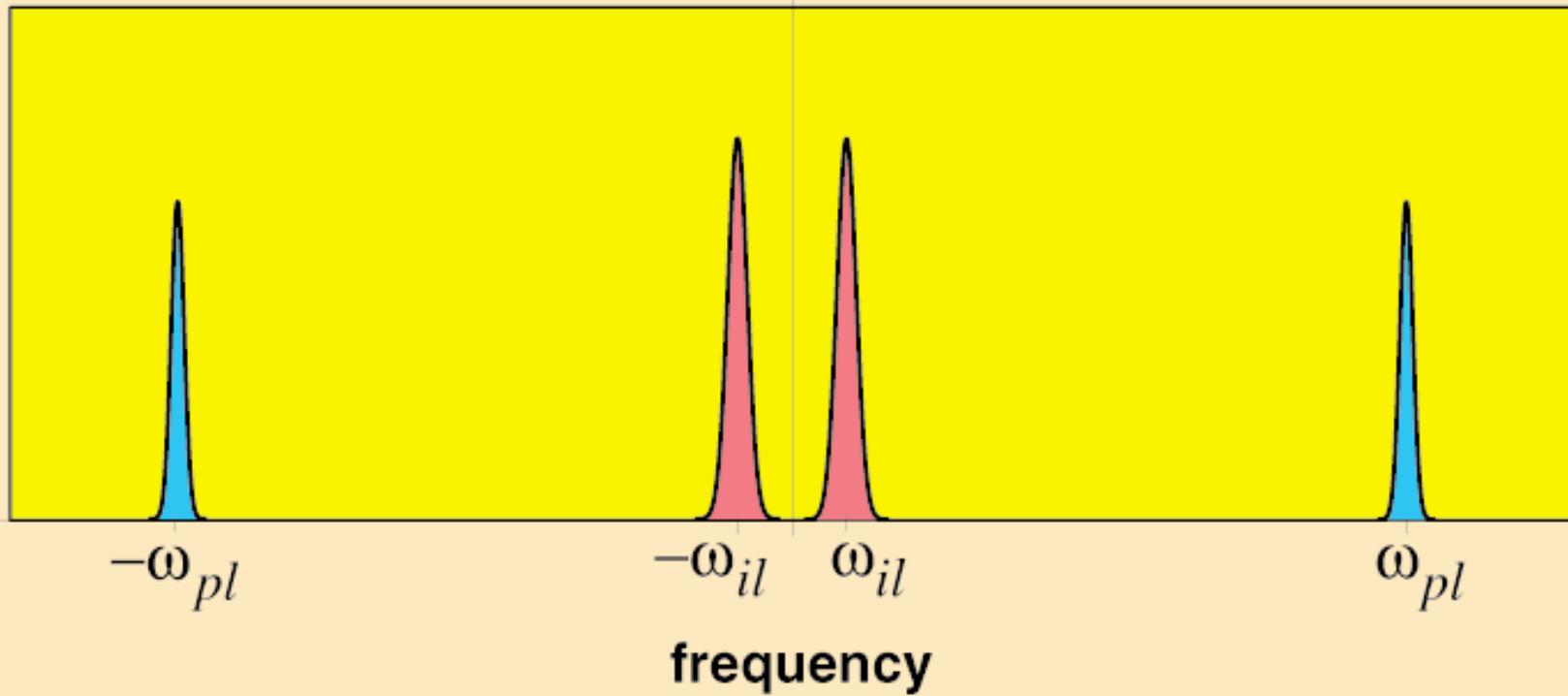
time

$$\text{ } \text{ } \text{ } \text{ } + \text{ } \text{ } \text{ } \text{ } =$$

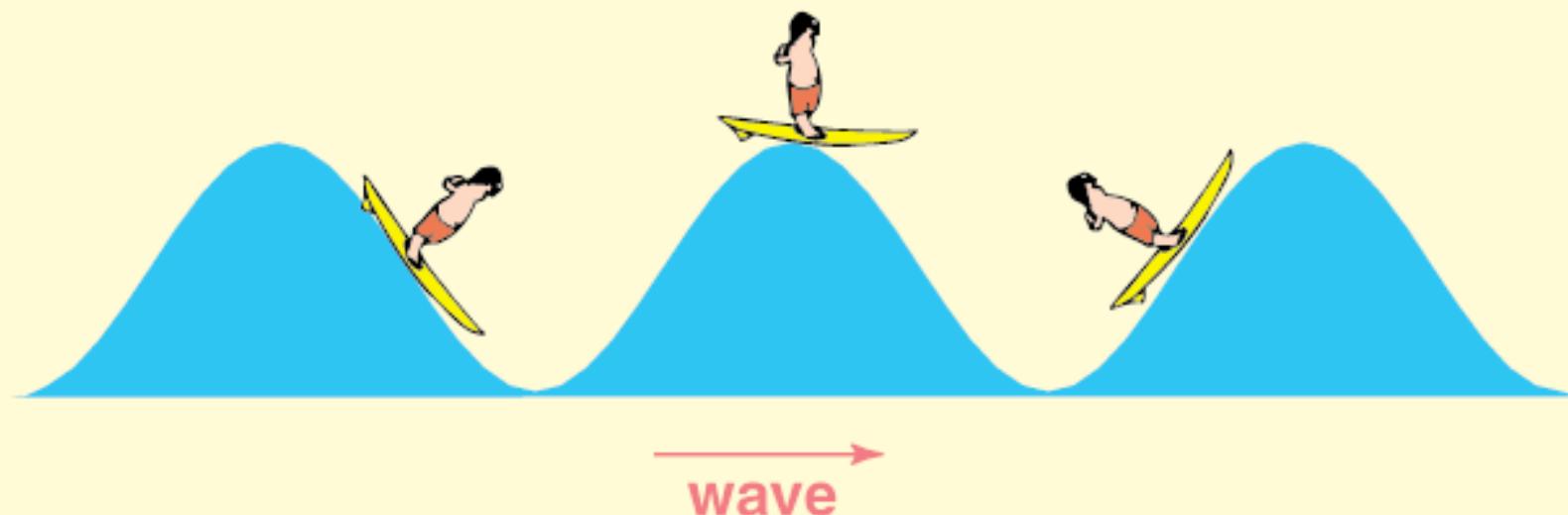


frequency

Plasma Wave Approach (cont'd)



Landau wave-particle interactions



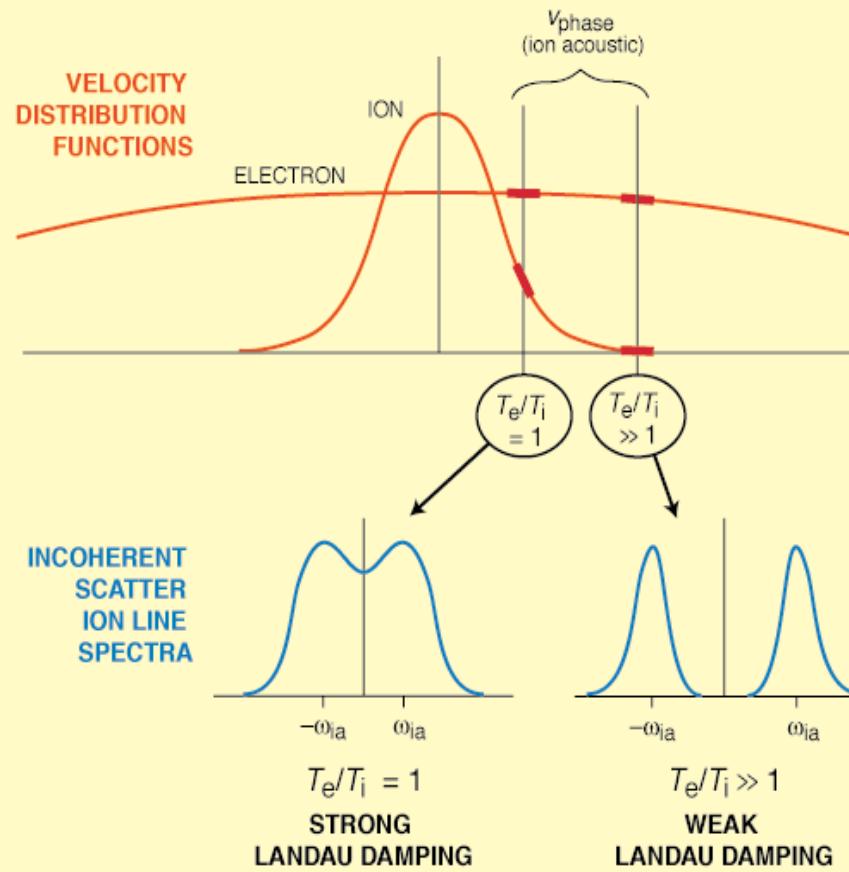
particle
gains
energy

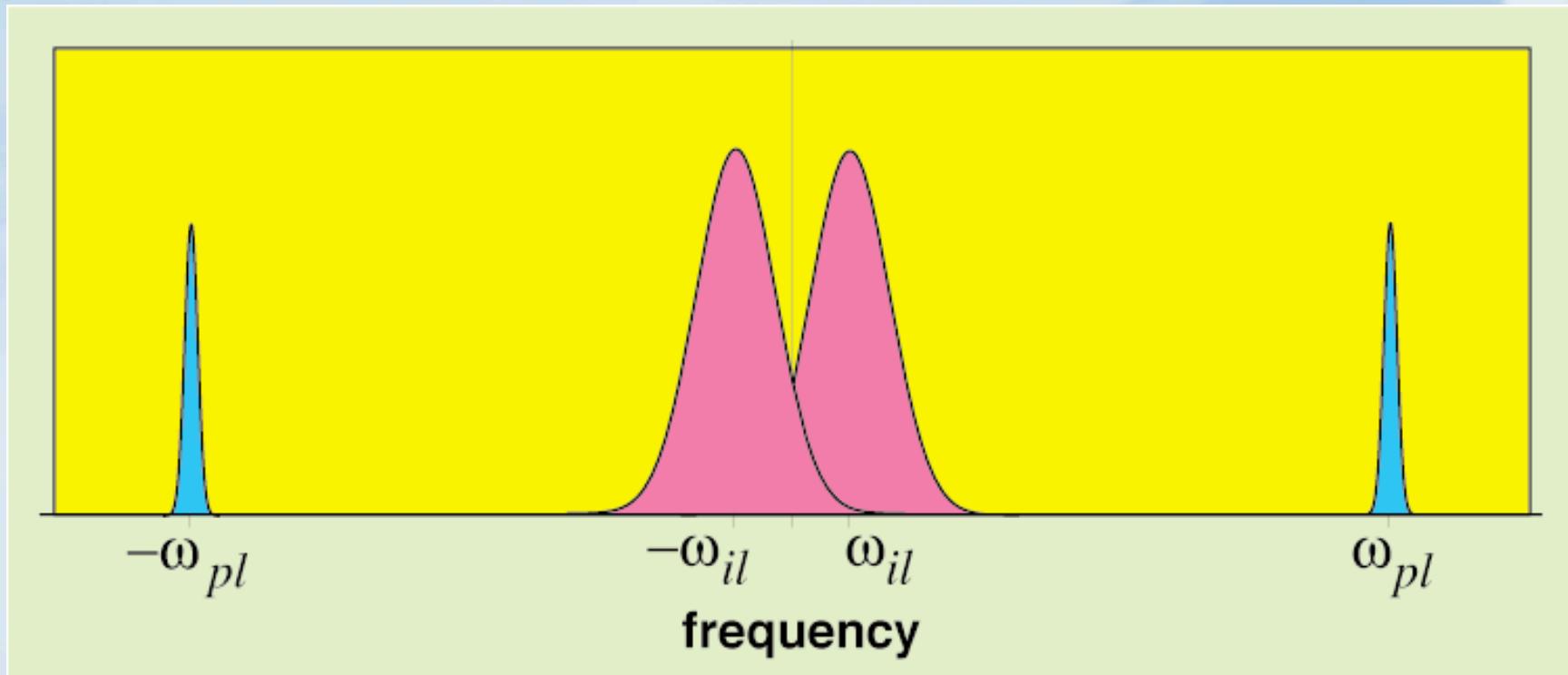
wave
gains
energy

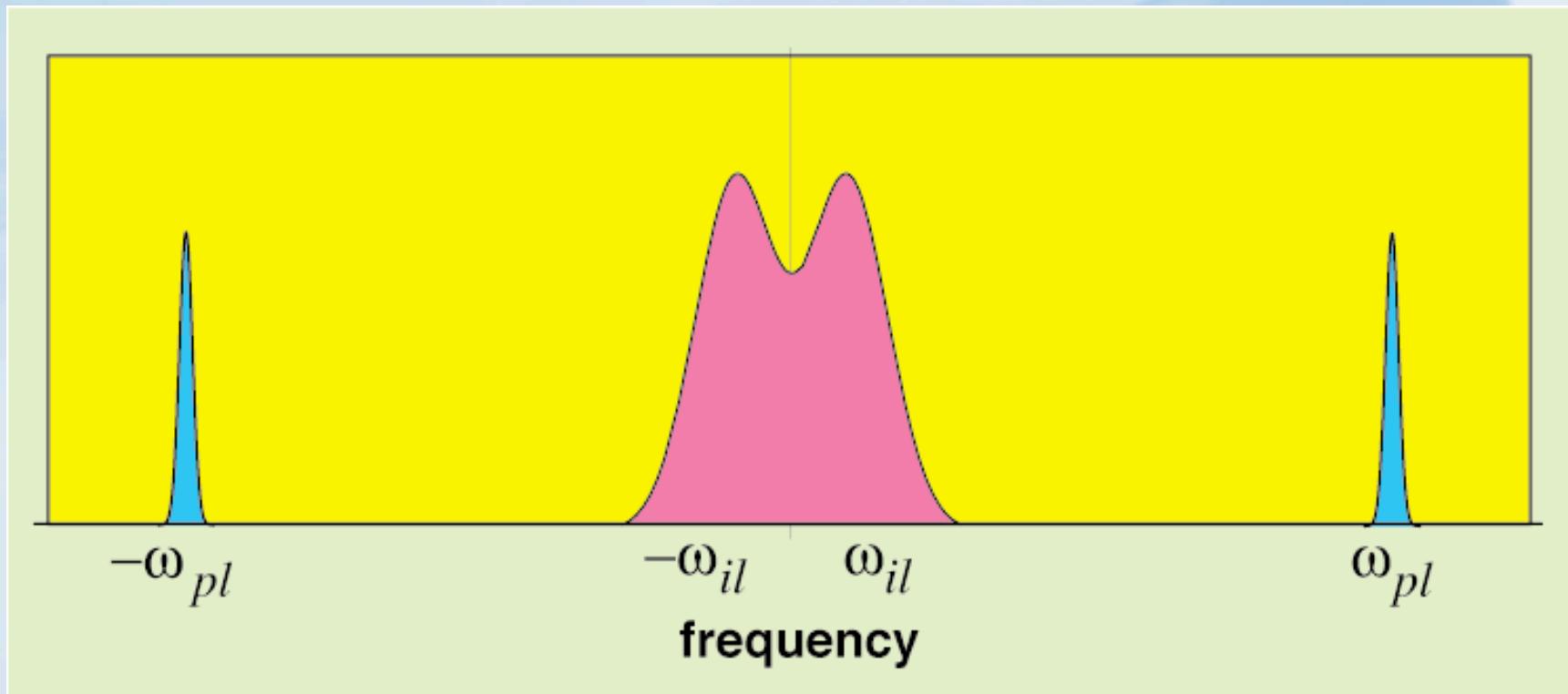
THE EFFECT OF LANDAU DAMPING ON THE INCOHERENT SCATTER ION LINE SPECTRUM

ION-Acoustic
Dispersion
Equation

$$\omega_{ia} = k v_{phase} = k \left(\frac{T_e + 3T_i}{m_i} \right)^{1/2}$$







Incoherent Scattering Spectrum

$$S_e(\mathbf{k}, \omega) = N_e \left| 1 - \frac{\chi_e(\mathbf{k}, \omega)}{\epsilon(\mathbf{k}, \omega)} \right|^2 \int d\mathbf{v} f_e(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v}) + N_i \left| \frac{\chi_e(\mathbf{k}, \omega)}{\epsilon(\mathbf{k}, \omega)} \right|^2 \int d\mathbf{v} f_i(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v})$$

Plasma line Ion line

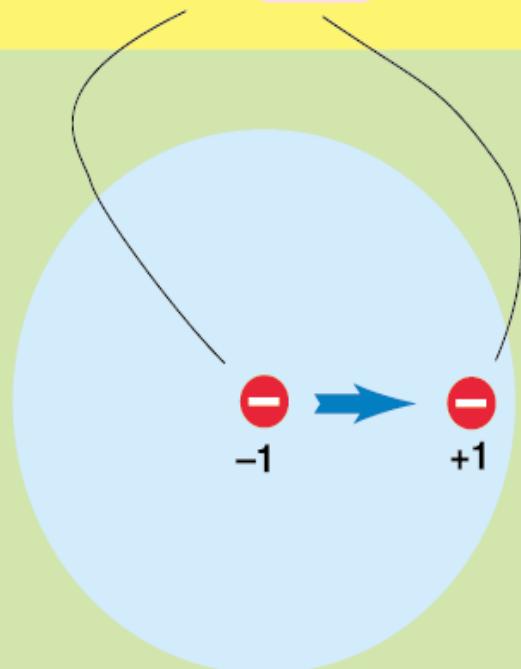
electric susceptibility $\chi_{e,i}(\mathbf{k}, \omega)$

dielectric constant function $\epsilon(\mathbf{k}, \omega)$

velocity distribution function $f_{e,i}(\mathbf{v})$

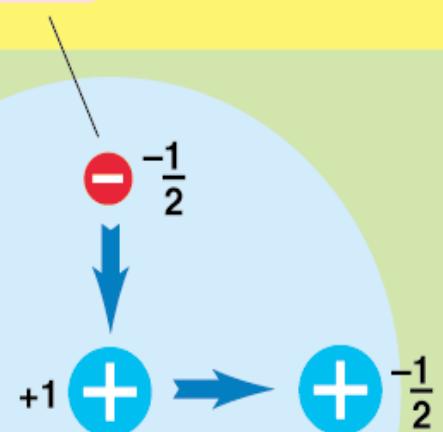
Plasma Line $S_{PL}(\mathbf{k}, \omega)$

$$S_e(\mathbf{k}, \omega) = N_e \left| 1 - \frac{\chi_e(\mathbf{k}, \omega)}{\epsilon(\mathbf{k}, \omega)} \right|^2 \int d\mathbf{v} f_e(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v}) + N_i \left| \frac{\chi_e(\mathbf{k}, \omega)}{\epsilon(\mathbf{k}, \omega)} \right|^2 \int d\mathbf{v} f_i(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v})$$



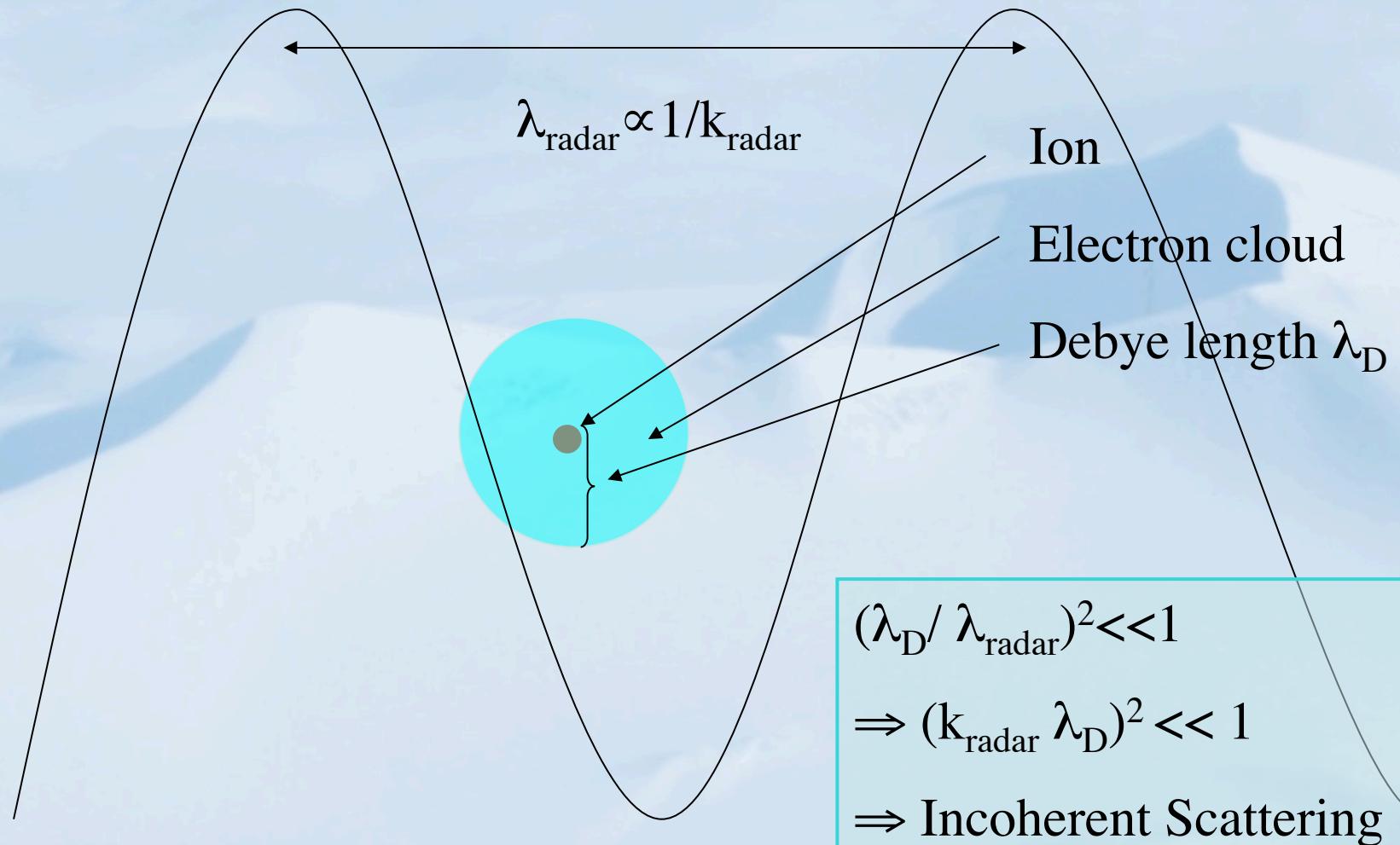
electron with cloud

Ion Line $S_{IL}(\mathbf{k}, \omega)$



ion with cloud

Debye length dependence





Plasma Line $S_{PL}(\mathbf{k}, \omega)$

Ion Line $S_{IL}(\mathbf{k}, \omega)$

$$S_e(\mathbf{k}, \omega) = N_e \left| 1 - \frac{\chi_e(\mathbf{k}, \omega)}{\epsilon(\mathbf{k}, \omega)} \right|^2 \int d\mathbf{v} f_e(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v}) + N_i \left| \frac{\chi_e(\mathbf{k}, \omega)}{\epsilon(\mathbf{k}, \omega)} \right|^2 \int d\mathbf{v} f_i(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v})$$

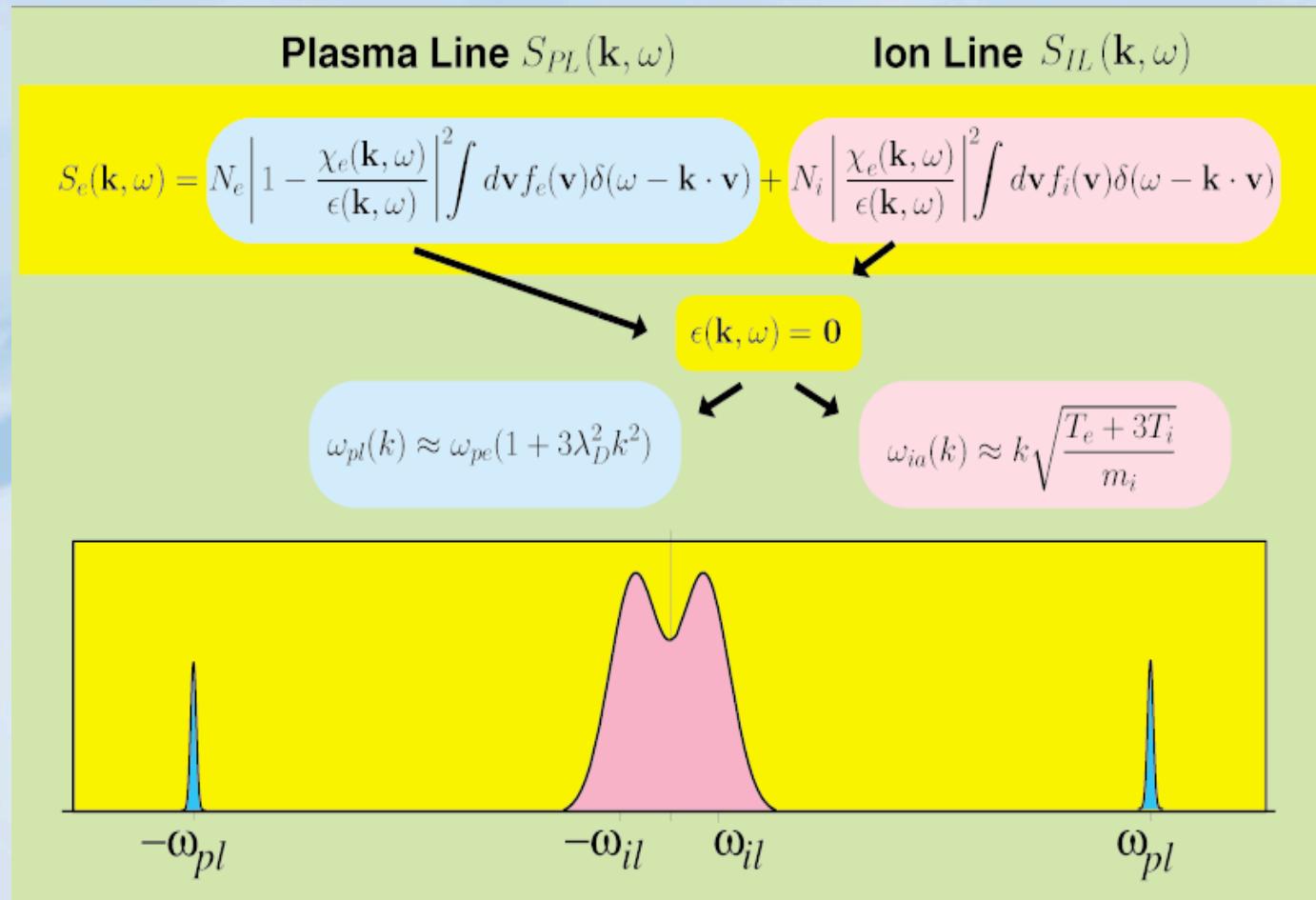
Plasma Line $S_{PL}(\mathbf{k}, \omega)$
Ion Line $S_{IL}(\mathbf{k}, \omega)$

$$S_e(\mathbf{k}, \omega) = N_e \left| 1 - \frac{\chi_e(\mathbf{k}, \omega)}{\epsilon(\mathbf{k}, \omega)} \right|^2 \int d\mathbf{v} f_e(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v}) + N_i \left| \frac{\chi_e(\mathbf{k}, \omega)}{\epsilon(\mathbf{k}, \omega)} \right|^2 \int d\mathbf{v} f_i(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v})$$

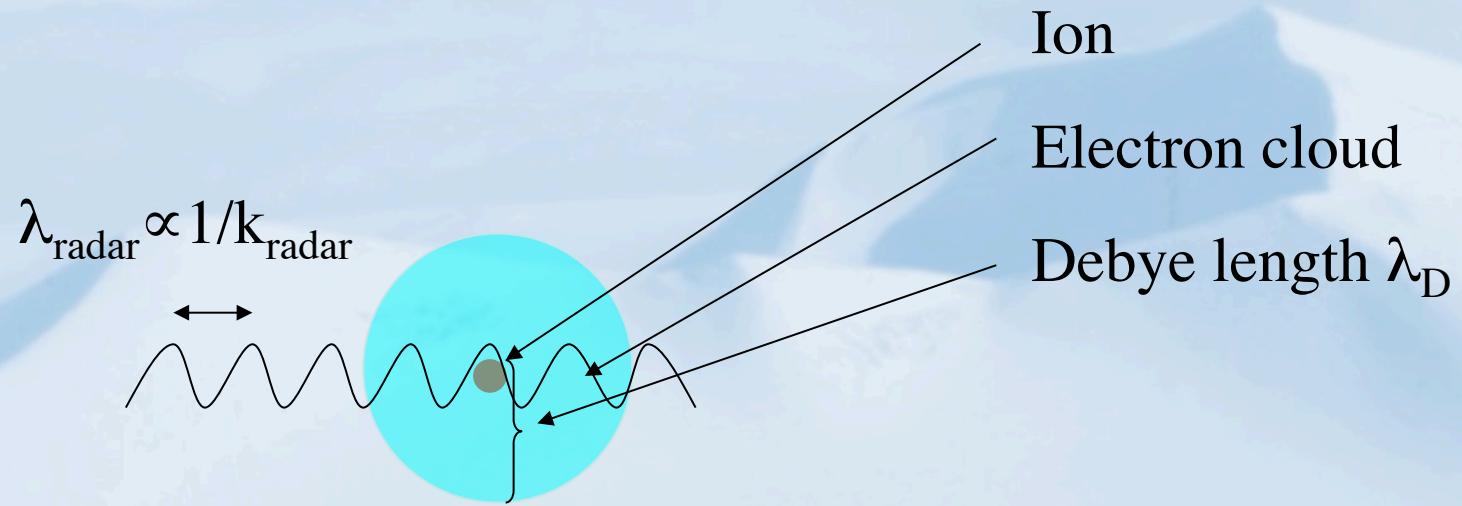
$$\epsilon(\mathbf{k}, \omega) = 0$$

$$\omega_{pl}(k) \approx \omega_{pe} (1 + 3\lambda_D^2 k^2)$$

$$\omega_{ia}(k) \approx k \sqrt{\frac{T_e + 3T_i}{m_i}}$$



Debye length dependence



$$(\lambda_D / \lambda_{\text{radar}})^2 > 1$$

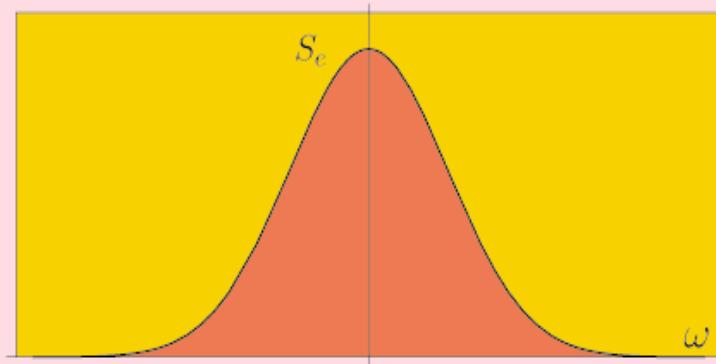
$$\Rightarrow (k_{\text{radar}} \lambda_D)^2 > 1$$

\Rightarrow No collective interactions

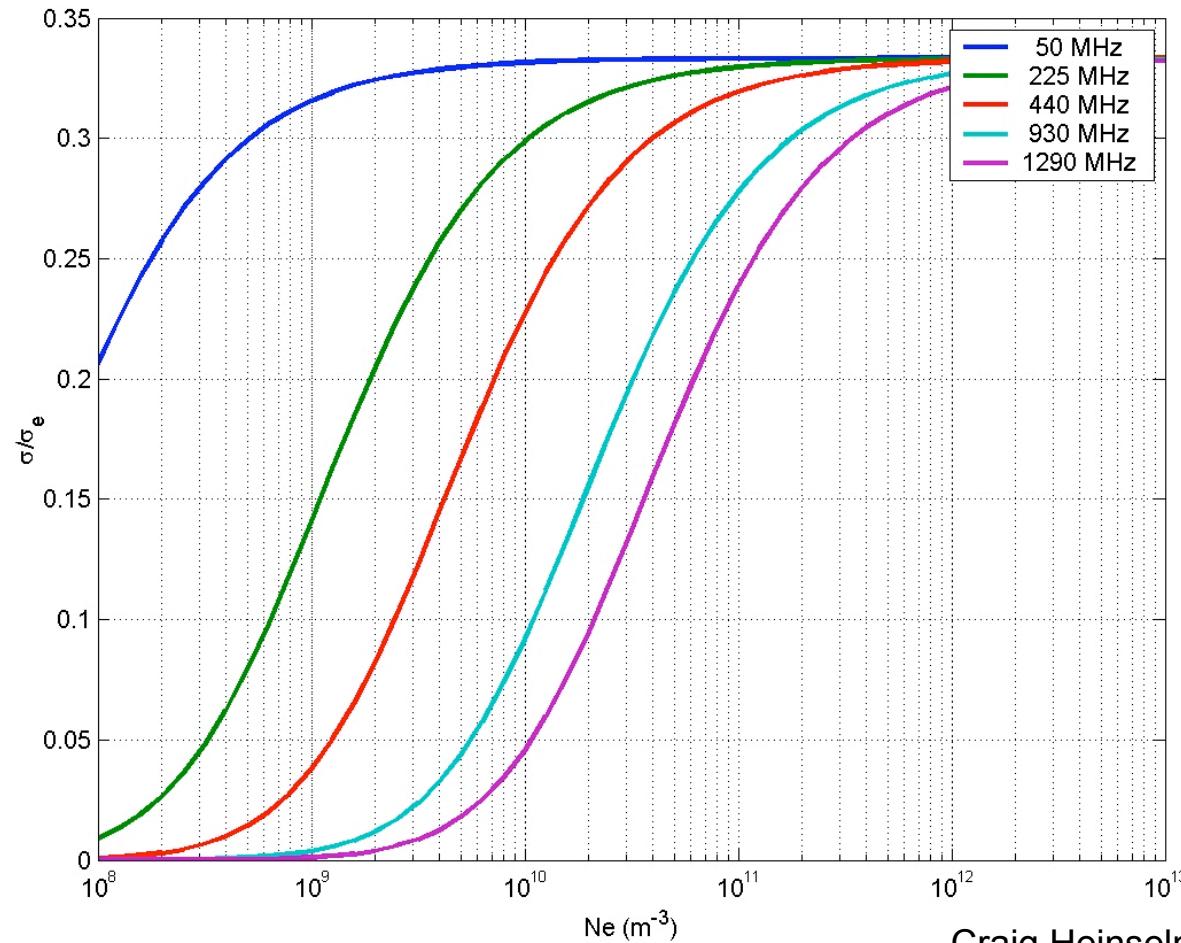
no collective interactions

$$S_e(\mathbf{k}, \omega) = N_e \left| 1 - \frac{\chi_e(\mathbf{k}, \omega)}{\epsilon(\mathbf{k}, \omega)} \right|^2 \int d\mathbf{v} f_e(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v}) + N_i \left| \frac{\chi_e(\mathbf{k}, \omega)}{\epsilon(\mathbf{k}, \omega)} \right|^2 \int d\mathbf{v} f_i(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v})$$

$$S_e(\mathbf{k}, \omega) = N_e \int d\mathbf{v} f_e(\mathbf{v}) \delta(\omega - \mathbf{k} \cdot \mathbf{v})$$



Debye Length Dependencies

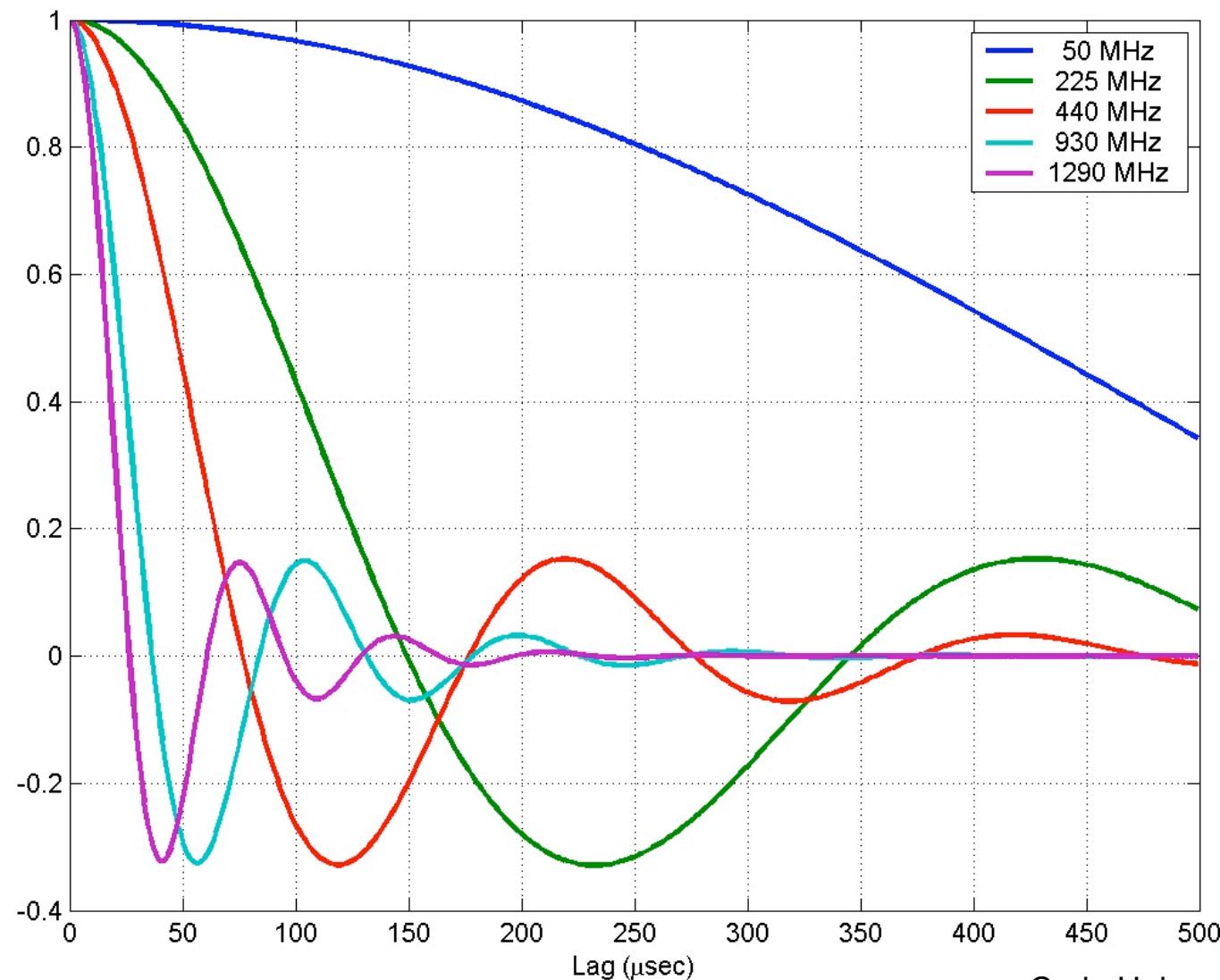


Parameters

T_i: 1000 K

T_e: 2000 K

Craig Heinselman



Parameters

$N_e: 10^{12} \text{ m}^{-3}$

$T_i: 1000 \text{ K}$

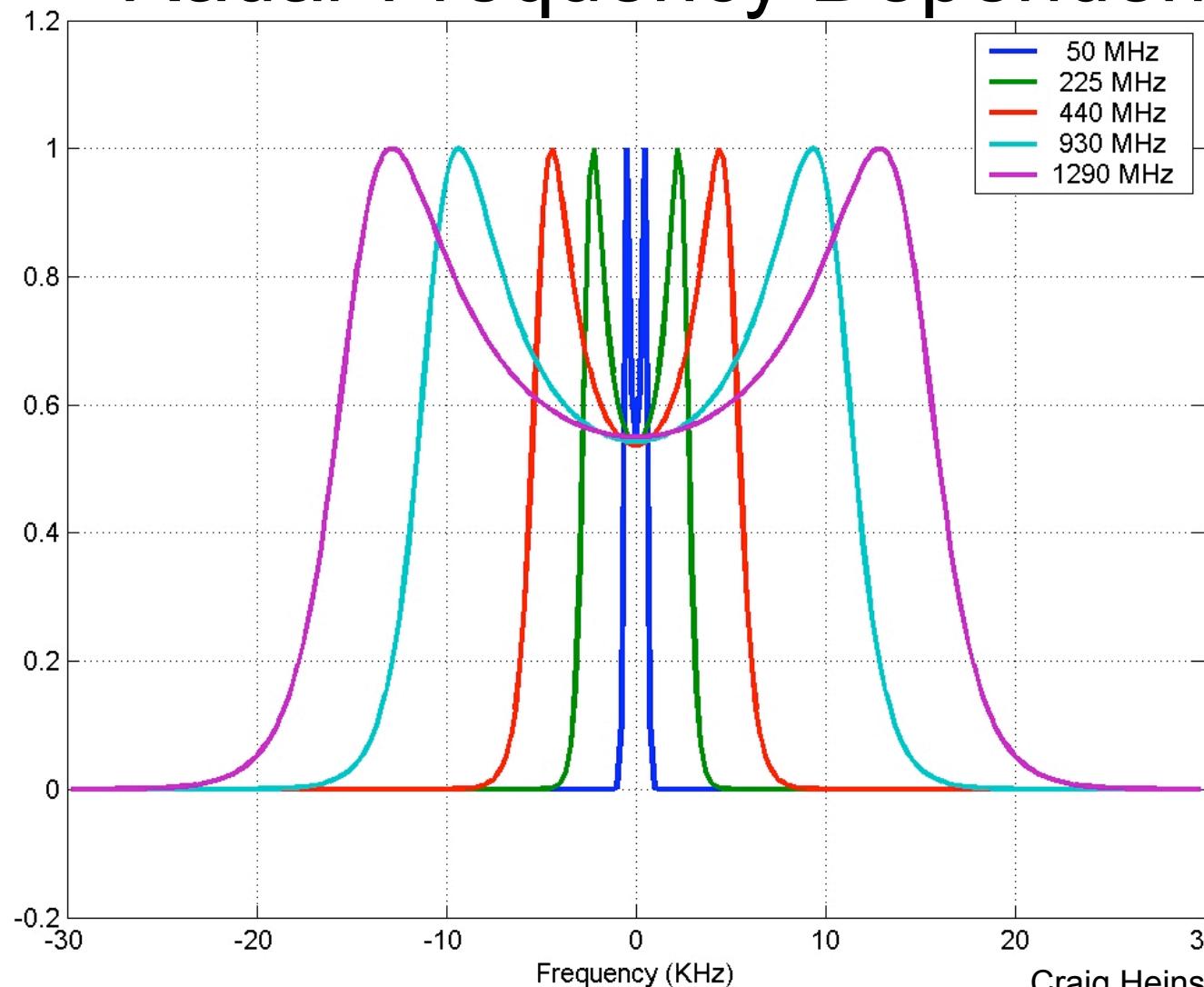
$T_e: 2000 \text{ K}$

Comp: 100% O^+

$v_{in}: 10^{-6} \text{ KHz}$

Craig Heinselman

Radar Frequency Dependencies



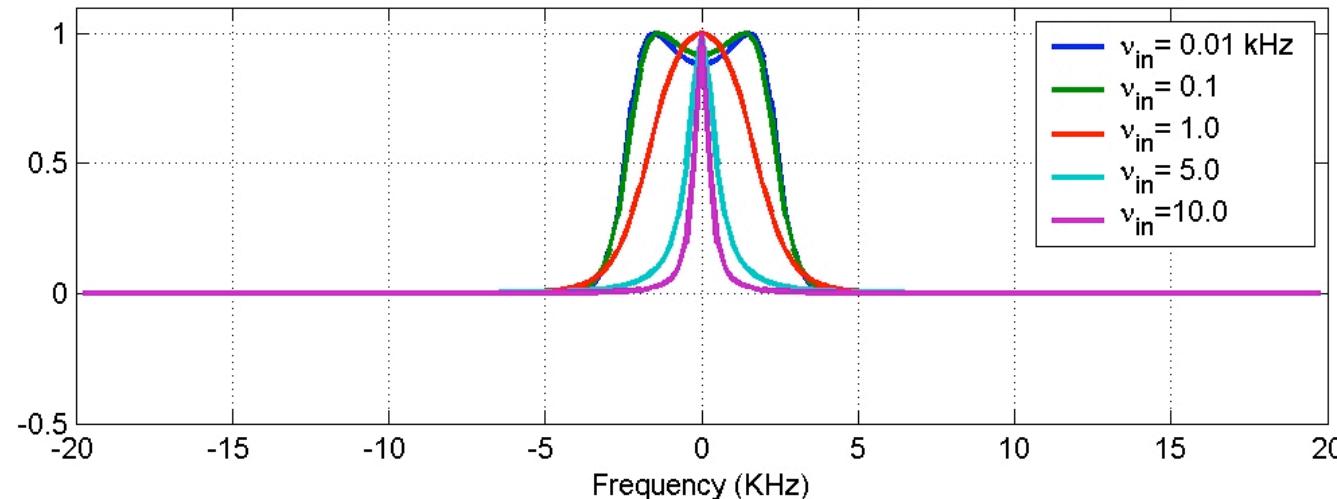
Parameters

$N_e: 10^{12} \text{ m}^{-3}$
 $T_i: 1000 \text{ K}$
 $T_e: 2000 \text{ K}$
 Comp: 100% O⁺
 $v_{in}: 10^{-6} \text{ KHz}$

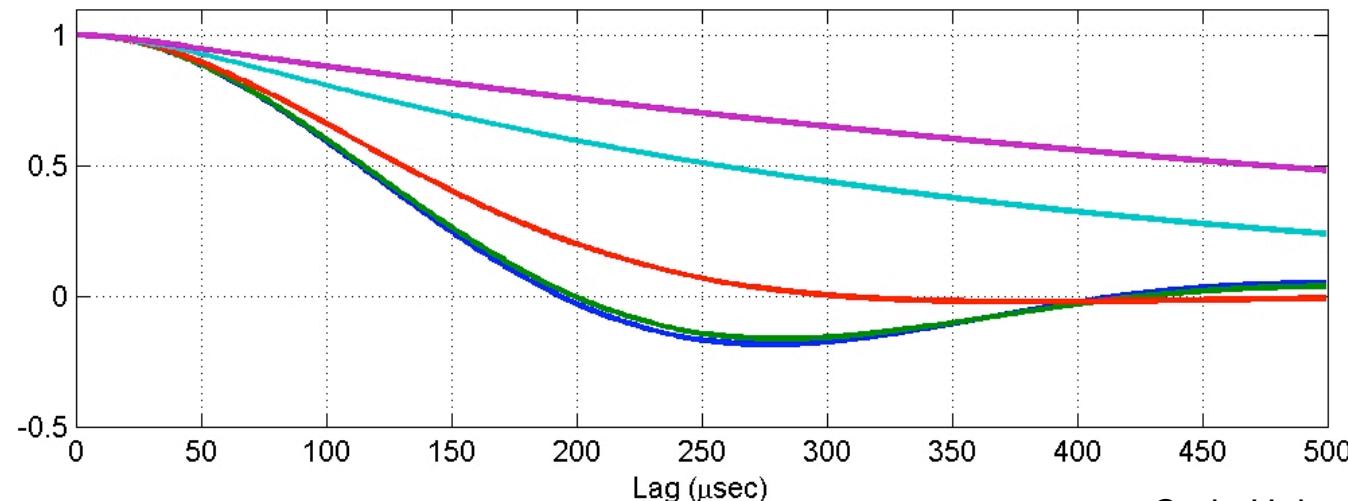


With the frequency of the radar chosen (which is a one time thing!), how do the spectra depend on geophysical parameters?

Ion-Neutral Collision Frequency

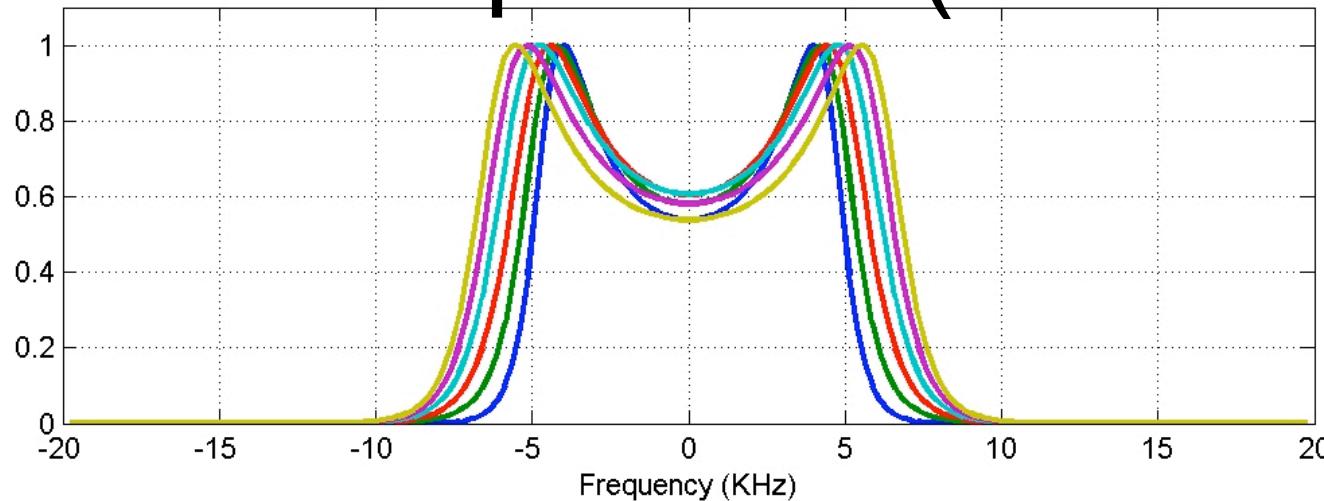


Parameters
 Freq: 449 MHz
 Ne: 10^{12} m^{-3}
 Ti: 500 K
 Te: 500 K
 Comp: 100% NO⁺



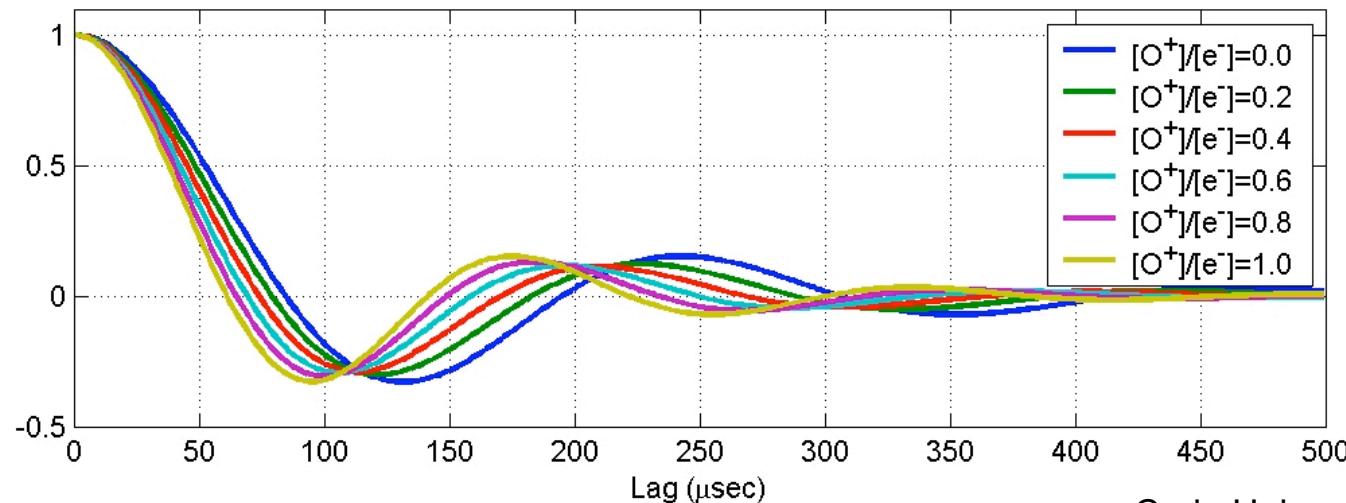
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Ion Composition (O^+ vs. NO^+)



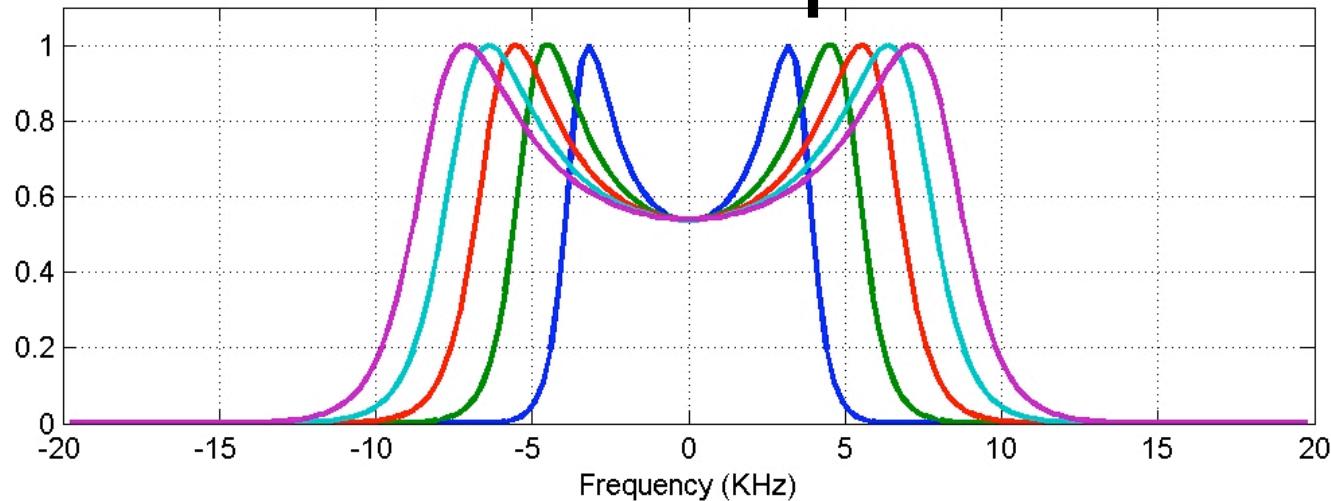
Parameters

- Freq: 449 MHz
- N_e : 10^{12} m^{-3}
- T_i : 1500 K
- T_e : 3000 K
- v_{in} : 10^{-6} KHz

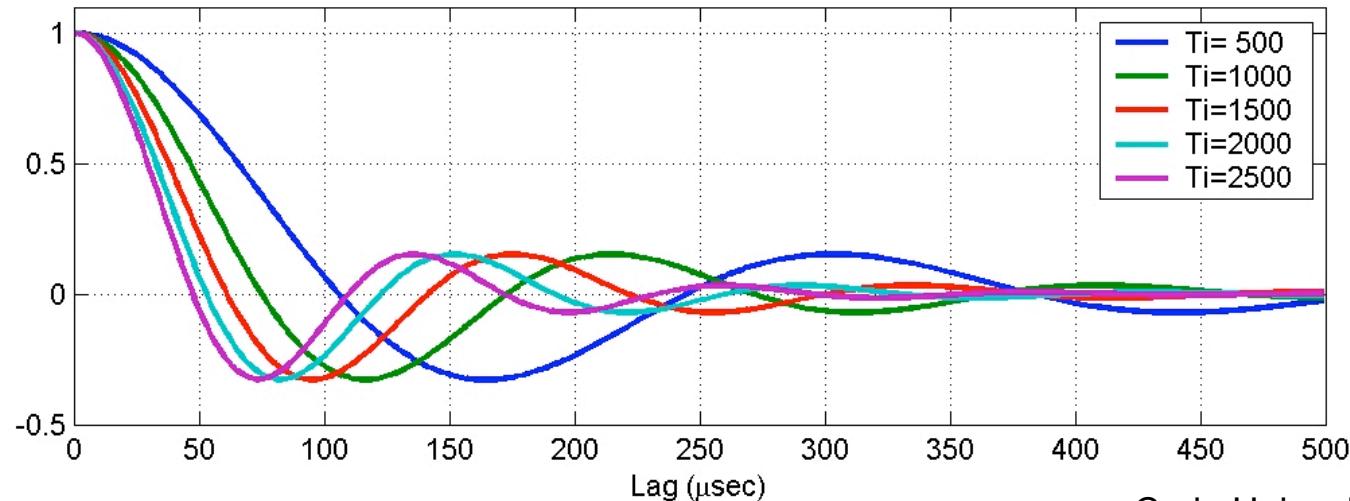


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Ion Temperature

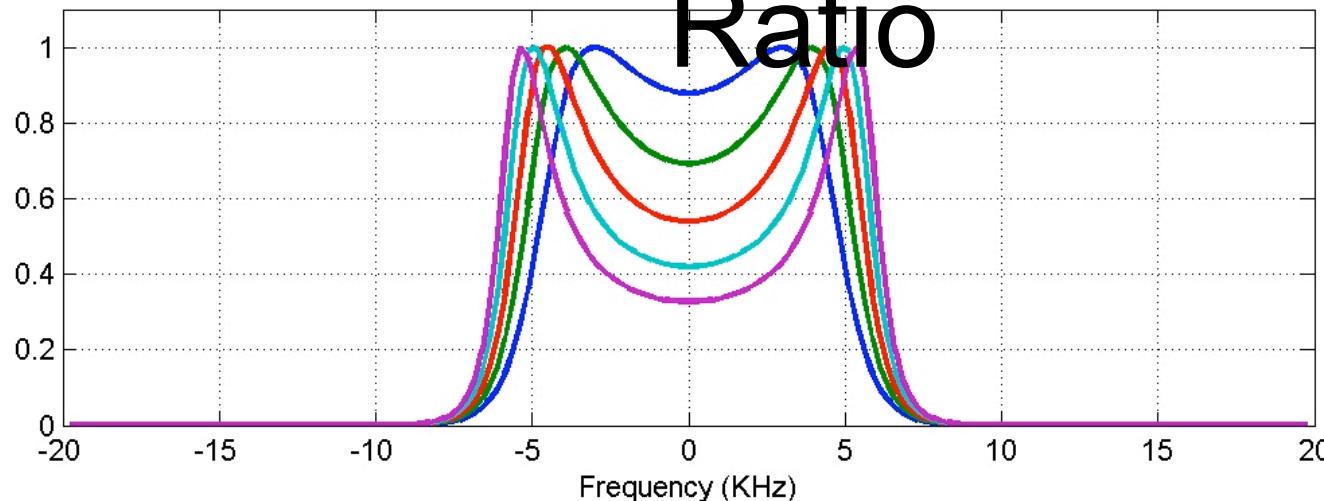


Parameters
 Freq: 449 MHz
 $N_e: 10^{12} \text{ m}^{-3}$
 $T_e: 2 * T_i$
 Comp: 100% O^+
 $v_{in}: 10^{-6} \text{ KHz}$



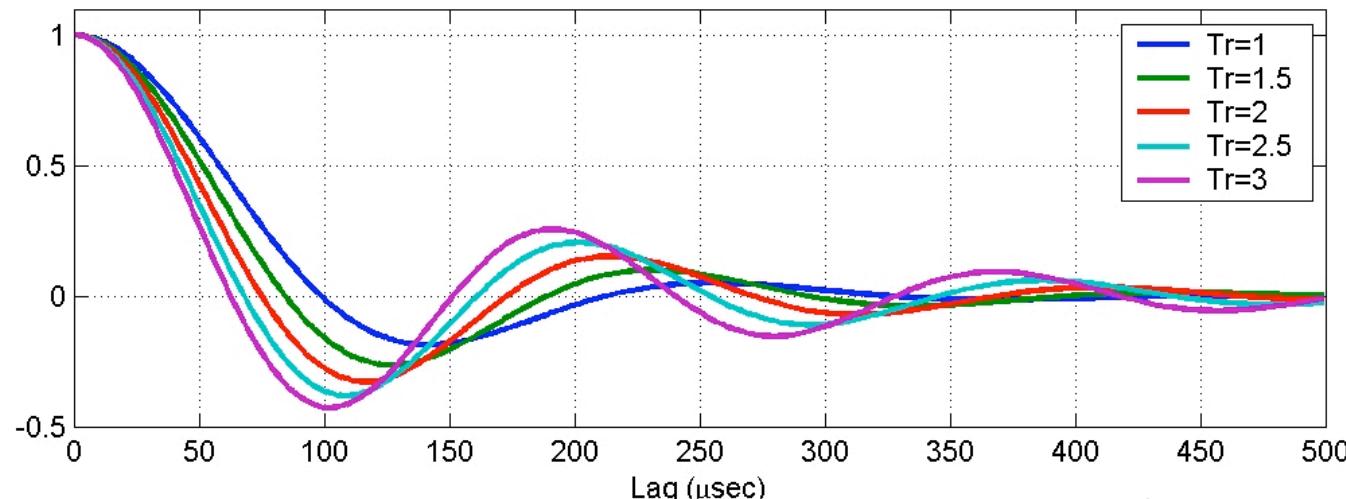
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Electron/Ion Temperature Ratio



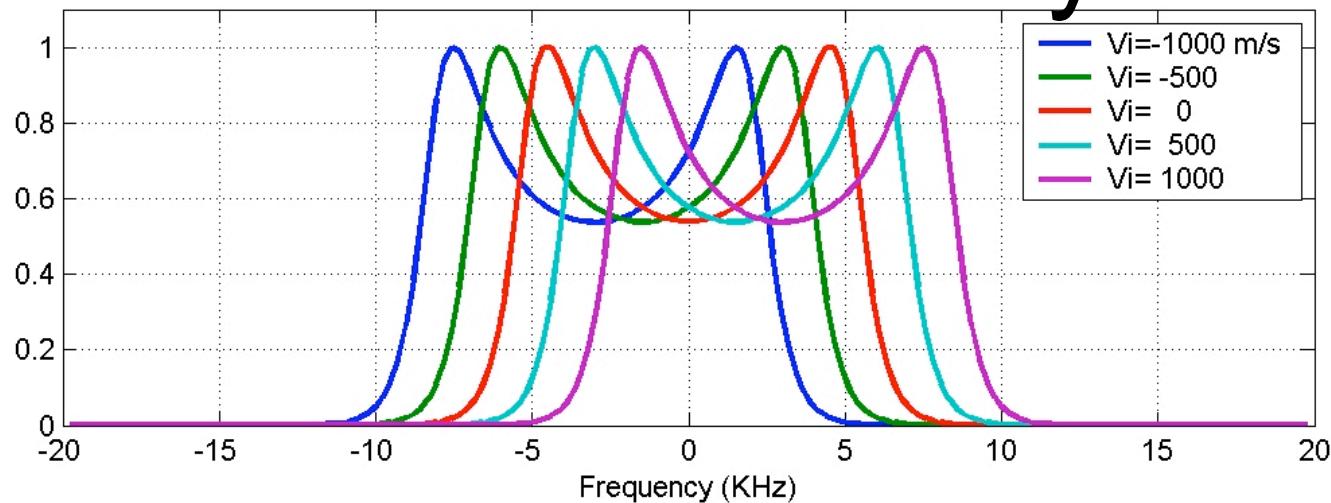
Parameters

- Freq: 449 MHz
- N_e : 10^{12} m^{-3}
- T_i : 1000 K
- Comp: 100% O⁺
- v_{in} : 10^{-6} KHz



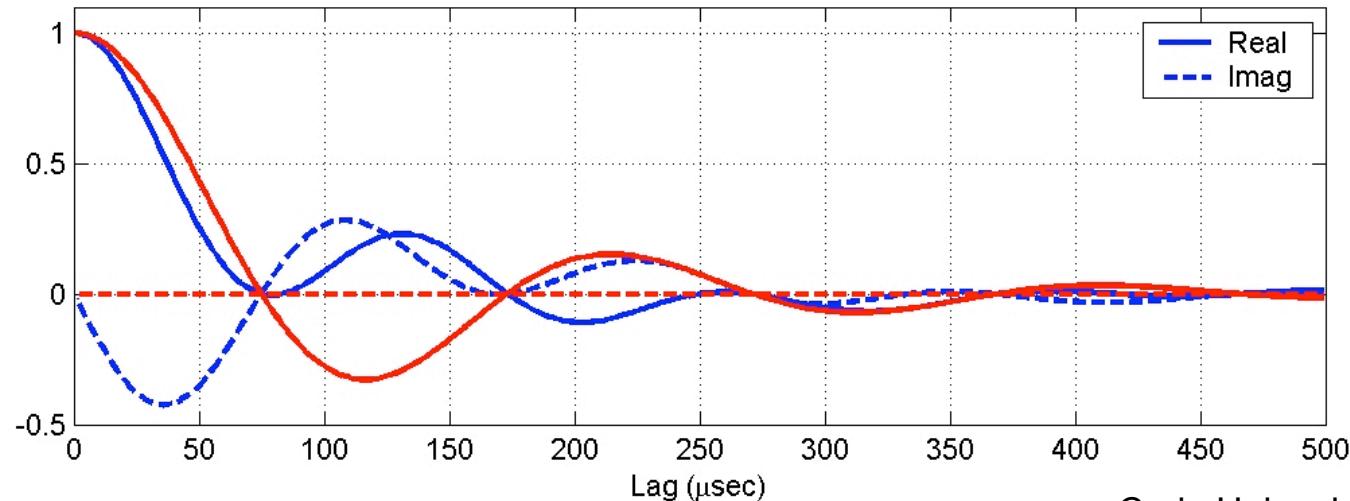
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Ion Velocity



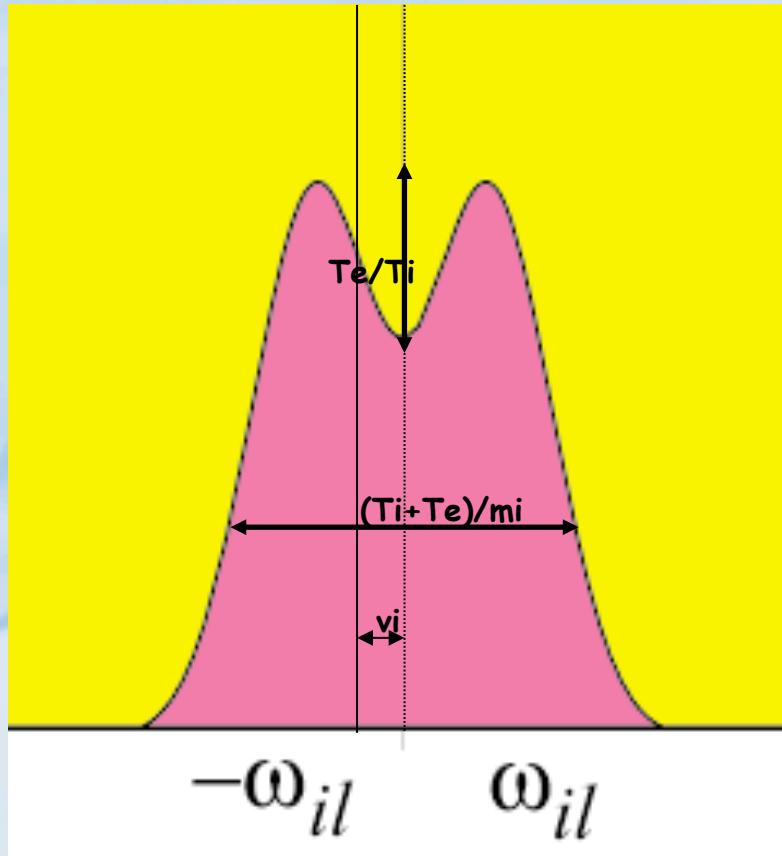
Parameters

Freq: 449 MHz
 Ne: 10^{12} m^{-3}
 Ti: 1000 K
 Te: 2000 K
 Comp: 100% O⁺
 v_{in} : 10^{-6} KHz



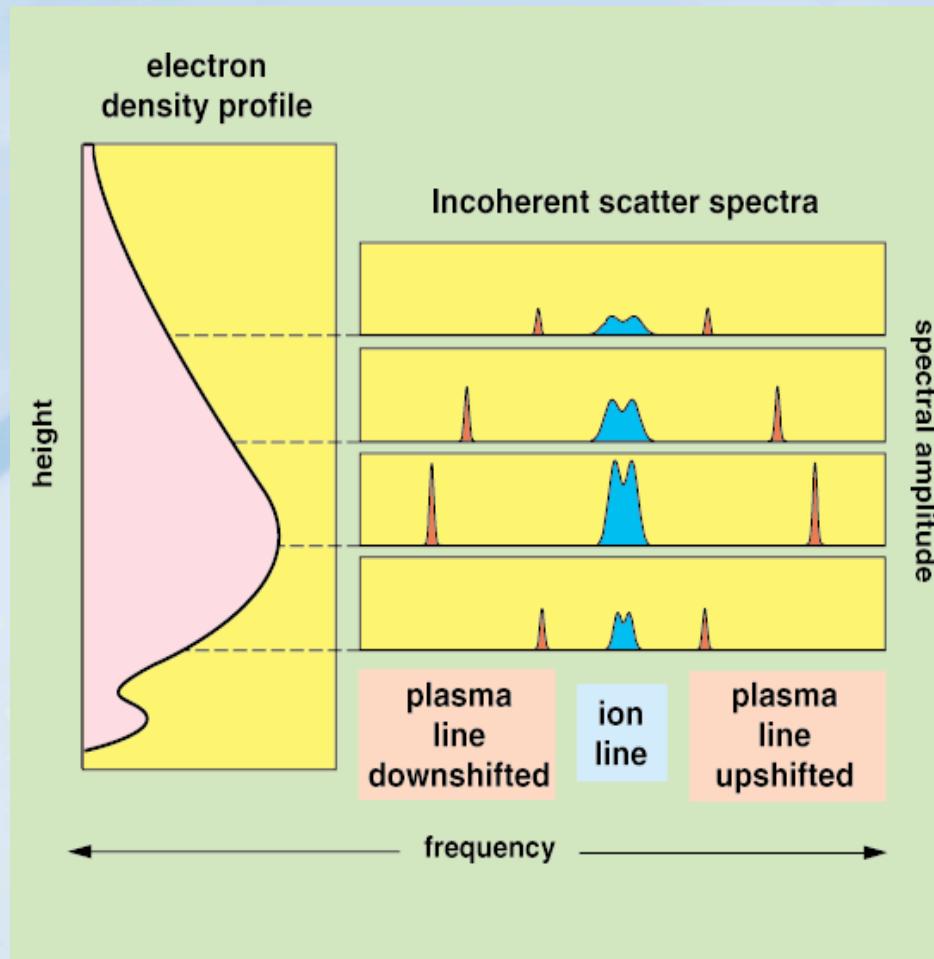
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...or to sum up...

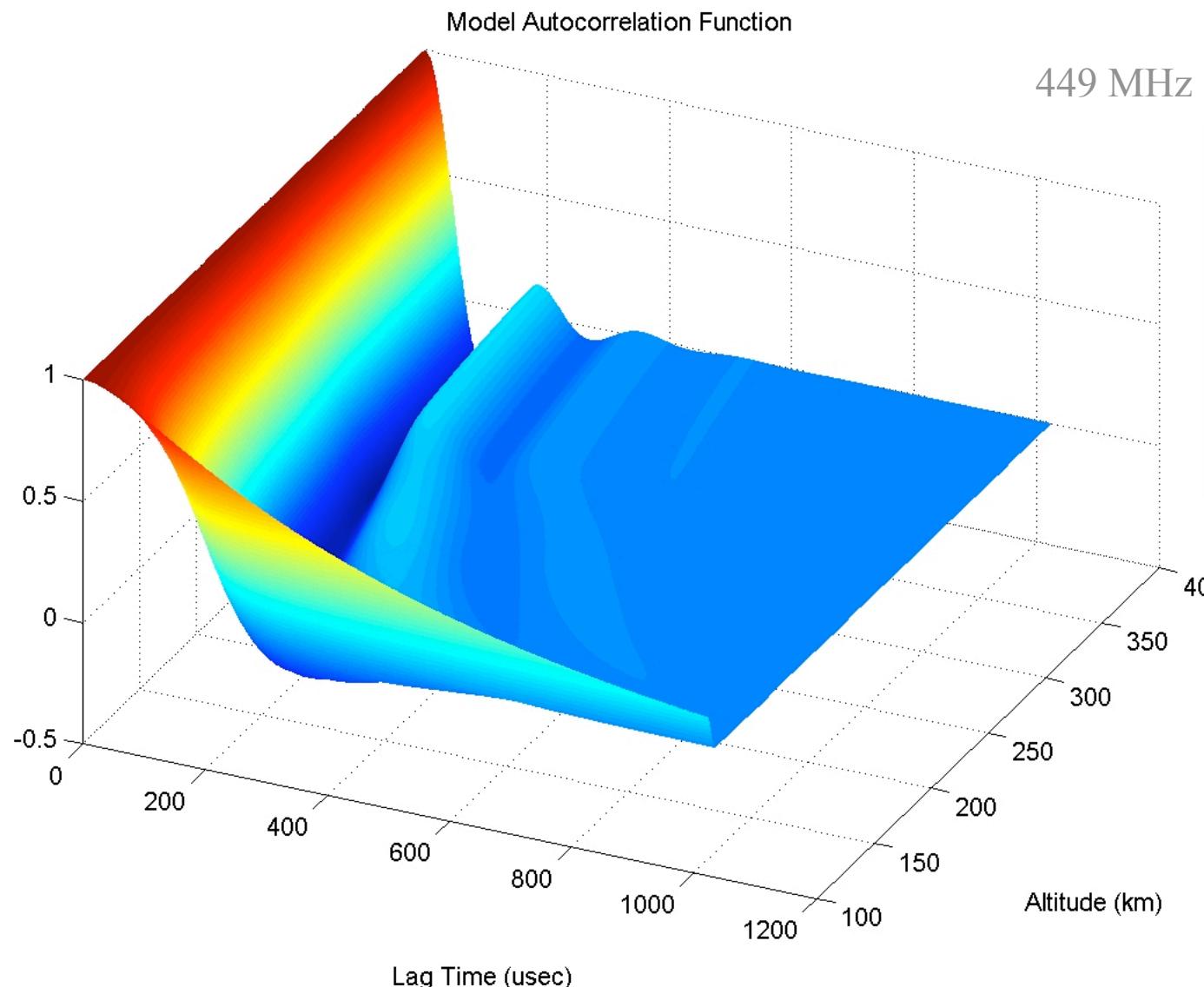


- Ion (and electron) temperature (T_i and T_e) to ion mass (m_i) ratio from the width of the spectra
- Electron to ion temperature ratio (T_e/T_i) from “peak_to_valley” ratio
- Electron (= ion) density from total area (corrected for temperatures)
- Ion velocity (v_i) from the Doppler shift

Spectral space as a function of altitude



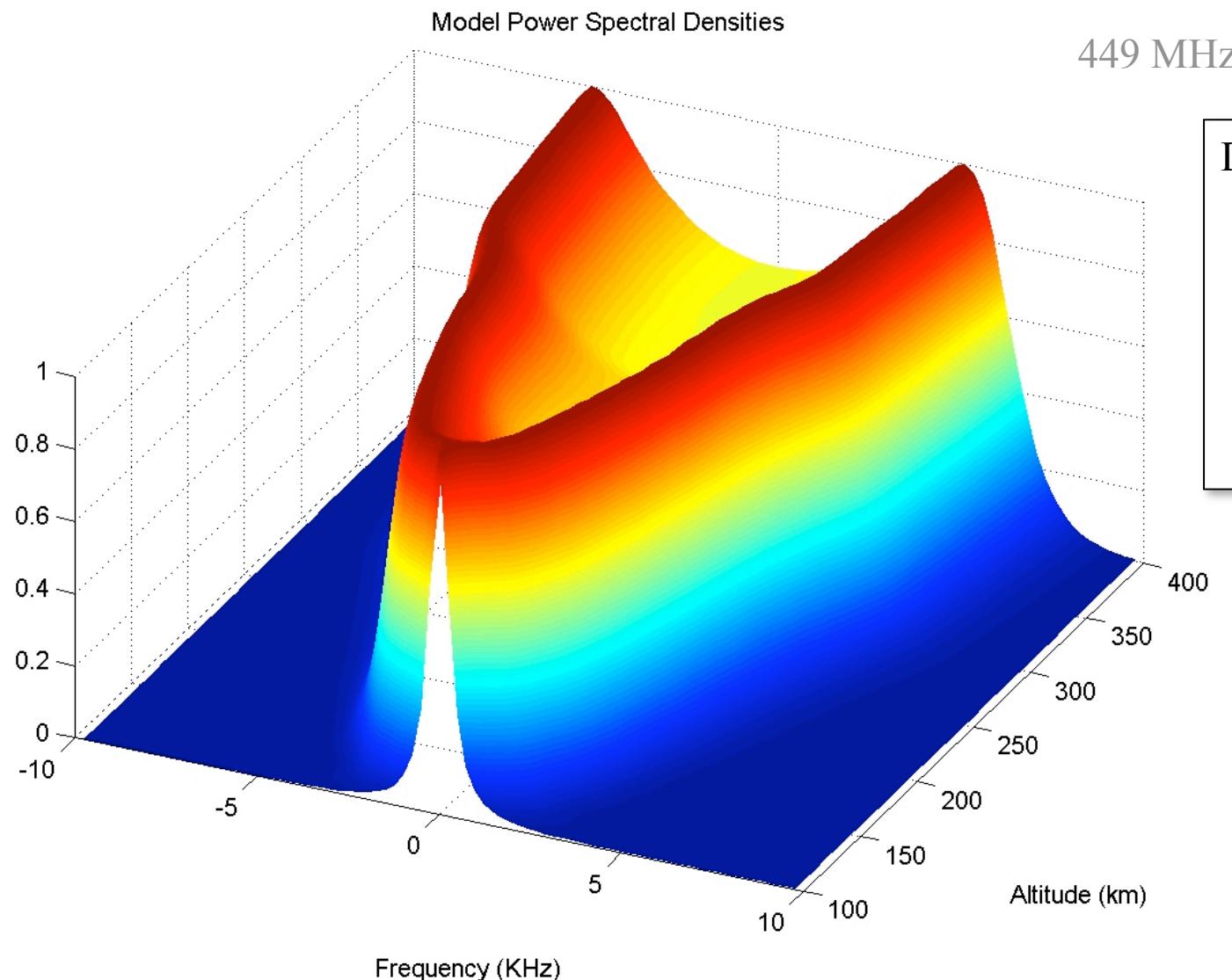
Incoherent Scatter Autocorrelation Functions



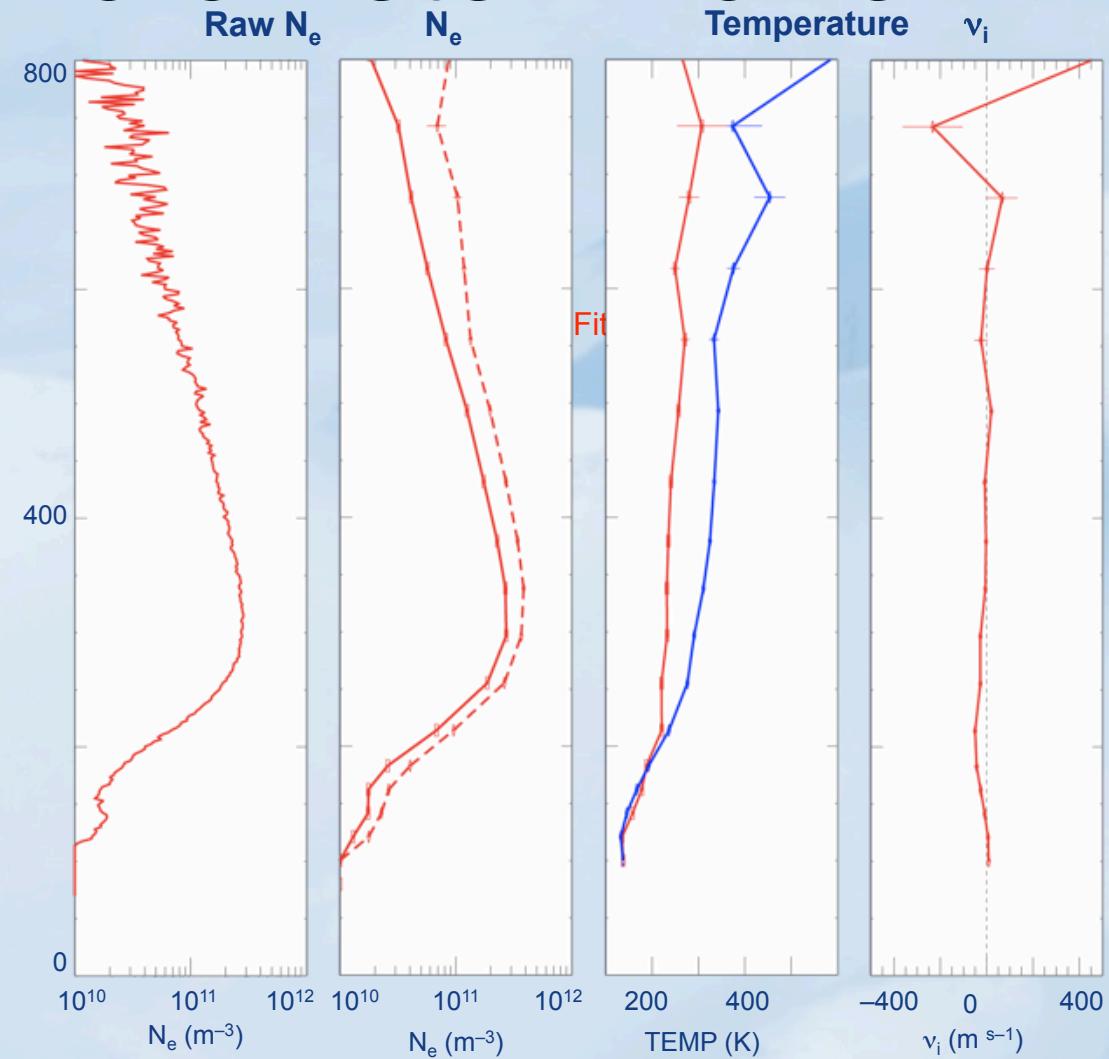
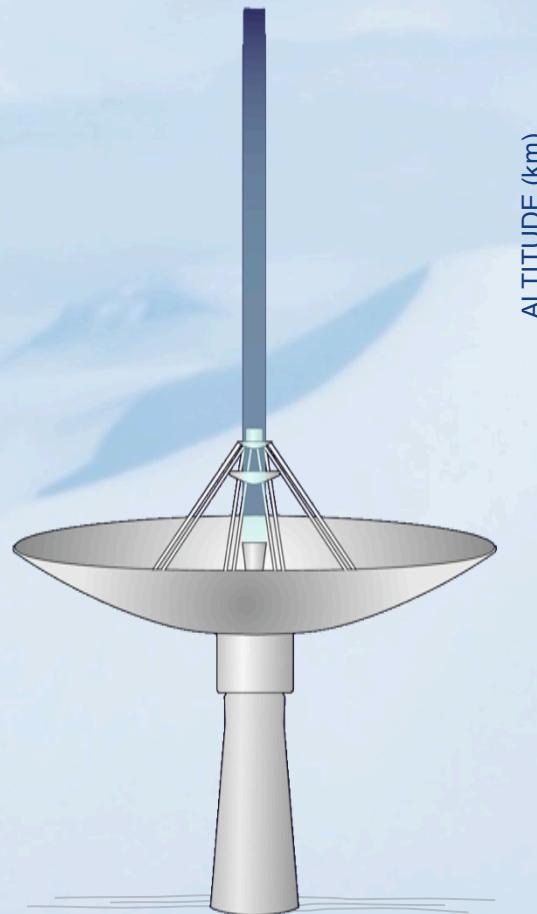
Dependencies:

- N_e
- T_e
- T_i
- V_i
- m_i
- v_{in}

Incoherent Scatter Power Spectra



Plasma Parameter Profile

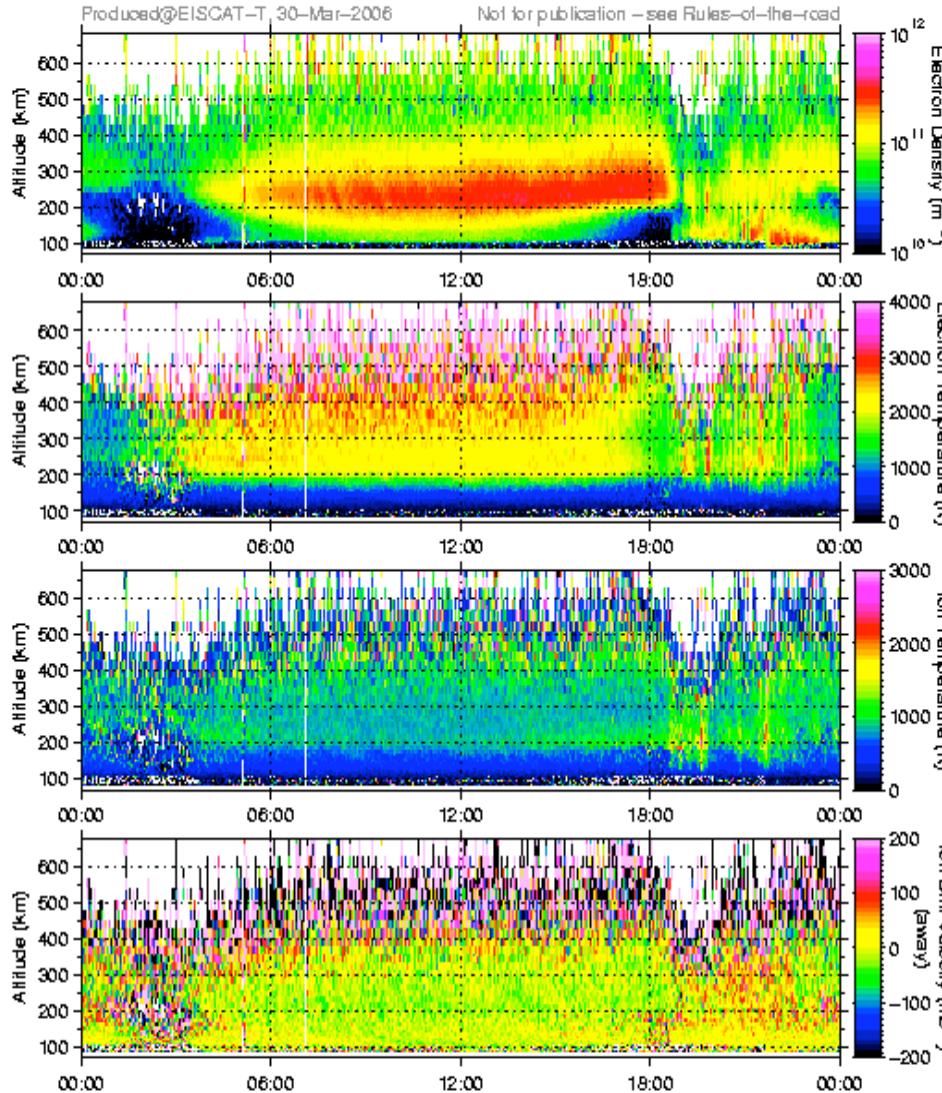


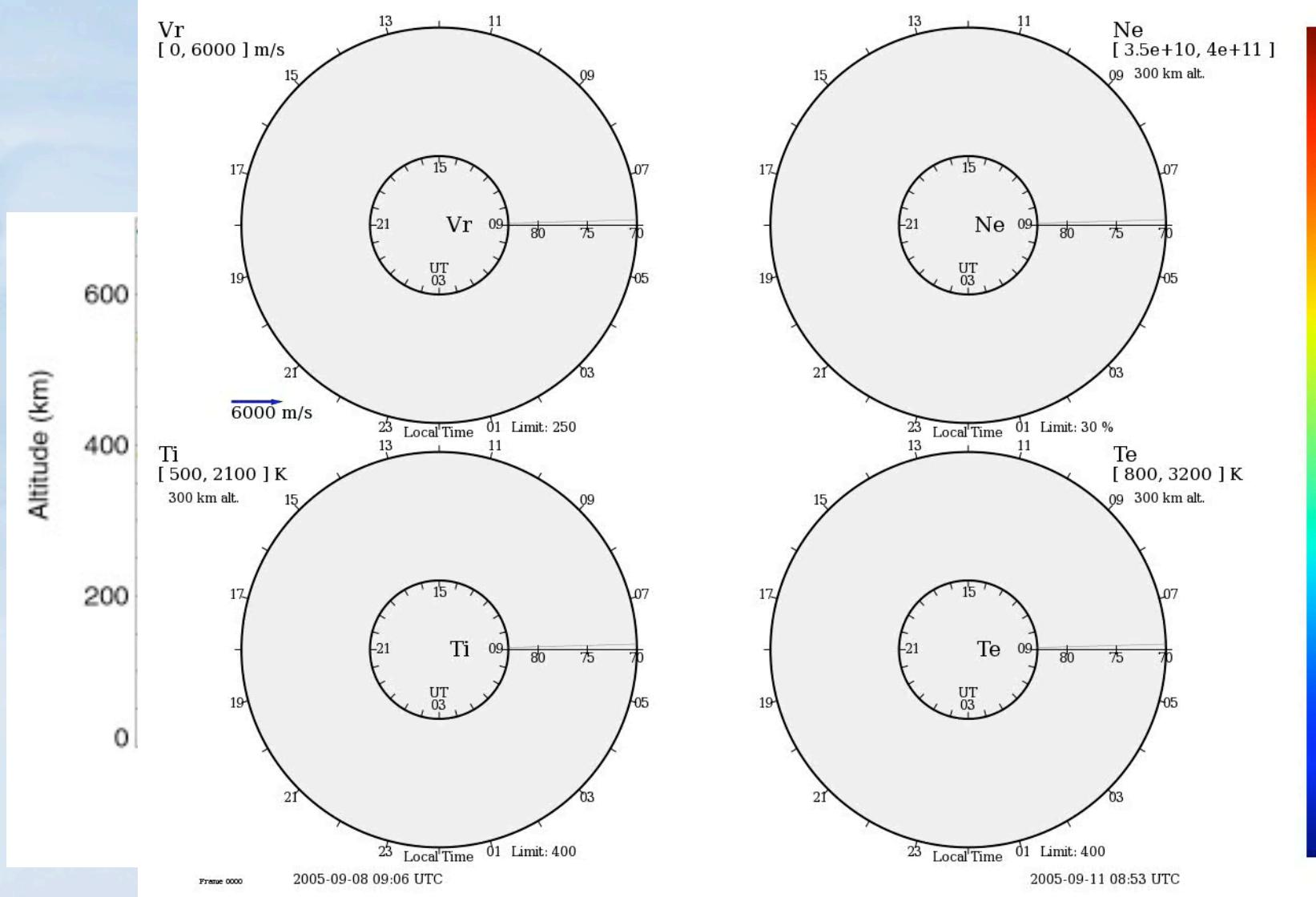


EISCAT Scientific Association

EISCAT UHF RADAR

CP, uhf, tau2pl, 29 March 2006







And this is the level data we
will work on in the MADRIGAL
session...