

# More on ISR Experiments, Data Reduction, and Analysis

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# Outline

## 1 ISR Pulses and Experiments

- The Nature of the IS Target
- *F*-Region Experiments
- *E*-Region Experiments
- *D*-Region Experiments
- System Info
- Beam Pointing

## 2 Level-0 Processing

- General
- Power Estimation
- ACF / Spectra Estimation

## 3 Level-1 Processing

- $N_e$  Estimation
- ACF / Spectral Fits
- ACF / Spectral Fits

## 4 Level-2 Processing

- Vector Velocities / Electric Fields
- *E*-Region Winds
- Collision Freqs. / Conductivities / Currents / Joule Heating
- *D*-Region Parameters

## Overspread Targets

(a.k.a, frequency and range aliased targets)

- For a target with a bandwidth  $B$ , you must sample at a rate  $F$  exceeding  $B$  (e.g., for IS,  $B \sim 40$  kHz).
- For a target which could be as far away as  $R_{max}$ , radar pulses must be at least  $2R_{max}/c$  apart.

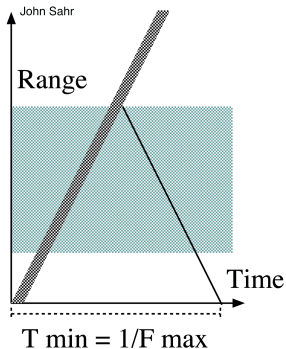
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Thus, there is a competition between distance and bandwidth

- $B < F < \frac{c}{2R_{max}}$
- or:  $B \frac{2R_{max}}{c} < 1$
- At 450 MHz,  $B \sim 40$  kHz,  
 $R \sim 750$  km (5 ms)  $\rightarrow$  highly overspread
- Do we get the range right or the spectrum right??





# Overspread Targets

How is this resolved?

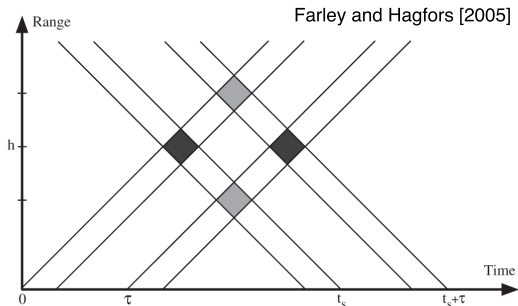
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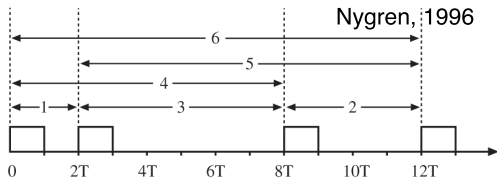
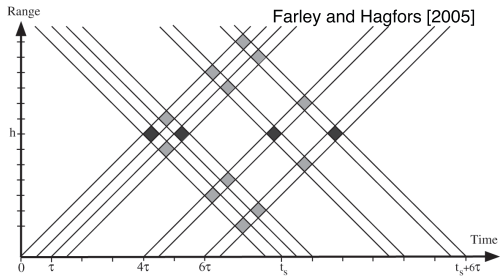
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Simplest scheme to measure correlation at a given lag - double pulse:

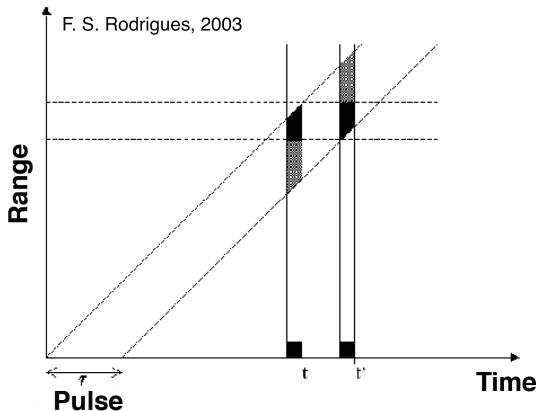


# Generalization - Multipulses



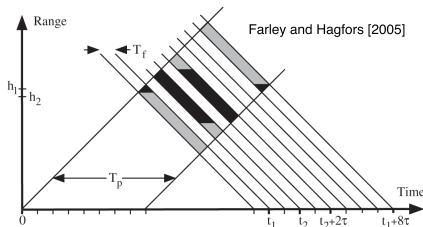
# Ambiguity Function

Long-pulse of length  $\tau$ , sampled at  $t$  and  $t'$  with a box-car impulse response.



- Range ambiguity function is box-car shaped.
- Lag ambiguity is triangular shaped.

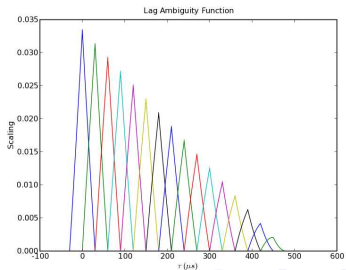
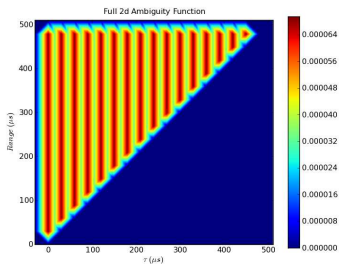
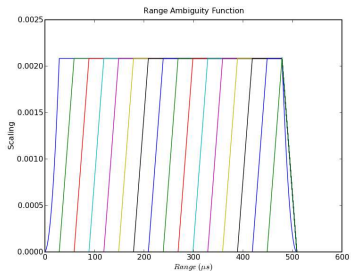
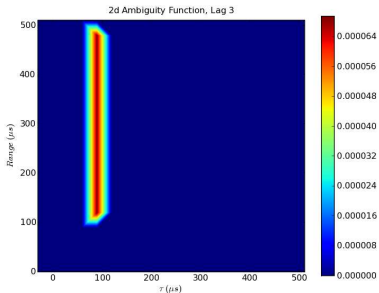
## Standard F-region Experiment - Long Pulse



- At high altitudes, use a single long pulse with mismatched filter (oversampled) to measure all lags of the ACF at once
- Sacrifice range resolution
- Typically use a  $480 \mu\text{s}$  pulse (*F* region) or 1 ms pulse (topside)

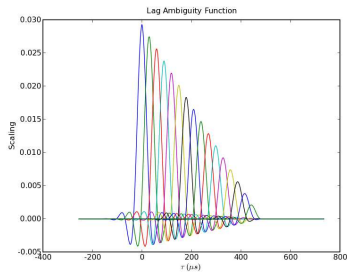
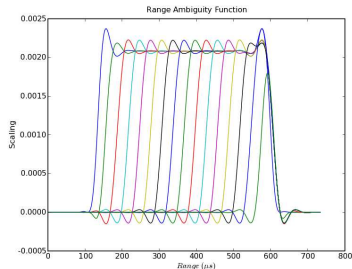
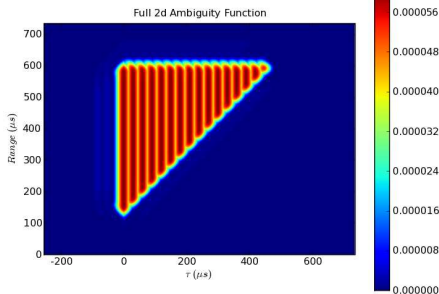
# Long Pulse Ambiguity Function

Ambiguity function with a boxcar filter. 480  $\mu\text{s}$  long pulse, 30  $\mu\text{s}$  sampling.



# Long Pulse Ambiguity Function

- Ambiguity function including filter effects.
- 480  $\mu\text{s}$  long pulse, 30  $\mu\text{s}$  sampling.
- With filter effects.



# Long Pulse Gating

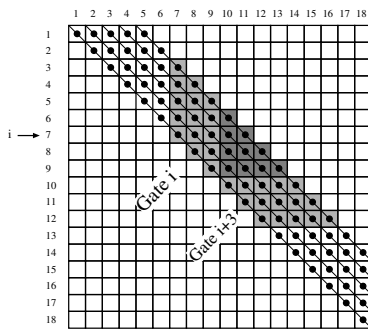
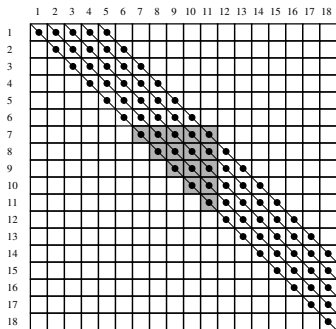
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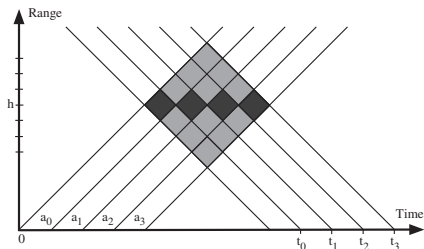
“Simple solution” - Gating using elements of the so-called lag-profile matrix.



*Nygren, 1996*

A better method - treat as an inverse problem: deconvolution or full profile methodologies. Active area of research.

## Standard E-region Experiment - Coded Pulse



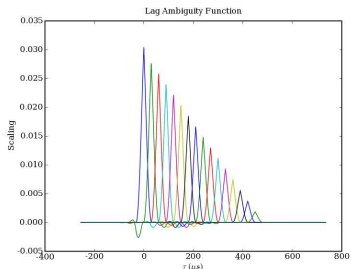
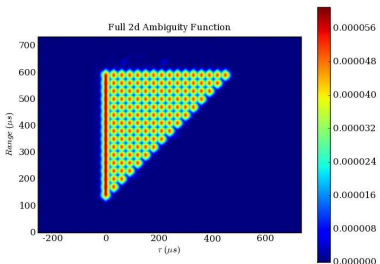
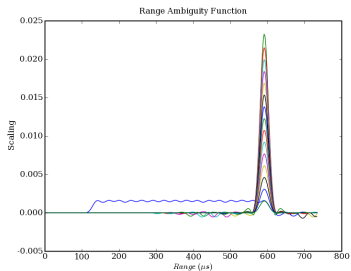
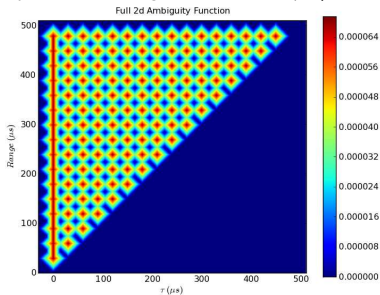
*Farley and Hagfors [2005]*

E.g., consider lag estimate using  $v(t_0)$  and  $v^*(t_1)$  - choose  $a_n$  such that clutter terms cancel.

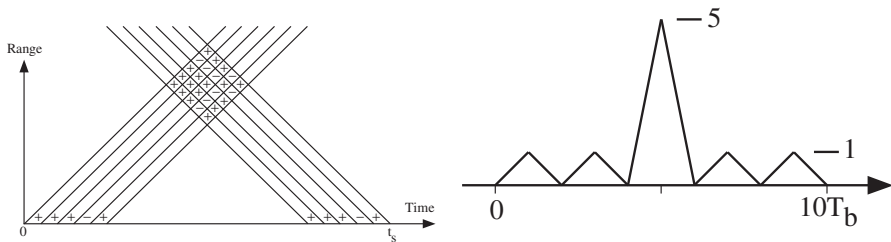
- At lower altitudes, we require better range resolution.
- For this, we utilize binary coded pulse ACF measurements (do not compress pulse or eliminate clutter like BC - eliminate correlation of clutter)
- Random (CLP) or alternating (cyclic codes)
- Standard experiment is  $480 \mu\text{s}$ , 16-baud (4.5 km), randomized strong code.
- Include an uncoded  $30 \mu\text{s}$  pulse for zero-lag normalization.

# Standard $E$ -region Experiment - Ambiguity Function

Ambiguity function including filter effects. 480  $\mu\text{s}$  (16-baud, 30  $\mu\text{s}$  baud, 32 pulse).



## Standard E/F-region Power Measurement



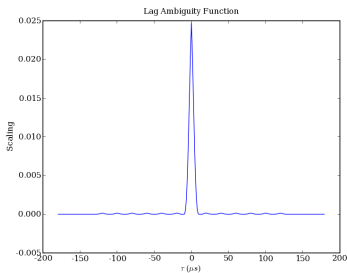
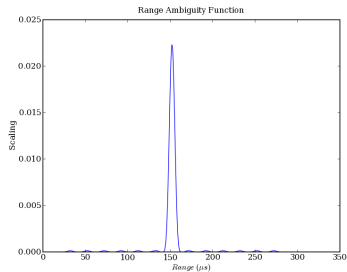
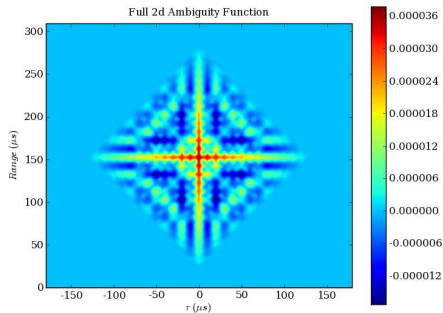
*Farley and Hagfors [2005]*

- Pulse compression code allow for high sensitivity, high range resolution power measurements.
- Plasma must remain correlated over pulse length (limits range of use for most systems).
- Typical code is 13-baud Barker code,  $130 \mu\text{s}$ .

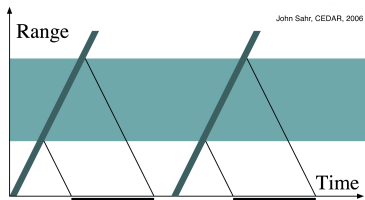
# E/F-region Power Measurement - Ambiguity Function

Ambiguity function including filter effects.

130  $\mu\text{s}$  (13-baud, 10  $\mu\text{s}$  baud, 5  $\mu\text{s}$  sampling).



## Standard *D*-region Experiments



- Long correlation times (narrow spectral widths) in the *D* region require pulse-to-pulse techniques
- We employ coded double-pulse techniques that give range resolutions up to 600 m and spectral resolutions up to 1 Hz.

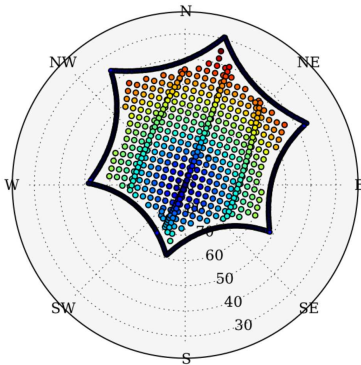
Mode	Pulse	Baud	$\delta R$	$\tau$	IPP	$\delta f$	Nyquist	$\delta t$
0	130 $\mu s$	10 $\mu s$	1.5 km	5 $\mu s$ (0.75 km)	2 ms	2 Hz	250 Hz	1 s
1	260 $\mu s$	10 $\mu s$	1.5 km	5 $\mu s$ (0.75 km)	4 ms	1 Hz	125 Hz	2.5 s
2	130 $\mu s$	10 $\mu s$	1.5 km	5 $\mu s$ (0.75 km)	2 ms	2 Hz	250 Hz	1.8 s
3	280 $\mu s$	10 $\mu s$	1.5 km	5 $\mu s$ (0.75 km)	3 ms	1.3 Hz	167 Hz	2.7 s
4	112 $\mu s$	4 $\mu s$	0.6 km	2 $\mu s$ (0.3 km)	3 ms	1.3 Hz	167 Hz	2.7 s

# System Information

- 128-panel AMISR system (upgraded from 96 in Sep. 07)
- Pulse-to-pulse phase capability
- $\sim 1.6$  MW peak Tx (upgraded from  $\sim 1.3$  MW)
- $\sim 10\%$  max duty cycle
- 4 reception channels
- Tx band 449-450 MHz
- 3.5 MHz max Rx bandwidth
- $4 \mu\text{s}$  min pulsewidth (freq. allocation limitation)
- Fully programmable, remotely operable/ed
- Graceful degradation - reliable operations



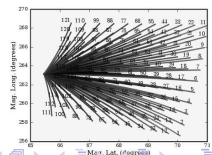
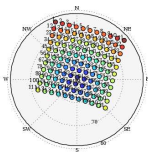
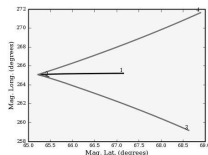
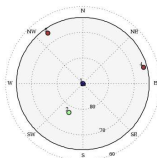
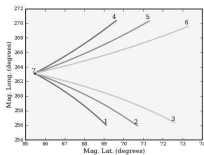
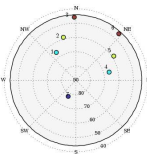
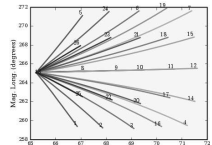
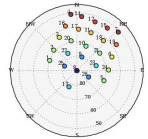
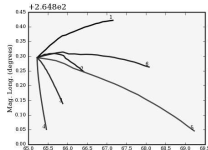
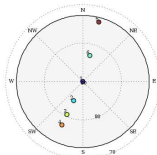
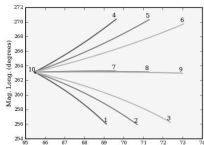
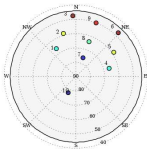
# Beam Pointing



- Range of pointing positions within grating lobe limits
- "Normal" experiments include  $\sim 1-10$  beams
- Main limitation is integration time / sensitivity



# Beam Pointing



# General

A typical experiment consists of:

- Data samples
- Noise samples
- Cal pulse samples

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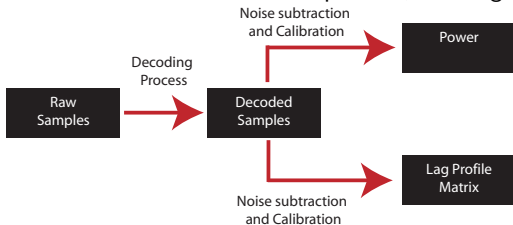
- Data samples
  - Noise samples
  - Cal pulse samples
- Interleaving of pulses (possibly on different frequencies)
  - Clutter considerations, Noise & Cal sample placement
  - Maximization of duty cycle
  - Beam pointing, Distribution of pulses, Integration time considerations
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# Power Estimation

Received power can be written as

$$P_r = \frac{P_t \tau_p}{r^2} K_{\text{sys}} \frac{N_e}{(1 + k^2 \lambda_D^2)(1 + k^2 \lambda_D^2 + T_r)} \text{ Watts}$$

where

$P_r$  - received power (Watts)

$P_t$  - transmit power (Watts)

$\tau_p$  - pulse length (seconds)

$r$  - range (meters)

$N_e$  - electron density ( $\text{m}^{-3}$ )

$k$  - Bragg scattering wavenumber (rad/m)

$\lambda_D$  - Debye length (m)

$T_r$  - electron to ion temperature ratio

$K_{\text{sys}}$  - system constant ( $\text{m}^5/\text{s}$ )

# Power Estimation

Received signal power needs to be calibrated to absolute units of Watts. To do this, we in general (a) take noise samples and (b) inject a calibration pulse at each AEU, which is then summed in the same way as the signal. The absolute calibration power in Watts is:

$$P_{cal} = k_B T_{cal} B \quad \text{Watts}$$

where

$k_B$  - Boltzmann constant (J/kg K)

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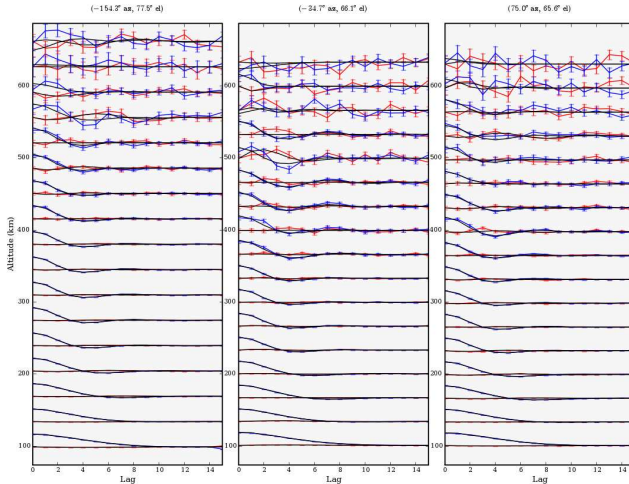
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The measurement of the calibration power (after noise subtraction) can then be used as a yardstick to convert the received power to Watts. This is done as,

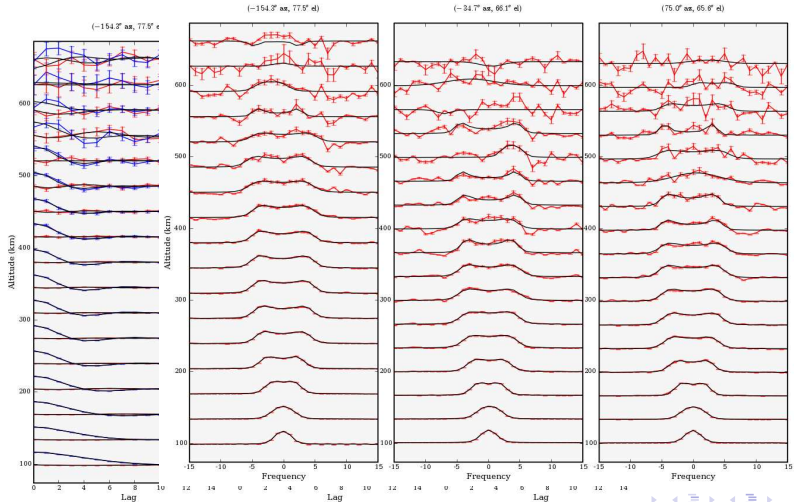
$$P_r = P_{cal} * (\text{Signal} - \text{Noise}) / (\text{Cal} - \text{Noise}) \quad \text{Watts}$$

# ACF / Spectra Estimation - E/F region

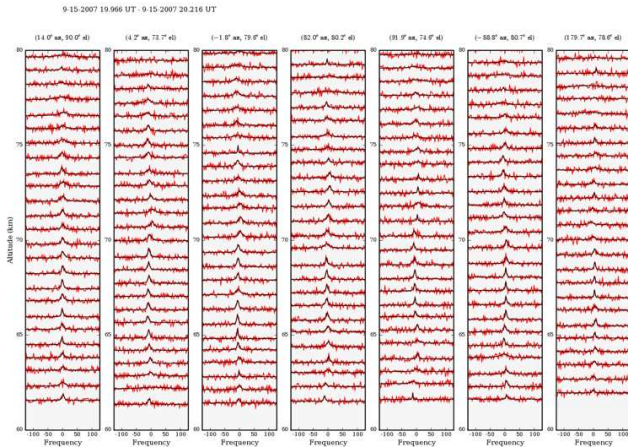




# ACF / Spectra Estimation - E/F region



# ACF / Spectra Estimation - $D$ region



# Electron Density

Recall,

$$P_r = \frac{P_t \tau_p}{r^2} K_{\text{sys}} \frac{N_e}{(1 + k^2 \lambda_D^2)(1 + k^2 \lambda_D^2 + T_r)} \text{ Watts}$$

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Within  $K_{\text{sys}}$  is embedded information on the gain, which for a phased-array varies with the look-angle off boresight, as well as the proximity to the grating lobe limits.

# Electron Density

$$f_r^2 \approx f_p^2 + \frac{3k^2}{4\pi^2} \frac{k_B T_e}{m_e} + f_c^2 \sin^2 \alpha$$

where

$f_r$  - plasma line frequency (Hz)

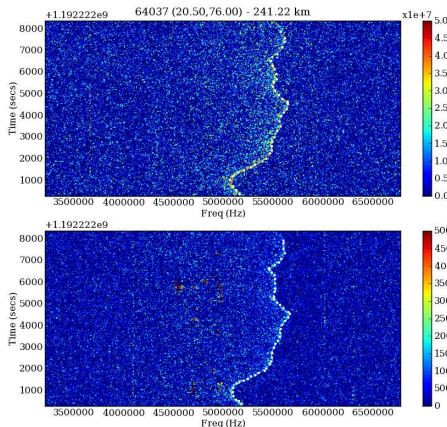
$f_p$  - plasma frequency (Hz)

$T_e$  - electron temperature (K)

$m_e$  - electron mass (kg)

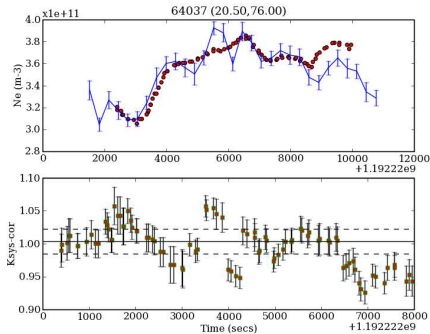
$f_c$  - electron cyclotron frequency (Hz)

$\alpha$  - magnetic aspect angle

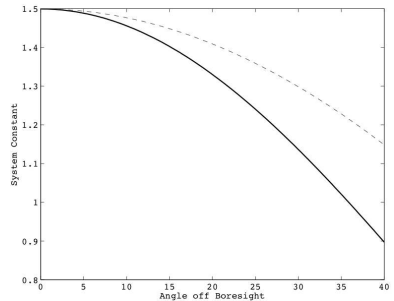


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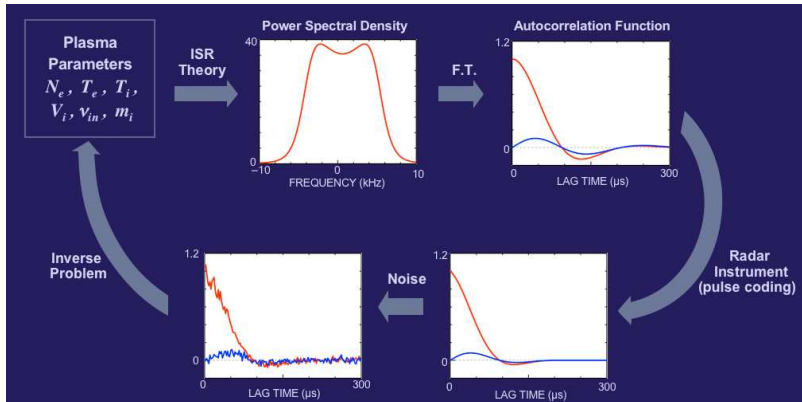
$$K_{\text{sys}} = A \cos^B(\theta_{BS}) \quad \text{m}^5/\text{s}$$



$\theta_{BS}$  - angle off boresight  
 $A, B$  - constants



# Fitting Spectra





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## General Complicating Factors:

- Range smearing
- Lag smearing
- Pulse coding effects / "Self"-clutter
- Clutter (geophysical and not - e.g., mountains, irregularities, turbulence, non-Maxwellian)
- Signal strength / statistics
- Time stationarity

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## Specific Based on Altitude:

- *F*-region/Topside - Light ion composition
- Bottomside - Molecular ion composition
- *E*-region - Collision frequency, Temperature
- *D*-region - Complete ambiguity

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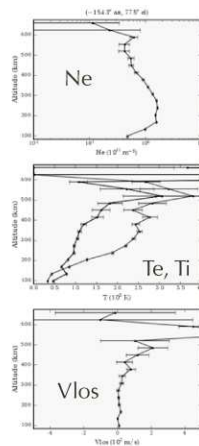
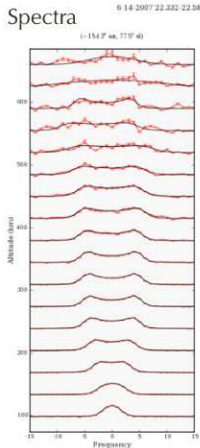
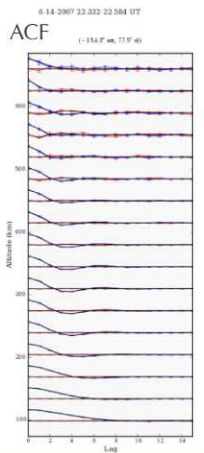
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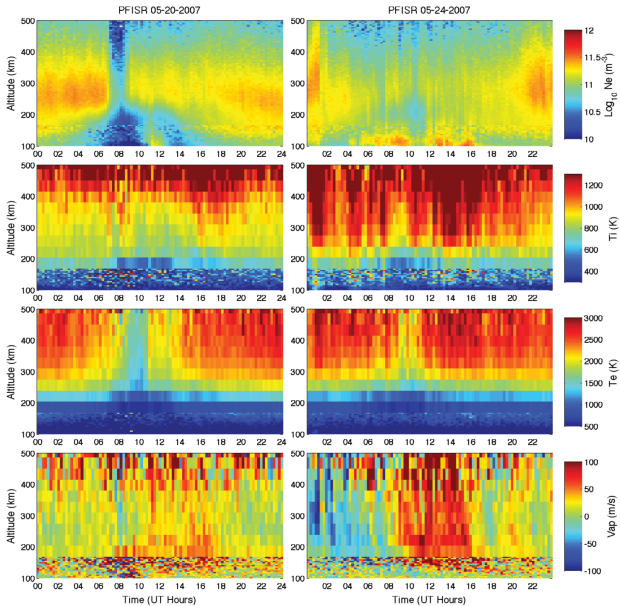
## Approach:

- *F*-region -  $T_e$ ,  $T_i$ ,  $v_{los}$ ,  $N_e$
- Bottomside - Assume a composition profile
- *E*-region -  $< \sim 105km$ , assume  $T_e = T_i$
- *D*-region - Fit a Lorentzian (width, Doppler,  $N_e$ )

# Fitting Spectra - Example



# Fitting Spectra - Example



## Vector Velocities - Preliminaries

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If we can neglect Earth curvature (“high enough” elevation angles),

$$\mathbf{k} = \begin{bmatrix} k_e \\ k_n \\ k_z \end{bmatrix} = \begin{bmatrix} \cos \theta \sin \phi \\ \cos \theta \cos \phi \\ \sin \theta \end{bmatrix}$$

where  $\theta$ ,  $\phi$  are elevation and azimuth angles, respectively.



## Vector Velocities - Preliminaries

For a local geomagnetic coordinate system we can use the rotation matrix,

$$R_{geo \rightarrow gmag} = \begin{bmatrix} \cos \delta & -\sin \delta & 0 \\ \sin I \sin \delta & \cos \delta \sin I & \cos I \\ -\cos I \sin \delta & -\cos I \cos \delta & \sin I \end{bmatrix}$$

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$$\mathbf{k} = \begin{bmatrix} k_{pe} \\ k_{pn} \\ k_{ap} \end{bmatrix} = \begin{bmatrix} k_e \cos \delta - k_n \sin \delta \\ k_z \cos I + \sin I (k_n \cos \delta + k_e \sin \delta) \\ k_z \sin I - \cos I (k_n \cos \delta + k_e \sin \delta) \end{bmatrix}.$$

## Vector Velocities - Two Point

Two LOS velocity measurements can be written as,

$$\begin{bmatrix} v_{los}^1 \\ v_{los}^2 \end{bmatrix} = \begin{bmatrix} k_{pe}^1 & k_{pn}^1 & k_{ap}^1 \\ k_{pe}^2 & k_{pn}^2 & k_{ap}^2 \end{bmatrix} \begin{bmatrix} v_{pe} \\ v_{pn} \\ v_{ap} \end{bmatrix}$$

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Can be solved for  $v_{pn}$  and  $v_{pe}$  assuming  $v_{ap} \approx 0$ ,

$$v_{pn} = \frac{v_{los}^1 - \frac{k_{pe}^1}{k_{pe}^2} v_{los}^2 - v_{ap} \left( k_{ap}^1 - k_{ap}^2 \frac{k_{pe}^1}{k_{pe}^2} \right)}{k_{pn}^1 \left( 1 - \frac{k_{pn}^2}{k_{pn}^1} \frac{k_{pe}^1}{k_{pe}^2} \right)} \approx \frac{v_{los}^1 - \frac{k_{pe}^1}{k_{pe}^2} v_{los}^2}{k_{pn}^1 \left( 1 - \frac{k_{pn}^2}{k_{pn}^1} \frac{k_{pe}^1}{k_{pe}^2} \right)}$$

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Implies that you need look directions with different  $\mathbf{k}$  vectors.

# Vector Velocities - Generalization

Multiple measurements can be written as,

$$\begin{bmatrix} v_{los}^1 \\ v_{los}^2 \\ \vdots \\ v_{los}^n \end{bmatrix} = \begin{bmatrix} k_{pe}^1 & k_{pn}^1 & k_{ap}^1 \\ k_{pe}^2 & k_{pn}^2 & k_{ap}^2 \\ \vdots & \vdots & \vdots \\ k_{pe}^n & k_{pn}^n & k_{ap}^n \end{bmatrix} \begin{bmatrix} v_{pe} \\ v_{pn} \\ v_{ap} \end{bmatrix} + \begin{bmatrix} e_{los}^1 \\ e_{los}^2 \\ \vdots \\ e_{los}^n \end{bmatrix}$$

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$$\hat{\mathbf{v}}_i = \Sigma_v \mathbf{A}^T (\mathbf{A} \Sigma_v \mathbf{A}^T + \Sigma_e)^{-1} \mathbf{v}_{los}$$

Error covariance,

$$\Sigma_{\hat{\mathbf{v}}} = \Sigma_v - \Sigma_v \mathbf{A}^T (\mathbf{A} \Sigma_v \mathbf{A}^T + \Sigma_e)^{-1} \mathbf{A} \Sigma_v = (\mathbf{A}^T \Sigma_e^{-1} \mathbf{A} + \Sigma_v^{-1})^{-1}$$

# Electric Fields

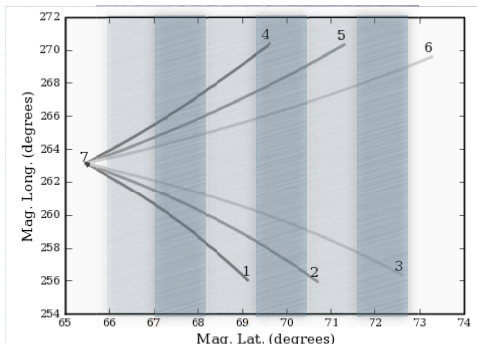
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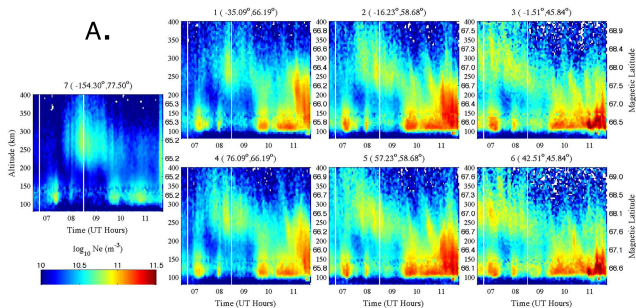
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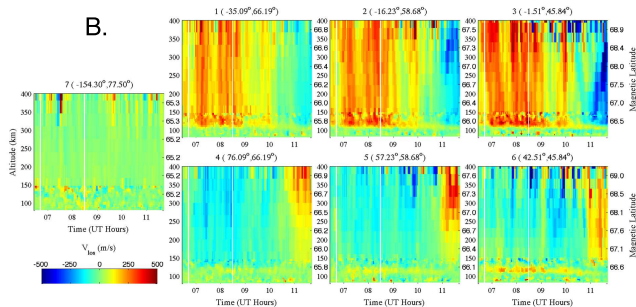
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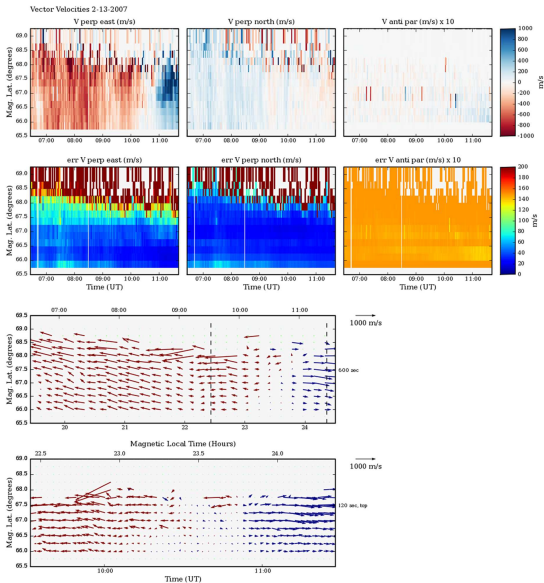
Electron Density



LOS Velocities

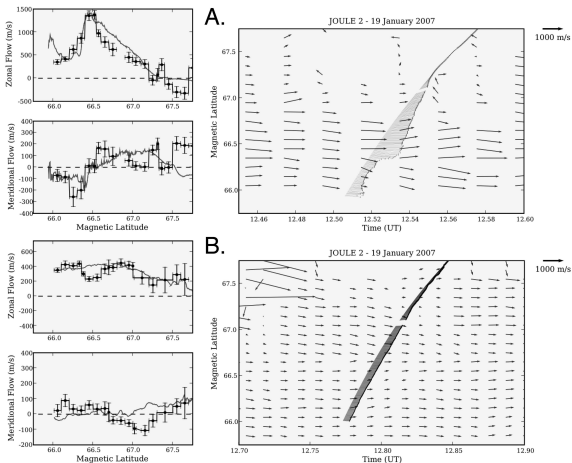
# Electric Fields - Example

## Resolved Vectors



# Electric Fields - Example

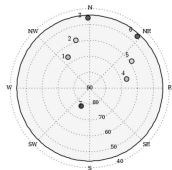
## Comparison to rocket-measured E-fields.



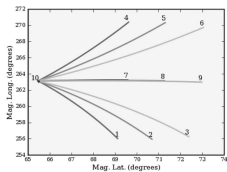
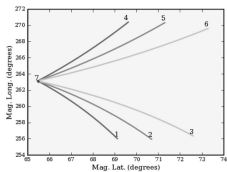
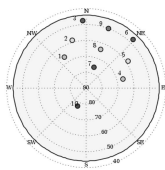
# Experiment Planning

The approach also allows for an efficient means of experiment planning, since the output covariance of the measurements is independent of the actual measurements.

A.



B.

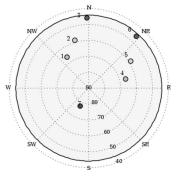




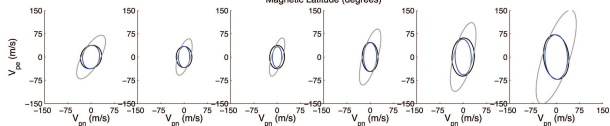
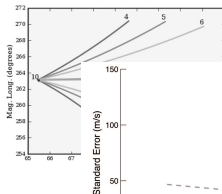
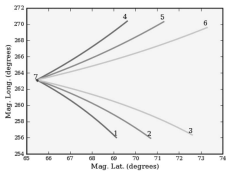
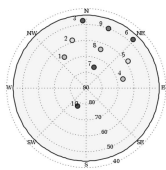
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The steady state ion momentum equations relate the vector velocities (as a function of altitude) to electric fields and neutral winds

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$$\mathbf{x} = [E_{pe} \ E_{pn} \ E_{||} \ u_{pe}^1 \ u_{pn}^1 \ u_{||}^1 \ u_{pe}^2 \ u_{pn}^2 \ u_{||}^2 \ \dots \ u_{pe}^n \ u_{pn}^n \ u_{||}^n]^T$$

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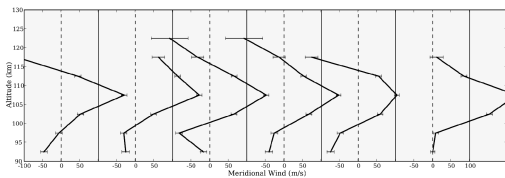
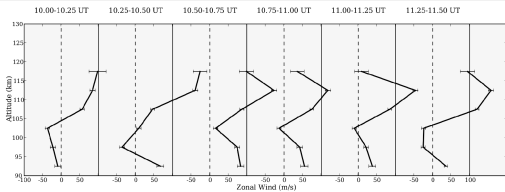
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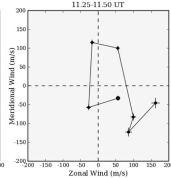
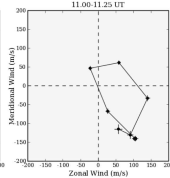
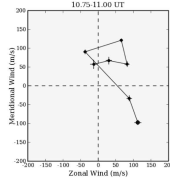
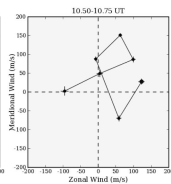
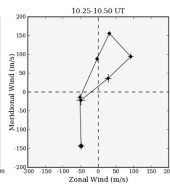
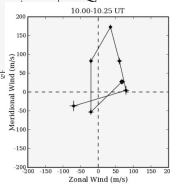
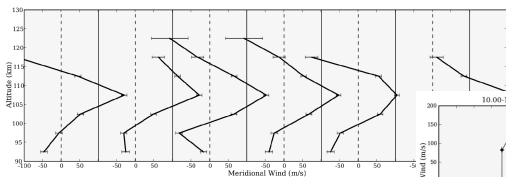
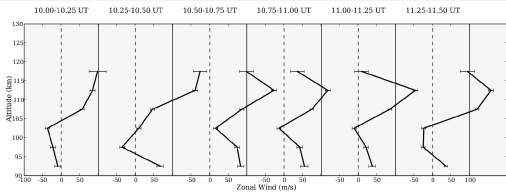
This allows for direct constraint of both the vertical wind and the parallel electric field, both of which we expect to be small.

$$\sum_v^{gmag} = J_{geo \rightarrow gmag} \sum_v^{geo} J_{geo \rightarrow gmag}^T$$

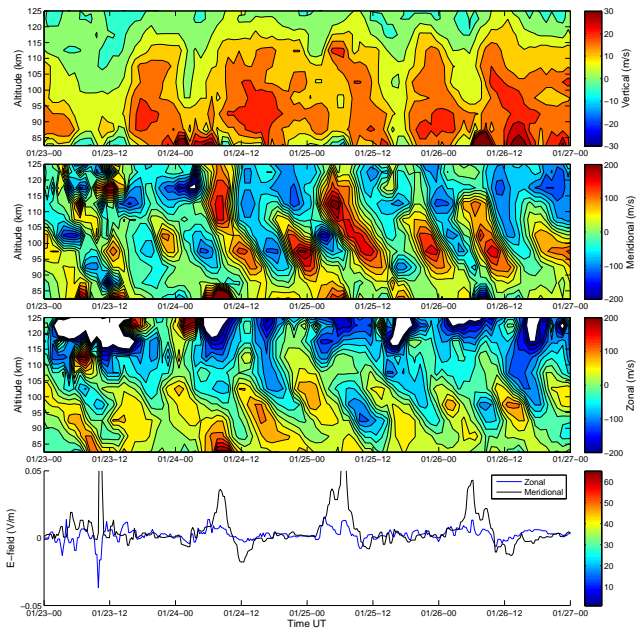
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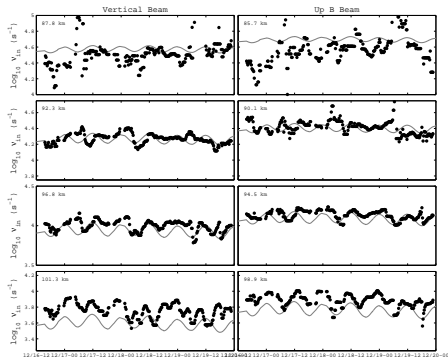
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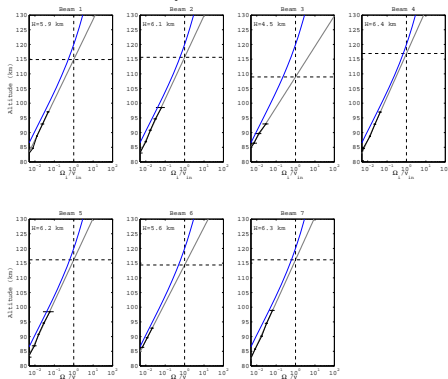
- 1 Direct fits at lower altitudes (spectral width  $\sim \propto T_n / \nu_{in}$ )
- 2 Examination of variation of LOS velocity with altitude

# Collision Frequency - Method 1

Semi-diurnal variation over several days.



Altitude profile and extrapolation.



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Perp-north and parallel components given by,

$$v_{\perp n} = \kappa_i (1 + \kappa_i^2)^{-1} (b_i E_{\perp e} + u_{\perp e}) + (1 + \kappa_i^2)^{-1} (b_i E_{\perp n} + u_{\perp n})$$

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Under strong convection (electric field) conditions, neglect winds

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$$v'_z \sim b_i(1 + \kappa_i^2)^{-1} [\kappa_i E_{\perp e} + E_{\perp n}] \cos I$$

If  $\kappa_i(z) = \kappa_0 e^{(z-z_0)/H}$ , vertical ion velocity will maximize at

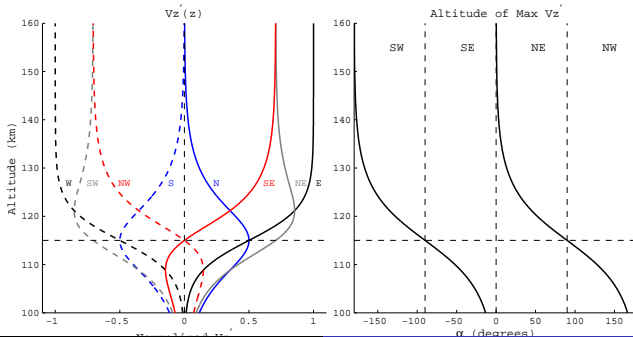
$$z_{\max} v'_z = z_0 + H \ln \kappa_0^{-1} + H \ln \left[ \frac{\cos \alpha \pm 1}{\sin \alpha} \right]$$

## Collision Frequency - Method 2 - Example

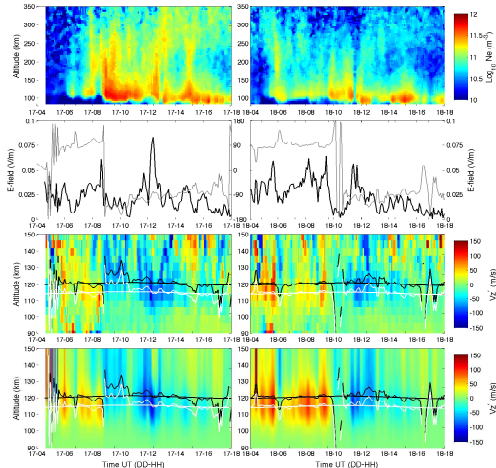
$$v'_z \sim b_i(1 + \kappa_i^2)^{-1} [\kappa_i E_{\perp e} + E_{\perp n}] \cos I$$

If  $\kappa_i(z) = \kappa_0 e^{(z-z_0)/H}$ , vertical ion velocity will maximize at

$$z_{\max} v'_z = z_0 + H \ln \kappa_0^{-1} + H \ln \left[ \frac{\cos \alpha \pm 1}{\sin \alpha} \right]$$

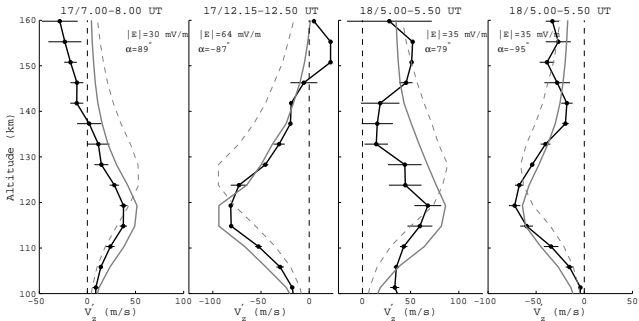


# Collision Frequency - Method 2

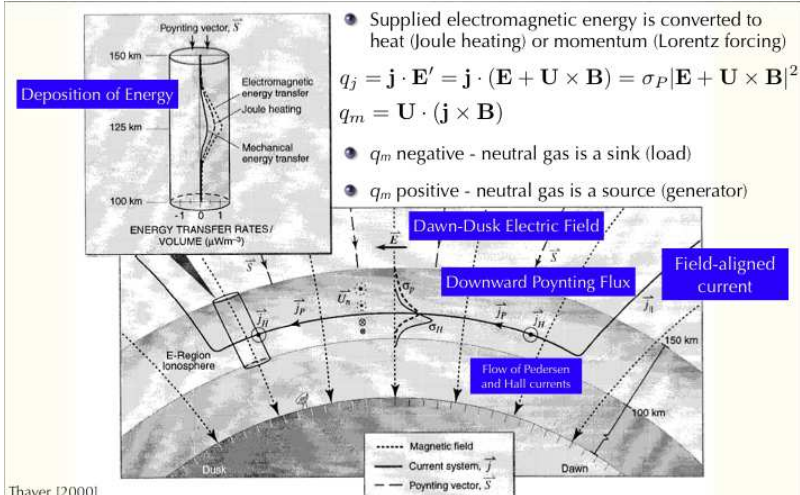


## Collision Frequency - Method 2

Profiles of  $v_z'$  during high convection conditions.  
 Dashed - with MSIS; Solid - scaled by a factor of 2.



# Conductivities / Currents / Joule Heating Rates

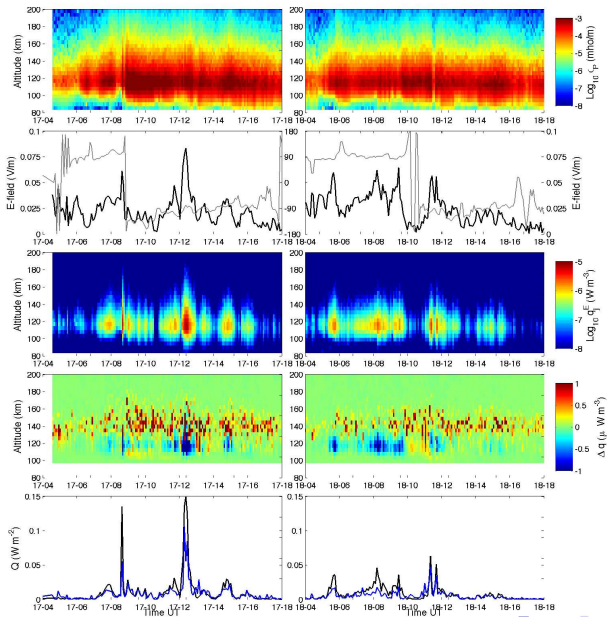


- Supplied electromagnetic energy is converted to heat (Joule heating) or momentum (Lorentz forcing)
- $q_j = \mathbf{j} \cdot \mathbf{E}' = \mathbf{j} \cdot (\mathbf{E} + \mathbf{U} \times \mathbf{B}) = \sigma_P |\mathbf{E} + \mathbf{U} \times \mathbf{B}|^2$
- $q_m = \mathbf{U} \cdot (\mathbf{j} \times \mathbf{B})$
- $q_m$  negative - neutral gas is a sink (load)
- $q_m$  positive - neutral gas is a source (generator)

Thayer [2000]

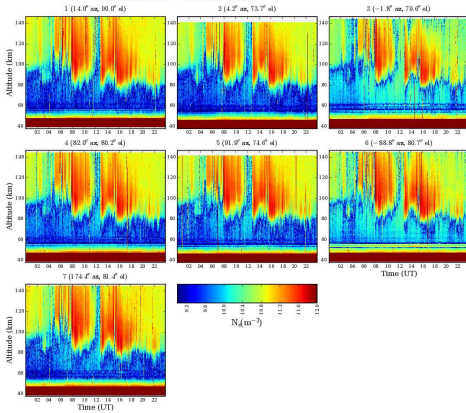


# Conductivities / Currents / Joule Heating Rates

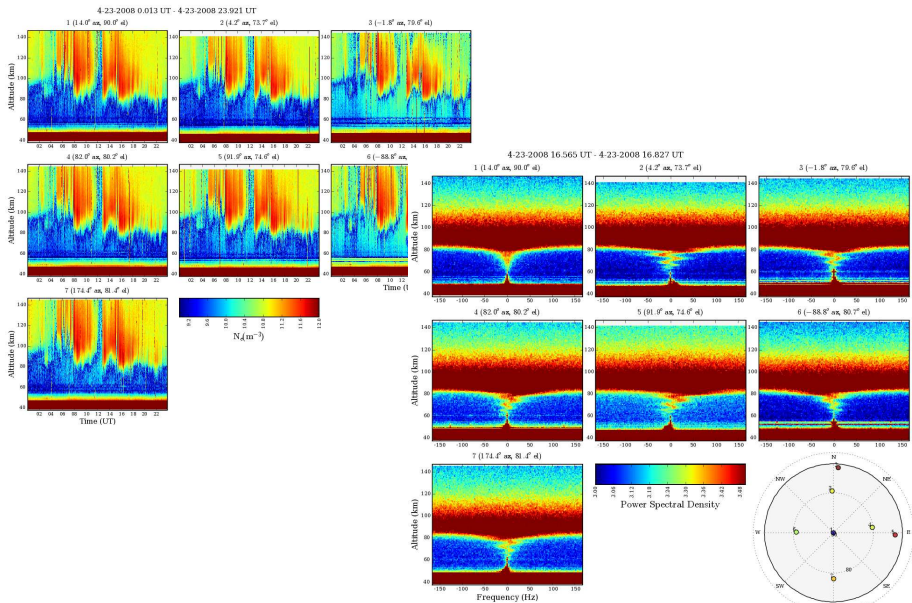


# D-Region Parameters - Raw Power and Spectra

4-23-2008 0.013 UT - 4-23-2008 23.921 UT

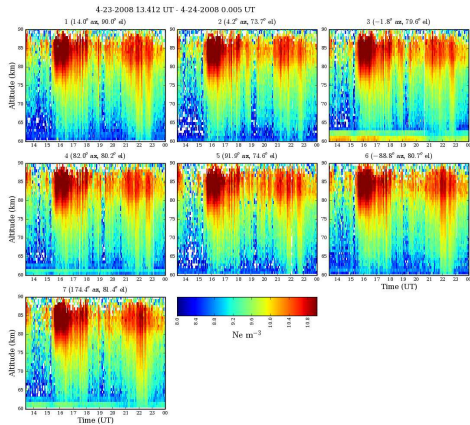


# D-Region Parameters - Raw Power and Spectra

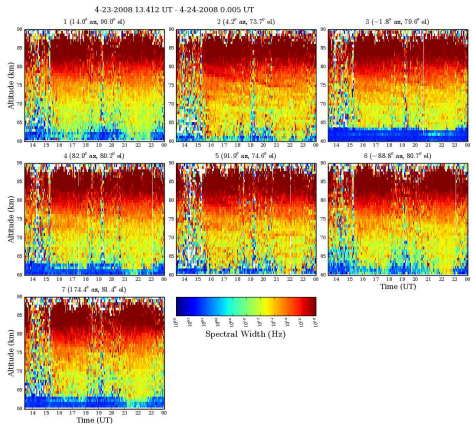
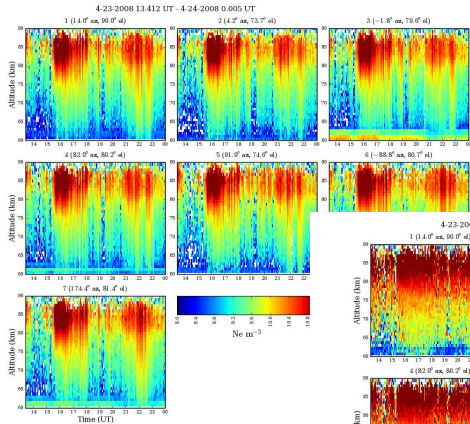




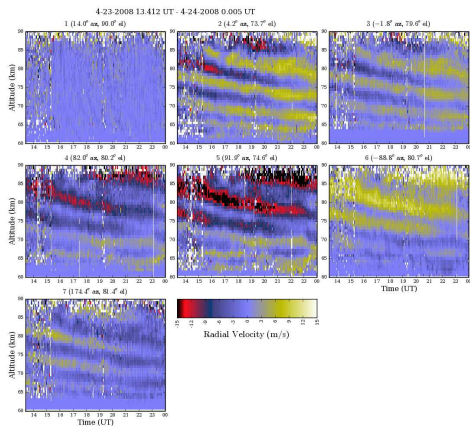
# D-Region Parameters - $N_e$ and Spectral Widths



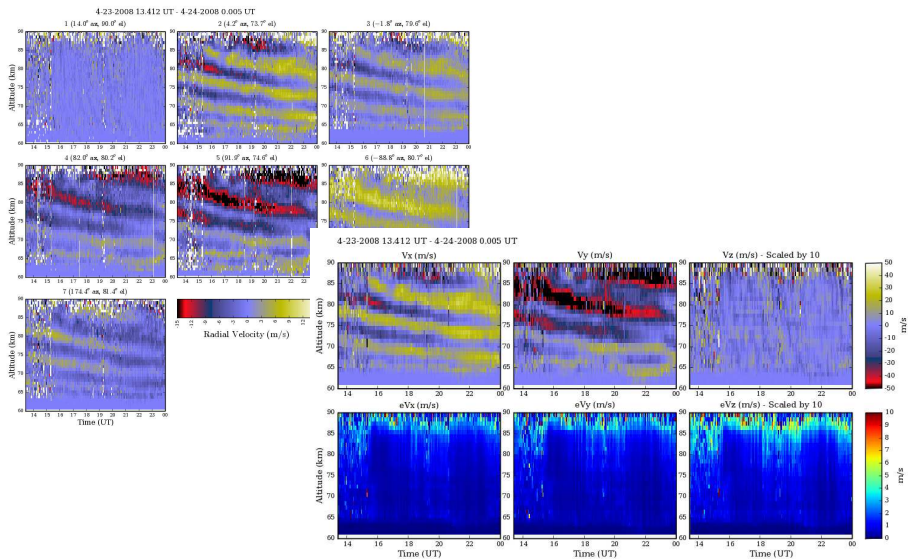
# D-Region Parameters - $N_e$ and Spectral Widths



# D-Region Parameters - Velocities and Winds



# D-Region Parameters - Velocities and Winds



# Future

- 1 Move towards full profile techniques
- 2 Take advantage of space and time information
- 3 Standardize approaches
- 4 Molecular ion composition, height-resolved plasma lines, topside parameters, etc.
- 5 Make these products available to interested users
- 6 Extend our arsenal of products (e.g., *D*-region momentum fluxes, higher altitude winds, etc.)